

A DIFFUSION MODEL OF AN
ISOTHERMAL TUBULAR FLOW REACTOR

by

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INTRODUCTION

The mechanism by which heat or mass is dispersed in tubular flow reactors, and other chemical process equipment through which a fluid is flowing, has recently gained a great deal of attention. In order to solve the problem of determining the conversion and mixing characteristics of a tubular flow reactor it is necessary to consider a rather complicated system of partial differential equations. The system is complicated because the velocity profile is not flat and because there are not only gradients of concentration but also of temperature and pressure. To obtain the general solution of this problem analytically is, of course, out of the question and the numerical solution is tedious and requires a very large computer. At the same time, however, there is the larger question of whether from a practical point of view the effects mentioned above are important enough to justify the effort needed to solve the problem. Therefore, it would be interesting to obtain some estimate of the magnitude of some of these effects.

In an attempt to exploit a technique which might give a partial insight into the behavior of tubular flow reactors, a flat velocity profile and a concentration gradient due to dispersion only is assumed. Earlier investigators introduced the concept of longitudinal diffusion to describe the phenomenon of dispersion in the direction of flow of the mean fluid stream, and the concept of radial diffusion to describe dispersion in the direction perpendicular to the mean fluid flow. This work will be limited to the study of axial dispersion.

A number of models have been proposed for the dispersion phenomenon associated with flow through tubular flow reactors, fluidized beds and packed beds, for example, the researches of Danckwerts (2) and Levenspiel and Smith (12) for flow through tubular flow reactors by the diffusion-type model, the work of Carberry and Bretton (1) for flow through fixed beds by the cell-model, the work of Miyauchi (14) by a cell with back-mixing, and the diffusion-dead space model by Koump (9) and Turner (20). All of these models idealize the three-dimensional problem into a one-dimensional transport equation.

The most commonly used model is the diffusion-type model. This model assumes that the overall rate of axial dispersion is proportional to the bulk concentration gradient of the diffusing component in the axial direction and is measured by the spreading of a concentration wave in an axial direction. This model is also used in this research to investigate the effect of axial mixing on the performance of isothermal tubular flow reactors. With this model, a boundary value problem of a second-order partial differential equation is obtained when a material balance is made with respect to a traceable component in the fluid.

In general, a change in boundary conditions may have a significant effect on the behavior of the solution of a partial differential equation. Hence, in order to avoid the possibility of error due to the improper choice of boundary conditions, one must prove that the boundary conditions used in solving a particular initial or boundary value problem are realistic, or that the solution of the particular differential equation is not sensitive to the choice of boundary conditions.

Various sets of boundary conditions have been suggested for use with the diffusion-type model to predict dispersion phenomenon associated with flow through tubular flow reactor. Danckwerts (2), for example, introduced a set of boundary conditions which satisfy the continuity requirement at both ends of a reactor. Many other investigators introduced sets of boundary conditions which do not conserve the continuity requirement at the inlet of reactor. Hulburt (8), for example, introduced a set of boundary conditions which neglect the diffusion at the inlet of the reactor. Levenspiel and Smith (12) further approximated a reactor of finite length by a reactor of infinite length. Although both of these two sets of boundary conditions are convenient for estimating the relative importance of reaction and diffusion rates, there has been no attempt made to justify numerically the quantitative relationships among the solutions obtained by use of these sets of boundary conditions. Another possible set of boundary conditions which meet the continuity requirement at both ends of the reactor, but hitherto unknown, is investigated in the course of this research.

The present study represents an effort to increase the theoretical knowledge of isothermal tubular flow reactors by investigating the effects of axial dispersion on the performance of a reactor based on a diffusion-type model. The differences and the relationships among the steady-state response curves against the various types of concentration perturbation obtained by use of the different sets of boundary conditions are numerically evaluated. Theoretical response curves of exit concentrations against impulse, step input and frequency response functions are presented in order to illus-

trate the method of determining mean dispersion coefficients from available dynamic experimental data.

LITERATURE SURVEY

As indicated in the previous section, different models of fluid flow and various sets of boundary conditions have in recent years been proposed in order to predict the effects of axial mixing on the performance of such process equipment as packed beds, fluidized beds and tubular flow reactors. Table 1 summarizes the work of previous investigators. Included in the table are the mathematical models and boundary conditions employed, the process equipment and experimental technique used, and the method of analysis of the data.

Danckwerts (2) presented a general discussion of how distribution functions for residence time can be defined and measured for pipes, packed beds and reactors from dynamic step-input experiments. He also reported the performance of a steady-state flow reactor based on the diffusion-type model with a first-order chemical reaction using a set of continuity preserving boundary conditions at both ends of the reactor. The effect of axial mixing on the first-order conversion is also compared with that of a plug flow reactor. The boundary conditions Danckwerts introduced will be denoted as B.C.(I).

Hulburt (8) investigated steady-state first-order kinetics using a diffusion model to study the effect of axial dispersion on the conversion. He presented a set of boundary conditions which neglect the effect of diffusion at the reactor entrance. This set of boundary conditions will be denoted as B.C.(II).

An analysis of these two boundary conditions was presented by Wehner and Wilhelm (24) for a steady-state flow reactor with axial diffusion and first-order reaction. They have indicated that the Hulburt solution (B.C.II) is convenient for estimating the relative

Table 1. Summary of previous investigations.

Investigators	Models	B.C.	Equipment	Experimental Technique	Method of Analysis
Danckwerts (2)	Diffusion-type	(I)	Reactors Blenders Packed tubes	Step input	
Kramers, et. al. (10)(11)	Diffusion-type	(I)(III)	Cylinder	Frequency response	
Levenspiel and Smith (12)	Diffusion-type	(III)	Reactors	No experiment	Variance, σ^2
Head, Hogen, and Walsh (7)	Diffusion-type	(III)	Tube	Impulse	Fourier Transforms
Ebach (4)	Diffusion-type	(III)	Packed beds	Frequency response; Pulse functions	Point source
Otake and Kunugita (16)	Diffusion-type	(II)	Packed beds	Step input	Approximate solution
Hulburt (8)	Diffusion-type	(II)	Reactors	Steady-state	
Deisler and Wilhelm (3)	Diffusion-type	(III)	Catalytic Reactors	Frequency response	Amplitude alone
Yagi and Miyauchi (25)	Diffusion-type	(I)	Reactor Fluidized bed	Impulse; Frequency response	Point source

Table 1. (con't.)

Investigators	Models	B.C.	Equipment	Experimental Technique	Method of Analysis
Strang and Geankoplis (18)	Diffusion-type	(II)	Packed bed	Frequency response	Amplitude alone
McHenry and Wilhelm (13)	Diffusion-type	(III)	Packed beds	Frequency response	Amplitude alone
Vincent (22)	Diffusion-type	(III)	Heat exchanger	Frequency response	
Romano and Pigford (17)	Diffusion-type	(III)	Fluidized beds	Frequency response	Fourier Transforms
Carberry and Bretton (1)	Cell-model	(III)	Fixed beds	Impulse	
Miyauchi (14)	Cell with back-mixing	General discussions		Step Impulse	
Turner (20)	Diffusion-dead space model	None	Packed beds	Frequency response	
Koump (9)	Diffusion-dead space model	(III)	Packed beds	Step input Frequency response Impulse	
von Rosenberg (23)	Diffusion-dead space model	(III)	Packed beds	Step input	

importance of reaction and dispersion rates, but is not consistent with the limit of a stirred tank reactor with infinite diffusivity in the reaction section. They have investigated the validity of these boundary conditions by simultaneous solution of three differential equations for the reaction section, fore and after sections. This investigation was based on the fact that axial dispersion may occur in the latter two sections.

Levenspiel and Smith (12) provided an analytical solution to the problem of fluid mixing in a continuous flow system without chemical reaction using the diffusion-type model but with a different set of boundary conditions. They presented a set of approximate boundary conditions by assuming that the effect of diffusion at the reactor entrance is negligible so that a discontinuity exists at this boundary. They further approximated a reactor with finite length by a reactor with infinite length. Based on this set of boundary conditions they showed two methods, a single point method (maximum concentration) and a method based on the variance of the concentration distribution (σ^2), for evaluating the axial dispersion coefficient from experimental measurements. They also worked out some examples to determine the conditions of applicability of the model. This set of boundary conditions will be denoted as B.C. (III).

In a short note, van der Laan (21) discussed an aspect of the work reported by Levenspiel and Smith (12). While Levenspiel and Smith (12) treated the diffusion-type approach by restricting the solutions to infinite lengths, van der Laan (21) treated the diffusional type flow in a more general manner so as to include the

case of finite length and of varying dispersion by applying suitable boundary conditions. He treated the general solution by means of a Laplace transformation and avoided the back transformation by applying momentum analysis, that is, the method of variance of the concentration distribution. An illustration was given for the evaluation of the dispersion coefficient using the results of calculations of mean residence time and variance.

One of the pioneer investigations of axial dispersion in process equipment was that of Kramers and Alberda (10). They investigated axial dispersion for the flow of water through a bed of Raschig rings. A frequency response technique was used, and the results were analyzed on the basis of the assumption that the rate of axial dispersion can be described by a diffusion-type model with B.C.(I). In order to avoid the complexities involved in obtaining an exact analytical solution, they have suggested that the performance function of the system be approximated by n completely mixed vessels connected in series. The feasibility of this approximation is open to question, except for the system involving a small degree of dispersion. Kramers, Honig, and Croockweit (11) extended their work to the study of longitudinal dispersion in liquid flow by using B.C.(II). Actual derivation of the analytical solution, however, was carried out by means of B.C.(III).

Yagi and Miyauchi (25) investigated the operational characteristics of continuous flow reactors in which the reactants are mixing and described the dispersion phenomenon in terms of mean longitudinal dispersion applying the diffusion-type model and B.C.(I). An impulse function was used, and the results were analyzed by means

of a single point method (i.e., maximum points). The data of Ebach (4) on impulse experiments were analyzed on the same basis as those of Yagi and Miyauchi (24). It is known that the use of a single datum in drawing conclusions for an entire experiment may lead to serious error because of the probable experimental error. Although the data shown in the work of Ebach (4) gave close agreement between impulse experimental values and calculated values at low degree of dispersion, the feasibility of using peak values of their impulse data at high degree of dispersion leaves much to be desired.

In experiments with the frequency response technique, Ebach (4) used amplitude data alone to analyze the mean dispersion coefficient. He stated that the axial dispersion coefficient calculated from phase shift data gave poor accuracy and reproducibility. The reasons for this fundamental fault were clearly illustrated in his paper. Strang and Geankoplis (18) studied the axial dispersion of liquids in packed beds by the same frequency response technique. The diffusion model was used, and the data were also analyzed using amplitude data alone.

Deisler and Wilhelm (3) also employed the frequency response technique to investigate both pore diffusion within particles and longitudinal dispersion in the fluid between particles in the laminar flow of gas mixtures. A diffusion-type model was used, and the analytical solution was derived based on B.C.(III). The axial dispersion was evaluated simultaneously with the diffusivity in solid and the adsorption constant. Applying a similar model, boundary conditions and experimental technique, McHenry and Wilhelm (13) studied the axial dispersion of binary gas flowing through beds of

inert spherical particles. The results were evaluated using amplitude data alone.

A very interesting paper was published by Head, Hougen and Walsh (7) in their study of the axial dispersion of water flowing through a vertical tube. A diffusion-type model was assumed together with B.C.(III). Experimental data were obtained by the pulse wave technique, and the results were analyzed by reducing pulse data to the frequency response form by means of a Fourier transformation. The reduced data were plotted in a Bode diagram to evaluate the mean dispersion coefficient by fitting experimental data and calculated values. A similar method of analyzing experimental data was reported by Romano and Pigford (17) in their study of the performance of a fluidized bed by the frequency response technique. The data obtained were reduced to pulse form by a Fourier transformation, and the resulting residence time distribution curve was used to study the mechanism of gas mixing in a fluidized bed, and the effects of gas velocity and bed height on the distribution function.

Otake and Kunugita (15) presented a general discussion of mixing characteristics of reactors without chemical reaction based on a diffusion-type model. Analytical solutions corresponding to the different sets of boundary conditions were presented in their paper for various types of transient inputs. In another paper (16), they investigated the performance of irrigated packed towers using a diffusion-type model and B.C.(II). A step signal was used, and the data were analyzed under the assumption that when the degree of mixing is very small their analytical solution, which is in the form of an n -th term expansion, can be simplified to only the first term.

In the previous section, it was pointed out that many other mathematical models have been proposed to predict the behavior of continuous flow systems. One of the most widely used models beside the diffusion-type model is the diffusion-dead space model. This model for axial dispersion postulates the existence of a stationary fluid region in the voids of the process equipment which can exchange mass with the moving fluid. In a purely theoretical investigation, Turner (20) presented a model of a packed bed which can be classified under the diffusion-dead space category. He introduced a concept hydrodynamically equivalent to the packed bed, which consists of a number of identical channels through which the fluid is flowing. Each of these channels has a number of "pockets" of stationary fluid which are arranged along the main channel. The fluid within the "pockets" is regarded as stationary. The mass transfer between the flowing fluid and the "pockets", and the mass transfer within the "pocket" volume is assumed to take place by molecular diffusion. The model proposed by Turner (20) appears to be too complex to be of practical use. In an attempt to simplify the model, Koump (9) presented the diffusion-dead space model under the assumptions that the axial dispersion of mass takes place by a combination of two mechanisms: a diffusion type process and a mass exchange between the moving fluid and the fluid within the stagnant void volume. The experiment was carried out by three different techniques: frequency response, impulse, and step inputs. The results were analyzed by both the diffusion-type model and the diffusion-dead space model.

Von Rosenberg (23) used the diffusion-dead space model to investigate the problem of the displacement of a fluid from a bed of inert solid by another fluid. He regarded the packed bed as an assembly of capillaries and used the solution of G. I. Taylor (19) for the case of negligible forward and rapid lateral diffusion as a basis of the analysis. He assumed that the axial dispersion takes place by the combined effects of velocity distribution in the pores and rapid diffusion in the radial direction. The concept of an axial dispersion coefficient, therefore, does not enter the picture in this approach.

Other models which are often used in the study of axial mixing are the cell-model and cell with back-mixing. The concept of the cell-model implies that the spread of injected tracer as a result of passage through a bed can be described by the equation of concentration-time behavior in the effluent of the n -th perfect mixing cell. The cell with back-mixing differs from the cell-model in that each one of the n perfect mixing cells is associated with a definite back-mixing.

Carberry and Bretton (1) investigated the longitudinal dispersion during the flow of water through beds of rings by the cell-model. Two different types of input signals were used in their work: an impulse function which injected the tracer instantaneously, and a rectangular input function which injected the tracer within a small given time. In both cases, the results were analyzed by fitting the shape of the experimental response curves to the calculated response curves based on the cell-model. Miyauchi (14) presented a purely theoretical analysis of the cell with back-mixing model.

It seems evident from this literature survey that the diffusion-type model has been employed almost exclusively because, to a first approximation, it gives a reasonable description of the phenomenon associated with mixing inside reactors. There seems to be, however, some disagreement on the method of analyzing experimental data and on which sets of boundary conditions to apply. All of the mathematical solutions are most certainly oversimplified and unrealistic from the standpoint of representing the true mechanism. Due to the simplification involved in the mathematical treatment on B.C.(III), it has been used most often without much quantitative justification.

THEORY

The purpose of this section is to derive the partial differential equation for the diffusion-type model of axial dispersion considered in this investigation, and to present the mathematical expressions for four sets of boundary conditions.

Diffusion Model

This model assumes that the rate of dispersion is proportional to the bulk concentration gradient of the diffusing component in the tubular reactor. Under this assumption, a material balance with respect to a traceable component in the fluid with chemical reaction occurring gives the following conservation of mass equation in vector form:

$$\frac{\partial C}{\partial t} + \text{div}(-D \text{ grad. } C) + \text{div}(\vec{u}C) + r(C) = 0 \quad (1)$$

where D is a dispersion coefficient or apparent diffusivity which characterizes the dispersion of fluids flowing with a velocity of \vec{u} , and C and t represent the concentration and time variables respectively. The first term of equation (1) represents the change in concentration with respect to time. The second term represents the net out-flow of component due to dispersion. The third term represents the velocity gradient while the fourth term represents the rate of chemical reaction.

The following assumptions are made:

1) Only the dispersion in the axial direction, z , is significant, that is, the physically three-dimensional problem of the

second term in equation (1) is idealized to a one-dimensional transport equation.

2) The coefficient D , in equation (1) is independent of position and concentration gradients as in the case of fully developed turbulent flow, and is characterized by \bar{D} .

3) Changes in volume are assumed not to take place, so that the velocity, \bar{u} , in equation (1) is the same at all cross-sections, and is characterized by \bar{u} .

4) Only the chemical reaction of first-order is considered. With these simplifications, equation (1) becomes the following partial differential equation:

$$\frac{\partial C}{\partial t} = \bar{D} \frac{\partial^2 C}{\partial z^2} - \bar{u} \frac{\partial C}{\partial z} - kC \quad (2)$$

where k is the reaction rate constant.

Equation (2) is a second-order linear partial differential equation which describes the dispersion phenomenon in terms of time, position and dispersion coefficient. As stated in the previous section, the solution of equation (2) depends on the initial and boundary conditions.

Boundary Conditions

Various boundary conditions have been suggested for use with equation (2). As stated in the previous section, boundary conditions may have a significant effect on the solution of equation (2). Hence, the boundary conditions used with equation (2) must be realistic and quantitatively justified. One of the purposes of this investigation is to find boundary conditions compatible with the solution

of equation (2).

Danckwerts (2) introduced a set of intuitive boundary conditions based on the equation of continuity. He proposed that, owing to the dispersion taking place, the concentration of reactant in the entering stream, C_1 , must be larger than the concentration just within the entrance of the reactor.

This was expressed mathematically as:

$$\bar{u}C_1 = \bar{u}C_{z \rightarrow 0^+} - \bar{D} \left(\frac{\partial C}{\partial z} \right)_{z \rightarrow 0^+} \quad (3)$$

At the outlet of a finite tubular reactor, he proposed the same relationship as equation (3):

$$\bar{u}C_{z \rightarrow L^-} = \bar{u}C_{z \rightarrow L^+} - \bar{D} \left(\frac{\partial C}{\partial z} \right)_{z \rightarrow L^+} \quad (3a)$$

The equation of continuity, however, necessitates equation (3a) to be expressed as:

$$\bar{u}C_{z \rightarrow L^-} = \bar{u}C_{z \rightarrow L^+} \quad (3b)$$

Substitution of equation (3b) into equation (3a) yields:

$$\bar{D} \left(\frac{\partial C}{\partial z} \right)_{z \rightarrow L^+} = 0 \quad (3c)$$

Since \bar{D} is different from zero, the term in parenthesis of equation (3c) has to satisfy the following expression:

$$\left(\frac{\partial C}{\partial z} \right)_{z = L} = 0 \quad (4)$$

Equations (3) and (4) will be designated as B.C.(I).

In an attempt to approximate the above B.C.(I), Hulburt (8) proposed a set of boundary conditions which neglect the effect of diffusion at the reactor entrance. Under this assumption, the second term in the right-hand side of equation (3) vanishes, and the following boundary condition at the inlet of the reactor can be obtained:

$$\bar{u}C_1 = \bar{u}C_{z \rightarrow 0+} \quad (5)$$

At the outlet, $z = L$, he used the same boundary condition proposed by Danckwerts. Equations (5) and (4) will be designated as B.C.(II).

Levenspiel and Smith (12) and several other investigators (14) (15) proposed a set of approximate boundary conditions. They not only neglected the effect of diffusion directly within the entrance of the reactor, but also treated a finite length of reactor as an infinite reactor, i.e., one value of the dispersion coefficient applied to the entire length and cross-section. This infinite length reactor boundary condition is expressed mathematically as:

$$\lim_{z \rightarrow \infty} C(z, t) = 0 \quad (6)$$

Equations (5) and (6) will be designated as B.C.(III).

Another set of boundary conditions which has never been employed previously is presented in this investigation. This set of boundary conditions preserves continuity just as those of Danckwerts. The only approximation made in this set of boundary conditions is that the finite length of reactor is treated as a section of an infinitely long reactor as in the case of B.C.(III). The feasibility

of this set of boundary conditions will be discussed in a later section in connection with steady-state first-order conversion and the response to various transient inputs. Equations (3) and (6) will be designated as B.C.(IV).

Table 2 summarizes the proposed boundary conditions for use with equation (2).

Table 2. Summary of proposed boundary conditions.

Boundary Conditions	Inlet ($z = 0$)	Outlet ($z = L$)
B.C.(I)	$\bar{u}C_1 = \bar{u}C_{z \rightarrow 0^+} - \bar{D}\left(\frac{\partial C}{\partial z}\right)_{z \rightarrow 0^+}$	$\left(\frac{\partial C}{\partial z}\right)_{z = L} = 0$
B.C.(II)	$\bar{u}C_1 = \bar{u}C_{z \rightarrow 0^+}$	$\left(\frac{\partial C}{\partial z}\right)_{z = L} = 0$
B.C.(III)	$\bar{u}C_1 = \bar{u}C_{z \rightarrow 0^+}$	$\lim_{z \rightarrow \infty}(z, t) = 0$
B.C.(IV)	$\bar{u}C_1 = \bar{u}C_{z \rightarrow 0^+} - \bar{D}\left(\frac{\partial C}{\partial z}\right)_{z \rightarrow 0^+}$	$\lim_{z \rightarrow \infty}(z, t) = 0$

MATHEMATICAL TREATMENT

The solutions of equation (2) which satisfy the four sets of boundary conditions tabulated in Table 2 of the previous section are derived by means of Laplace transformations. This has the advantage of transforming a partial differential equation into an ordinary linear differential equation. To derive the expressions for the exit concentrations against various types of transient inputs, the back transformation is carried out for each of the solutions by means of tables of integral transforms (5) or the Heaviside partial fraction theorem.

This section will be divided into four parts to show how the solutions are obtained for both ideal and non-ideal flow in steady-state and unsteady-state isothermal tubular flow reactors. Only the final results will be given here, and the details of the derivations will be given in Appendix I, II, III, and IV.

Steady-state Non-ideal Flow

The differential equation for the steady-state tubular flow reactor based on the diffusion-type model with axial dispersion and first-order chemical reaction can be obtained from equation (2) by letting the accumulation be zero:

$$0 = D \frac{d^2 C}{dz^2} - \bar{u} \frac{dC}{dz} - kC \quad (7)$$

The following transformations of the variables in Equation (7) lead to the dimensionless form of the differential equation:

$$\frac{d^2 \gamma}{d\eta^2} - 2M \frac{d\gamma}{d\eta} - 2MR\gamma = 0$$

(8)

where $\gamma = \frac{C}{C_1}$; $\eta = \frac{z}{L}$; $M = \frac{\bar{u}L}{2\bar{D}}$ and $R = \frac{kL}{\bar{u}}$

The four sets of boundary conditions tabulated in Table 2 of the previous section are transformed likewise and are summarized in Table 3.

Table 3. Transformed forms of boundary conditions.

Boundary Conditions	Inlet ($\eta = 0$)	Outlet ($\eta = 1$)
B.C.(I)	$\frac{d\gamma}{d\eta} = 2M [\gamma(0^+) - 1]$	$\frac{d\gamma}{d\eta} = 0$
B.C.(II)	$\gamma = 1$	$\frac{d\gamma}{d\eta} = 0$
B.C.(III)	$\gamma = 1$	$\lim_{\eta \rightarrow \infty} \gamma = 0$
B.C.(IV)	$\frac{d\gamma}{d\eta} = 2M [\gamma(0^+) - 1]$	$\lim_{\eta \rightarrow \infty} \gamma = 0$

The details of the solutions of equation (2) which satisfy the above boundary conditions (Table 3) will be given in Appendix I. The results obtained for each of the above four boundary conditions are respectively:

$$\gamma_I(\eta) = 2 \left\{ \frac{(1+\beta)\exp[M(1+\beta) + M(1-\beta)\eta] - (1-\beta)\exp[M(1+\beta)\eta + M(1-\beta)]}{(1+\beta)^2 \exp[M(1+\beta)] - (1-\beta)^2 \exp[M(1-\beta)]} \right\} \quad (9)$$

$$\gamma_{II}(\eta) = \frac{(1+\beta)\exp[M(1+\beta) + M(1-\beta)\eta] - (1-\beta)\exp[M(1+\beta)\eta + M(1-\beta)]}{(1+\beta)\exp[M(1+\beta)] - (1-\beta)\exp[M(1-\beta)]} \quad (10)$$

$$\gamma_{III}(\eta) = \exp[M(1-\beta)\eta] \quad (11)$$

$$\gamma_{IV}(\eta) = \frac{2\exp[M(1-\beta)\eta]}{1+\beta} \quad (12)$$

where $\beta = \sqrt{1 + \frac{2R}{M}}$

Steady-state Ideal Flow

In practice neither plug flow nor completely mixed flow will occur; there always will be some degree of dispersion due to viscous effects and molecular or eddy diffusion. In order to establish the theoretical limitations involved with regard to the degree of mixing, it is necessary to derive the equations for the ideal flow reactor.

For the completely mixed flow reactor, where $D = \infty$ and $M = 0$, and if the effective volume of the reacting mixture, V , and the volumetric flow rate, v , are constant, the material balance is the same as for a single uniformly stirred tank reactor and may be written as:

$$V \frac{dC}{dt} = vC_1 - vC - kC^n \quad (13)$$

where n is the order of reaction.

For the calculation of the steady-state concentration profile and the conversion for the case of a first-order chemical reaction, equation (13) may be simplified.

$$0 = \frac{v}{V} C_1 - \frac{v}{V} C - \frac{k}{V} C \quad (14)$$

Equation (14) can be transformed into a dimensionless form by the substitutions used for (8) to obtain:

$$R\gamma' + \gamma' - 1 = 0 \quad (15)$$

or

$$\gamma' = \frac{1}{1 + R} \quad (16)$$

For the steady-state plug flow reactor, where $D = 0$ and $M = \infty$, the first term in the left-hand side of equation (8) vanishes to give

$$\frac{d\gamma'}{d\eta} = -R\gamma' \quad (17)$$

The solution of equation (17) is:

$$\gamma'(\eta) = e^{-R\eta} \quad (18)$$

Because of the theoretical limits, which were pointed out early in this section, all the concentrations and conversion calculated from equations (9) through (12) have to lie between the results of equations (16) and (18), which give the minimum and maximum values respectively.

Unsteady-state Non-ideal Flow

When the deviation of the concentration of the traceable component from the steady-state value due to a perturbation of inlet concentration, $C_1(t)$, only is considered, the initial condition may be set as $C(z,0) = 0$. The transfer functions, $W(s)$, defined as the ratio between the Laplace transforms of the inlet and the outlet concentrations, using equation (2) for each set of four boundary conditions are shown in equations (19) through (22). The details of the derivations are shown in Appendix II.

$$W(s) = \frac{4\beta'}{(1+\beta')^2 \exp[-M(1-\beta')] - (1-\beta')^2 \exp[-M(1+\beta')]} \quad (19)$$

$$W(s) = \frac{2\beta'}{(1+\beta') \exp[-M(1-\beta')] - (1-\beta') \exp[-M(1+\beta')]} \quad (20)$$

$$W(s) = \exp[M(1-\beta')] \quad (21)$$

$$W(s) = \frac{2\exp[M(1-\beta')]}{1+\beta'} \quad (22)$$

The basic forms of the inlet concentration variations, $C_1(t)$, and their Laplace transforms, $C_1(s)$, for impulse and step input are:

	<u>$C_1(t)$</u>	<u>$C_1(s)$</u>
Impulse	$C_1 u^{\delta}(t)$	C_1
Step Input	$C_1 u(t)$	C_1/s

The transforms of the exit concentration, $C_2(L,s)$, are then obtained systematically as the products of the transfer functions, $W(s)$, and the $C_1(s)$ as follows:

Impulse	$C_2(s) = C_1 W(s)$
Step Input	$C_2(s) = C_1 W(s)/s$

The inverse transforms of $C_2(s)$ to obtain the response of exit concentration, $C_2(t)$, against the variation of feed concentration, $C_1(t)$, may be carried out systematically by means of the Heaviside partial fraction theorem. The results are tabulated in Table 4 for each of the boundary conditions. The details of the derivations are given in Appendix III.

Table 4. Response of exit concentration from an isothermal tubular flow reactor.

Transient
Inputs

B.C.(I)

Impulse

$$\frac{C_2^V}{Q} = \sum_{n=1}^{\infty} \frac{2 \delta_n (M \sin \delta_n + \delta_n \cos \delta_n)}{(M^2 + 2M + \delta_n^2)} \times \text{Exp} \left[M - \left(\frac{M^2 + \delta_n^2 + 2MR}{2M} \right) \tau \right] \quad (23)$$

Step
Input

$$\frac{C_2(t)}{C_1} = A - 4 \sum_{n=1}^{\infty} \frac{M \delta_n (M \sin \delta_n + \delta_n \cos \delta_n)}{(M^2 + 2M + \delta_n^2)(M^2 + \delta_n^2 + 2MR)} \times \text{Exp} \left[M - \left(\frac{M^2 + \delta_n^2 + 2MR}{2M} \right) \tau \right] \quad (27)$$

Frequency
Response

$$A. R. = \frac{b e^M}{(\alpha_1^2 + \alpha_2^2)^{\frac{1}{2}}} \quad \phi = - \tan^{-1} \left(\frac{\alpha_2}{\alpha_1} \right) + a \quad (31)$$

Table 4. (Cont.)

Transient
Inputs

B.C.(II)

Impulse

$$\frac{C_2 V}{Q} = \sum_{n=1}^{\infty} \frac{-\delta_n'^2}{(M + M^2 + \delta_n'^2) \cos \delta_n'}$$

$$\times \text{Exp} \left[M - \left(\frac{M^2 + \delta_n'^2 + 2MR}{2M} \right) \tau \right] \quad (24)$$

Step
Input

$$\frac{C_2(t)}{C_1} = B - 2 \sum_{n=1}^{\infty} \frac{\delta_n' \sin \delta_n' (M^2 + \delta_n'^2)}{(M^2 + \delta_n'^2 + M) (M^2 + \delta_n'^2 + 2MR)}$$

$$\times \text{Exp} \left[M - \left(\frac{M^2 + \delta_n'^2 + 2MR}{2M} \right) \tau \right] \quad (28)$$

Frequency
Response

$$A. R. = \frac{M}{(\alpha_3^2 + \alpha_4^2)^{\frac{1}{2}}}$$

$$\phi = a - \text{Tan}^{-1} \left(\frac{\alpha_4}{\alpha_3} \right) \quad (32)$$

Table 4. (Cont.)

Transient
Inputs

B.C.(III)

Impulse

$$\frac{C_2 V}{Q} = \left(\frac{M}{2\pi} \right)^{\frac{1}{2}} (\mathcal{C})^{-1.5} \\ \times \text{Exp} \left[M - \frac{1}{2} \left(M \mathcal{C} + \frac{M}{\mathcal{C}} + 2R \mathcal{C} \right) \right] \quad (25)$$

Step
Input

$$\frac{C_2(t)}{C_1} = \frac{1}{2} \left\{ e^{M+(M^2 + 2MR)^{\frac{1}{2}}} \right. \\ \times \text{Erfc} \left[\frac{1}{2} \left(\frac{2M}{\mathcal{C}} \right)^{\frac{1}{2}} + \frac{1}{2} (2M \mathcal{C} + 4R \mathcal{C})^{\frac{1}{2}} \right] \\ \left. + e^{M-(M^2 + 2MR)^{\frac{1}{2}}} \right. \\ \times \text{Erfc} \left[\frac{1}{2} \left(\frac{2M}{\mathcal{C}} \right) - \frac{1}{2} (2M \mathcal{C} + 4R \mathcal{C})^{\frac{1}{2}} \right] \left. \right\} \quad (29)$$

Frequency
Response

$$A. R. = e^{M-x}$$

$$\phi = -y \quad (33)$$

Table 4. (Cont.)

Transient
Inputs

B.C. (IV)

Impulse

$$\frac{C_2^V}{Q} = \left(\frac{2M}{\pi} \right)^{\frac{1}{2}} (\tau)^{-\frac{1}{2}} \times \text{Exp} \left[M^{-\frac{1}{2}} \left(M\tau + \frac{M}{\tau} + 2R\tau \right) \right] \\ - M \text{Exp}(2M - R\tau) \text{Erfc} \left[\left(\frac{M}{2\tau} \right)^{\frac{1}{2}} + \left(\frac{M\tau}{2} \right)^{\frac{1}{2}} \right] \quad (26)$$

Step
Input

$$\frac{C_2(t)}{C_1} = \int_0^t \left\{ \left(\frac{2M}{\pi} \right) (\lambda)^{-\frac{1}{2}} \text{Exp} \left[M^{-\frac{1}{2}} \left(M\lambda + \frac{M}{\lambda} - 2R\lambda \right) \right] \right. \\ \left. - M \text{Exp}(2M - R\tau) \text{Erfc} \left[\left(\frac{M}{2\tau} \right)^{\frac{1}{2}} + \left(\frac{M\tau}{2} \right)^{\frac{1}{2}} \right] \right\} d\lambda \quad (27)$$

Frequency
Response

$$A. R. = \frac{2 e^{M-x}}{\left[(1 + x/M)^2 + \sin^2 a \right]^{\frac{1}{2}}} \\ \phi = -y - \text{Tan}^{-1} \left(\frac{\sin a}{1 + x/M} \right) \quad (34)$$

The amplitude ratios and phase shifts of the frequency response functions caused by a sinusoidal input of amplitude C_1 and angular frequency ω , are obtained by means of a Fourier transformation. Thus, a performance function, defined as the ratio of the Fourier transform of the output to that of the input, is obtained by replacing s by $j\omega$ in equations (19) through (22) for each of the boundary conditions respectively, where $j = \sqrt{-1}$. From these performance functions, the computable forms of the amplitude ratios and phase shifts are obtained by changing into polar coordinates. These are also summarized in Table 4. The details of derivations will be given in Appendix IV.

Unsteady-state Ideal Flow

For a completely mixed reactor, where $D = \infty$ and $M = 0$, the Laplace transform of equation (13) with respect to time with zero initial concentration gives a transfer function of the following form:

$$W(s) = \frac{1}{1 + (s+k)\theta} \quad (35)$$

Substitution of the corresponding transient inputs and their Laplace transforms into equation (35) and inverse transformation of the resulting equations gives:

$$\text{Impulse} \quad \frac{C_2(t)V}{Q} = e^{-(1+R)\tau} \quad (36)$$

$$\text{Step Input} \quad \frac{C_2(t)}{C_1} = \frac{1}{1+R} \left[1 - e^{-(1+R)\tau} \right] \quad (37)$$

Frequency Response

$$[P.F.]_{M=0} (j\omega) = \frac{1}{(1+R)^2 + (\omega\theta)^2} e^{-j\tan^{-1} \left(\frac{\omega\theta}{1+R} \right)} \quad (38)$$

where P.F. = the performance function.

Q = quantity of tracer injected in the impulse function.

V = effective volume of reactor

$$\tau = t/\theta$$

For the unsteady-state plug flow reactor, where $\bar{D} = 0$ and $M = \omega$, the first term on the right-hand side of equation (2) vanishes, so that

$$\frac{\partial C}{\partial t} = -\bar{u} \frac{\partial C}{\partial z} - kC \quad (39)$$

The Laplace transform of equation (39) with respect to time with zero initial concentration gives a transfer function of the following form:

$$W(s) = e^{-R-s\theta} \quad (40)$$

Substitution of the corresponding transient inputs and their Laplace transforms into equation (40) and inverse transformation of the resulting equations gives:

$$\text{Impulse } \frac{C_2(t) V}{Q} = e^{-R} \delta(\tau - \tau_0) \quad (41)$$

$$\text{Step Input } \frac{C_2(t)}{C_1} = \begin{cases} 0 & 0 < \tau < 1 \\ e^{-R} & \tau > 1 \end{cases} \quad (42)$$

$$\text{Frequency Response } [P.F.]_{M=\infty} (j\omega) = e^{-R} e^{-j(\omega\theta)} \quad (43)$$

NUMERICAL COMPUTATION

The numerical computations for steady-state solutions represented by equations (9) through (12), (16) and (18) are straight forward and those for the frequency response solutions represented by equations (31) through (34), (38) and (43) are easy to perform by use of a computer, and no elaboration is required. This section is primarily concerned with the computations for unsteady-state transient inputs, impulse and step inputs, represented by equations (23), (24), (27) and (28). All of these four equations are in series form which requires that a tolerable error be established such that the error term of the series does not exceed 10^{-5} .

The first step in calculating numerical values of these equations was to determine the method of convergence of the equations:

$$\cot \delta = \frac{1}{2} \left(\frac{\delta}{M} - \frac{M}{\delta} \right) \quad \text{for } \delta_n \quad (44)$$

$$\tan \delta' = + \frac{\delta'}{M} \quad \text{for } \delta'_n \quad (45)$$

For this purpose, the Newton-Raphson iterative formula was used, namely,

$$\begin{aligned} \delta_n^{[q+1]} &= \delta_n^{[q]} - \frac{f(\delta_n)}{f'(\delta_n)} \\ &= \delta_n^{[q]} + \frac{\cot \delta_n^{[q]} - \frac{1}{2} \frac{\delta_n^{[q]}}{M} - \frac{M}{\delta_n^{[q]}}}{1 + \cot^2 \delta_n^{[q]} + \frac{1}{2} \left(\frac{1}{M} + \frac{M}{\delta_n^{[q]2}} \right)} \end{aligned} \quad (46)$$

with a first approximation of the following:

$$\begin{aligned}\delta_1^{[0]} &= 2 \\ \delta_n^{[0]} &= \delta_{n-1}^{[0]} + \pi \quad n = 2, 3, 4, \dots\end{aligned}\tag{47}$$

Equation (47) holds for all values of M .

The n -th term of equations (23), (24), (27) and (28) were then calculated with the δ and δ' thus obtained from equations (44) and (45), and the infinite series is truncated after the n -th term such that $\left| \left(\frac{C_{2V}}{Q} \right)_n / \sum_{n=1}^n \left(\frac{C_{2V}}{Q} \right)_n \right|$ is smaller than the maximum tolerable error of 10^{-5} .

The block diagrams and computer programs for use with the IBM 650 computers are presented in Table 5 for equations (23), (24) and (44) and in Table 6 for equations (27), (28) and (45) for $R = 0$ only.

Table 5 - a Block diagram for use with IBM 650 for computing equations (23), (24) and (44).

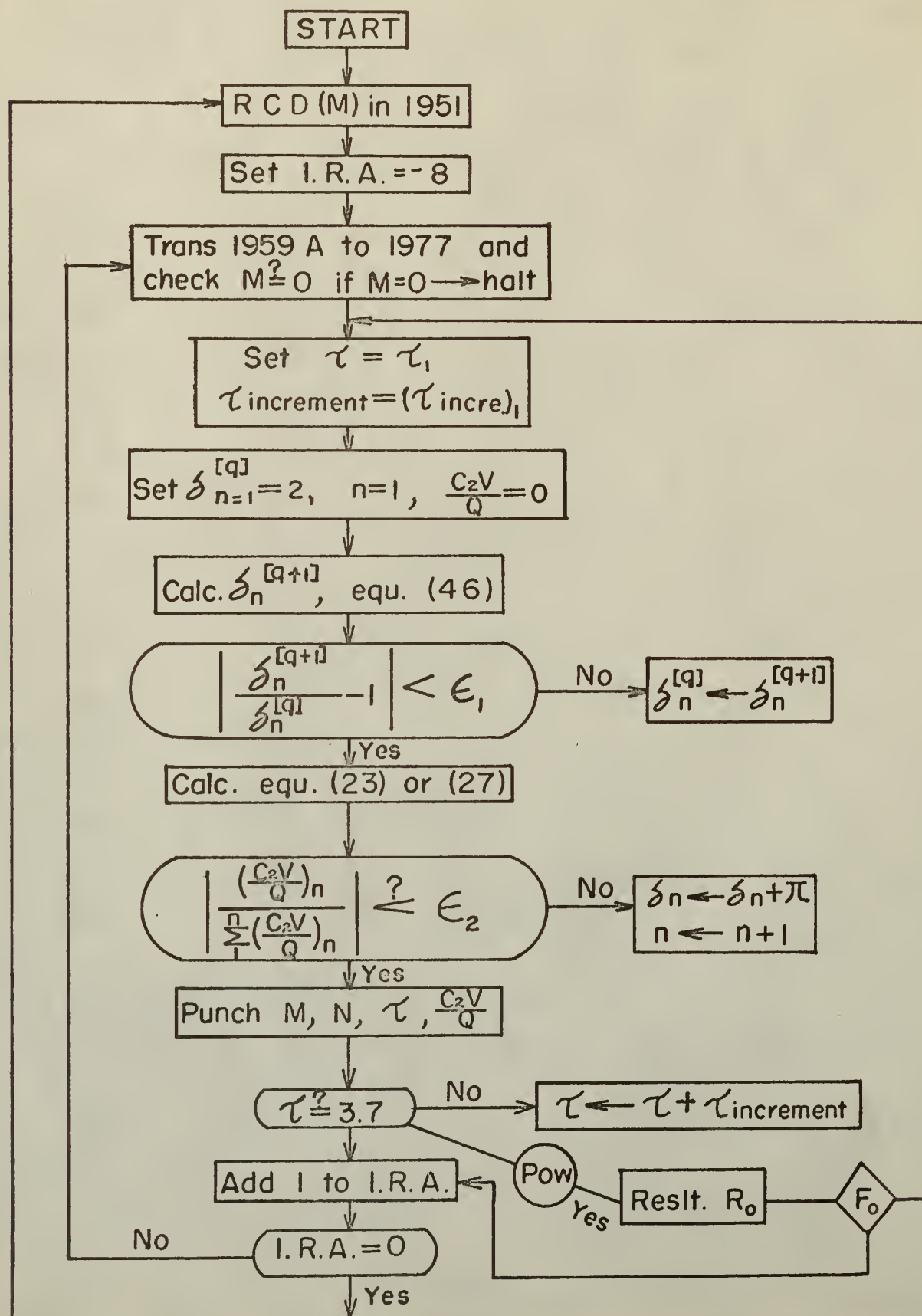


Table 5-b. Computer programs for use with IBM 650 for computing equations (23) (24) and (44)

MODIF	LDD	1582		1572	69	1582	1592
	STD	1570		1592	24	1570	1553
	32	1500	0700	1582	32	1500	0700
	LDD	1563		1553	69	1563	1573
	STD	1571	DEL17	1573	24	1571	0107
1563 = Number of R				1563	00	0003	0000
Program Step No. 202, Connect to 1570				0277	34	0250	1570
	FAD	1500	0700	1570	32	1500	0700
	RAU	1570		1560	60	1570	1580
	AUP	1590		1580	10	1590	1551
	STU	1570		1551	21	1570	1561
	RAU	1571		1561	60	1571	1581
	SUP	1590		1581	11	1590	1552
	STU	1571		1552	21	1571	1562
	NZU	NO	BACK to prog.	1562	44	0030	0777
1590 = Fixed point one				1590	00	0001	0000
	R			1500	00	0000	0000
	R			1501	10	0000	0050
	R			1502	10	0000	0051
	R			1503	50	0000	0051
	R			1504	10	0000	0052
	R			1505	50	0000	0052
	RSA	0003	MODIF	0151	81	0008	1572
Program step No. 255, Connect to 1560				0727	71	1877	1560

EXP01	STO EXP02	15	0000	24	0003	0006
	STU EXP03	16	0006	21	0010	0013
	LOO 8007	17	0013	69	8007	0019
	STO EXP04	18	0019	24	0022	0025
	RAC 0000	19	0025	88	0000	0031
	RAU EXP03	20	0031	60	0010	0015
	SMI EXP05	21	0015	46	0018	0069
EXP06	FSB LNTEN	22	0069	33	0072	0049
	SMI EXP07	23	0049	46	0002	0053
	AXC 0001	24	0053	58	0001	0069
EXP07	FAD LNTEN	25	0002	32	0072	0099
EXP05	SXC 0001	26	0018	59	0001	0024
	FAD LNTEN	27	0024	32	0072	0149
	SMI EXP05	28	0149	46	0018	0099
EXP08	STU EXP09	29	0099	21	0010	0063
	RAL EXP04	30	0063	21	0068	0021
	LDO 8007	31	0021	65	0022	0027
	STO EXP04	32	0027	69	8007	0033
	RAC 8002	33	0033	24	0022	0075
	LDO FPONE	34	0075	88	8002	0083
	STO EXP10	35	0083	69	0036	0039
	STO EXP11	36	0039	24	0042	0045
EXP12	STO EXP11	37	0045	24	0048	0001
	RAU EXP10	38	0001	60	0042	0047
	FAO EXP09	39	0047	32	0068	0095
	STU EXP10	40	0095	21	0042	0145
	RAU EXP11	41	0145	60	0048	0103
	FAO FPONE	42	0103	32	0036	0113
	STU EXP11	43	0113	21	0048	0051
	RAU EXP09	44	0051	60	0068	0023
	FMP EXP03	45	0023	39	0010	0060
	FOV EXP11	46	0160	34	0048	0098
	STU EXP09	47	0098	21	0068	0071
	FOV EXP10	48	0071	34	0042	0092
	FSB SIZE8	49	0092	33	0195	0121
	SMI	50	0121	46	0074	0001
	RAU EXP10	51	0074	60	0042	0097
	AUP EXP04	52	0097	10	0022	0003
FPONE	10 0000	53	0036	10	0000	0051
SIZE8	10 0000	54	0195	13	0000	0043
LNTEN	23 0258	55	0072	23	0258	5151
EQO0CB	STO CSV02	56	0050	24	0153	0056
	NZE CSV03	002	57	0056	45	0110
CSV04	RAU FPONE	003	58	0011	60	0036
CSV03	SMI CSV05	004	59	0110	46	0163
CSV05	FAO TWPDI	005	60	0163	32	0016
	NZE	006	61	0043	45	0046
	SMI CSV05	007	62	0046	46	0163
	FSB ONEPI	008	63	0100	33	0203
	NZE CSV07	009	64	0029	45	0032
CSV08	RSU FPONE	010	65	0133	61	0036
CSV06	FSB TWPDI	011	66	0014	33	0016
	NZE	012	67	0093	45	0096
	SMI	013	68	0096	46	0199
	FAD ONEPI	014	69	0199	32	0203
	NZE CSV07	015	70	0079	45	0032
CSV07	BTU CSV09	016	71	0032	21	0086
	RSU FPONE	017	72	0089	61	0036
	STU CSV10	018	73	0041	21	0146
	STU CSV11	019	74	0249	21	0004
	STL CSV12	020	75	0007	20	0061
EQO0CA	STO CSV02	021	76	0150	24	0153
	NZE	022	77	0160	45	0160
	SMI CSV14	023	78	0160	46	0213
CSV14	FAD TWPDI	024	79	0213	32	0016
	NZE	025	80	0143	45	0196
	SMI CSV14	026	81	0196	46	0213
	FSB ONEPI	027	82	0200	33	0203
	NZE CSV16	028	83	0129	45	0082
CSV15	FSB TWPDI	029	84	0144	33	0016
	NZE	030	85	0193	45	0246
	SMI	031	86	0246	46	0299
	FAD ONEPI	032	87	0299	32	0203
	NZE CSV16	033	88	0172	45	0082
CSV16	STU CSV09	034	89	0082	21	0086
	RSU 8003	035	90	0139	61	8003
	STU CSV10	036	91	0147	21	0146
	STU CSV11	037	92	0349	21	0004
	LOO FPONE	038	93	0189	69	0033
	STO CSV12	039	94	0189	24	0061
CSV13	RAU CSV13	040	95	0064	60	0061
	FAD FPONE	041	96	0065	32	0036
	BTU CSV17	042	97	0263	21	0118
	FAD FPONE	043	98	0171	32	0036
	STU CSV12	044	99	0313	21	0061
	RSU CSV10	045	100	0164	61	0146
	FMP CSV09	046	101	0101	39	0086
	STU CSV10	047	102	0136	39	0086
	FDV CSV11	048	103	0111	21	0146
	RAM 8003	049	104	0399	34	0004
	RAU 8002	050	105	0054	67	8003
	FSB SIZE8	051	106	0161	60	8002
	SMI	052	107	0119	33	0195
	RAU CSV11	053	108	0221	46	0124
	FAD CSV10	054	109	0125	60	0004
	STU CSV11	055	110	0009	32	0146
CSV18	RAU CSV11	056	111	0073	21	0004
	FPONE	057	112	0073	60	0004
FPONE	10 0000	058	113	0124	10	0000
FTW0	20 0000	059	114	0036	20	0000
TW0PI	62 8318	060	115	0250	62	8318
SIZE8	10 0000	061	116	0016	62	8318
ONEPI	31 4159	062	117	0016	31	4159
XSD05	50 0000	063	118	0203	31	4159
FP38	38 0000	064	119	0049	50	0000
FP01	10 0000	065	120	0300	38	0000
SIZEA	10 0000	066	121	0350	10	0000
SIZE8	10 0000	067	122	0400	10	0000
EQNO1	LDO LINKA	068	123	0450	10	0000
	STO LINKA	069	124	0500	10	0000
LINKC	NOP 0000	070	125	0197	69	0004
EQNO2	LDO LINKD	071	126	0197	24	0550
	STO LINKA	072	127	0044	00	0000
LINKD	FOV M202	073	128	0044	69	0245
START	RCD 1951	074	129	0198	24	0550
	RSB 0008	075	130	0245	34	0248
	FOV CSV12	076	131	0253	70	1951
	FDV CSV17	077	132	0151	81	0008
DEL17	RAU 1959	078	133	0168	34	0061
	NZU	079	134	0186	34	0118
	STU 1977	080	135	0107	60	3959
	LDO XSD05	081	136	0363	44	0017
		082	137	0017	21	1977
		083	138	0030	69	0300

	STO X INCR		137	0303	24	0156	0059
	STO X	DEL15	138	0059	24	1979	0132
DEL15	RQU FPONE		139	0132	61	0036	0091
	STU DELTA		140	0091	21	0296	0449
	RAB 0000	DEL19	141	0449	82	0000	0005
DEL19	RAC 0000		142	0005	88	0000	0211
	LDD 1998		143	0211	69	1998	0201
	STO N		144	0201	24	1978	0081
	STO Y	OEL20	145	0081	24	1980	0183
DEL20	NZB 0001	OEL21	146	0183	42	0236	U037
	AXC 0001	OEL22	147	0236	58	0001	0142
DEL21	RAU DELTA		148	0037	60	0296	0251
	FAD DNEPI		149	0251	32	0203	0229
	STU DELTA	DELO6	150	0229	21	0296	0499
DEL06	RAU DELTA		151	0499	60	0296	0301
	LDD	E00CA	152	0301	69	0104	0150
	STU SIN		153	0104	21	0008	0261
	RAU DELTA		154	0261	60	0296	0351
	LDD	E00C8	155	0351	69	0154	0050
	STU COS		156	0154	21	0058	0311
	FDV SIN		157	0311	34	0008	0108
	STU COT		158	0100	21	0012	0115
	FMP 8003		159	0115	39	8003	0169
	STU COT39		160	0169	21	0174	0077
	RAU FPDNE		161	0077	60	0036	0141
	FDV M		162	0141	34	1977	0127
	STU DELO1		163	0127	21	0182	0035
	RAU M		164	0035	60	1977	0131
	FDV DELTA		165	0131	34	0296	0346
	FDV DELTA		166	0346	34	0296	0396
	FAD DELO1		167	0396	32	0182	0109
	FDV FPTWD		168	0109	34	0250	0600
	FAD COT39		169	0600	32	0174	0401
	FAD FPONE		170	0401	32	0036	0413
	STU DELO2		171	0413	21	0268	0271
	RAU DELTA		172	0271	60	0296	0451
	FDV M		173	0451	34	1977	0177
	STU DELO3		174	0177	21	0232	0085
	RAU M		175	0085	60	1977	0181
	FDV DELTA		176	0181	34	0296	0446
	F3B DELO3		177	0446	33	0232	0159
	FDV FPTWD		178	0159	34	0250	0650
	FAD COT		179	0650	32	0012	0239
	FDV DELO2		180	0239	34	0268	0318
	FAD DELTA		181	0318	32	0296	0123
	STU DELO4		182	0123	21	0028	0231
	FDV DELTA		183	0231	34	0296	0496
	F3B FPONE		184	0496	33	0036	0463
	RAM 8003		185	0463	67	8003	0321
	RAU 8002		186	0321	60	8002	0279
	F3B SIZEA		187	0279	33	0450	0227
	BMI DELO5		188	0227	46	0080	0281
	LDD DELO4		189	0281	69	0028	0331
	STO DELTA	DEL06	190	0331	24	0296	0499
DEL05	AXC 0001		191	0080	58	0001	0286
	LDD DELO4		192	0286	69	0028	0381
	STO DELTA		193	0381	24	0296	0549
	RAU DELTA		194	0549	60	0296	0501
	FMP 8003		195	0501	39	8003	0055
	STU DELO7		196	0055	21	0210	0513
	RAO M		197	0513	60	1977	0431
	FMP 8003		198	0431	39	8003	0135
	DEL DELO7		199	0135	34	0210	0087
	STU TEMPX		200	0087	21	0192	0295
	FDV M		201	0295	34	1977	0277
	FDV EPTWD		202	0277	34	0250	0700
	STU 1200 C		203	0700	21	7200	0353
	RAU M		204	0353	60	1977	0481
	FMP FPTWD		205	0481	39	0250	0750
	FAD TEMPX		206	0750	32	0192	0219
	STU DELO10		207	0219	21	0224	0327
	RAU COS		208	0327	60	0058	0563
	FMP DELTA		209	0563	39	0296	0546
	STU DELO11		210	0546	21	0800	0403
	RAO SIN		211	0403	60	0008	0613
	FMP M		212	0613	39	1977	0377
	FAD DELO11		213	0377	32	0800	0427
	FMP DELTA		214	0427	39	0296	0596
	FMP FPTWD		215	0596	60	0250	0850
	FDV DELO10		216	0850	34	0224	0274
DEL22	STO 1000 C	DEL22	217	0274	21	7000	0142
	LDD 1200 C		218	0142	69	7200	0453
	STD M202		219	0453	24	0248	0551
	LDD 1000 C		220	0551	69	7000	0503
	STO NEXTM		221	0503	24	0206	0209
	RAU N		222	0209	60	1978	0233
	FAD FPDNE		223	0233	32	0036	0663
	STU N		224	0663	21	1978	0531
	RAU M202		225	0531	60	0248	0553
	FMP X		226	0553	39	1979	0329
	STU DELO9		227	0329	21	0034	0137
	RAU M		228	0137	60	1977	0581
	F3B DELO9		229	0581	33	0034	0361
	LDD	EXP01	230	0361	69	0214	0000
	FMP NEXTM	LINKA	231	0214	39	0206	0510
LINKA	STU DEL12		232	0148	21	0052	0105
	FAD Y		233	0105	32	1980	0157
	STU Y		234	0157	21	1980	0283
	RAO OEL12		235	0283	60	0052	0207
	FDV Y		236	0207	34	1980	0130
	RAM 8003		237	0130	67	8003	0187
	RAU 8002		238	0187	60	8002	0345
	F3B SIZEA		239	0345	33	0500	0377
DEL13	BMI OEL13	OEL20	240	0477	46	0180	U183
	PCH 1977		241	0180	71	1977	0527
	RAU X		242	0527	60	1979	U333
	FAD X INCR		243	0333	32	0156	0383
	STU X		244	0383	21	1979	0282
	AXB 0001		245	0282	52	0001	0038
	F3B FP38		246	0038	33	0350	0377
	NZU	DEL14	247	0377	44	0631	0332
	FAD FP38		248	0631	32	0350	0627
	F3B FP01		249	0627	33	0400	0677
	NZU DEL19	DEL16	250	0677	44	0005	0382
DEL16	RAO X INCR		251	0382	60	0156	0411
	FMP FPTWD		252	0411	39	0250	0900
	STO X INCR	DEL19	253	0900	21	0156	0005
DEL14	PCH 1877		254	0332	71	1877	U727
	PCH 1877		255	0727	71	1877	U777
	AXA 0001		256	0777	50	0001	0433
	HZA DEL17	START	257	0433	40	0107	0253
DEL18	HLT 9999	9999	258	0218	01	9999	9999

Table 6-a. Block diagram for use with IBM 650 for computing equations (27) and (28).

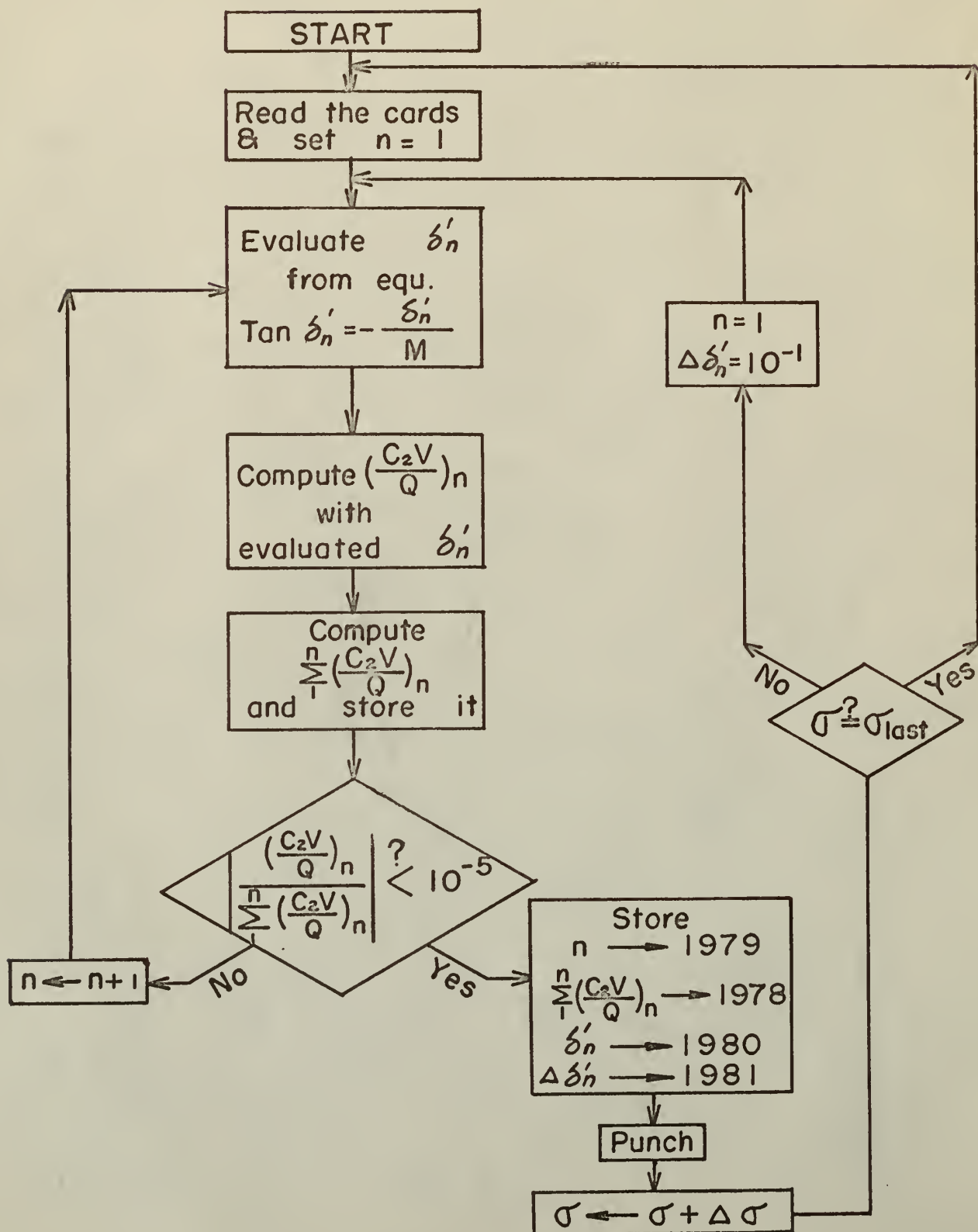


Table 6-b. Block diagram for use with IBM 650 for computing equation (45).

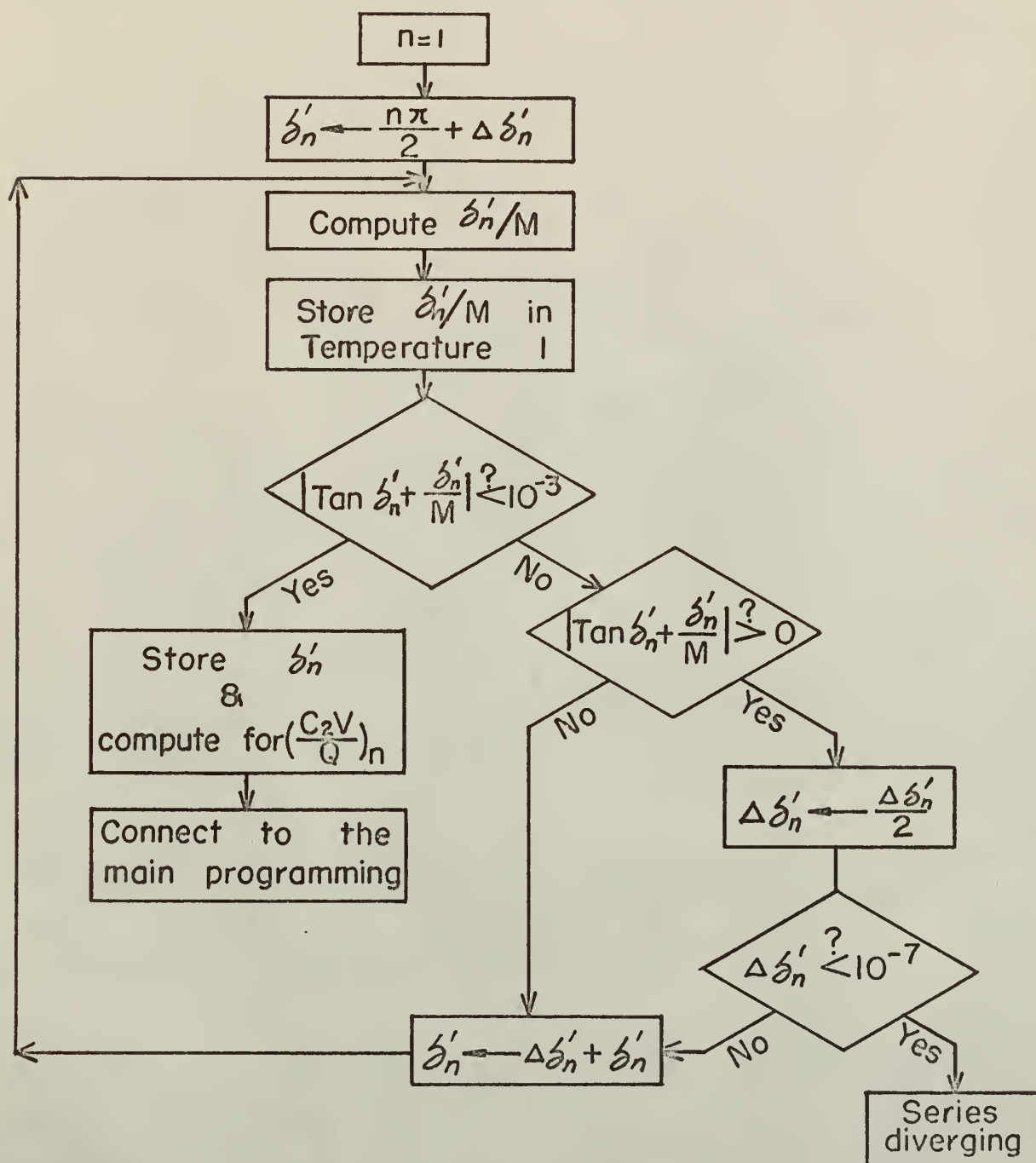


Table 6 - c. Computer program for use with
IBM 650 for computing equations (27) & (45)

	BLR	1000	1999	1	0000	00	0000	0000
PRES1	10	0000	0048	2	0000	10	0000	0048
INCOT	10	0000	0050	3	0050	10	0000	0050
DECOT	10	0000	0050	4	0100	10	0000	0050
HAFPI	15	7079	6351	5	0150	15	7079	6351
FP043	10	0000	0043	6	0200	10	0000	0043
ACUCY	10	0000	0044	7	0250	10	0000	0044
MXTAU	36	0000	0051	8	0350	36	0000	0051
D TAU	50	0000	0049	9	0400	50	0000	0049
DIVGE	99	9999	9999	10	0450	99	9999	9999
ONE 5	10	5000	0051	11	0500	10	5000	0051
FP 3	30	0000	0051	12	0550	30	0000	0051
START	RCO	1951		13	0600	70	1951	0001
SET	RAU	1STAU	SET	14	0001	60	0300	0005
	STU	1977		15	0005	21	1977	0030
	LDO	DECOT		16	0030	69	0100	0003
	STD	INCOT		17	0003	24	0050	0053
	LDO	1951		18	0053	69	1951	0004
	STO	1984		19	0004	24	1984	0037
	LDO	FPONE		20	0037	69	0040	0043
	STD	1979		21	0043	24	1979	0032
AK100	RAU	1979	AK100	22	0032	60	1979	0033
	FMP	DNEPI		23	0033	39	0036	0086
	STU	N PI		24	0086	21	0030	0093
	RAU	1979		25	0093	60	1979	0083
	FMP	FPTWO		26	0083	39	0136	0186
	FSB	FPONE		27	0186	33	0040	0017
	FMP	HAFPI		28	0017	39	0150	0650
	FAO	INCOT		29	0650	32	0050	0027
	FSB	N PI		30	0027	33	0090	0067
	NZU		RETUN	31	0067	44	0021	0022
	BMI		RETUN	32	0021	46	0024	0022
	FAO	N PI		33	0024	32	0090	0117
COMPU	STU	LOCTO	COMPU	34	0117	21	0072	0025
	FDV	1951		35	0025	34	1951	0051
	STU	TEMP1		36	0051	21	0006	0009
	RAU	LOCTO		37	0009	60	0072	0077
	LDO		EOOCB	38	0077	69	0080	0133
	STU	COSIN		39	0080	21	0034	0087
	STU	LOCTO		40	0087	60	0072	0127
	LDO		EOOCA	41	0127	69	0130	0183
	FOV	COSIN		42	0130	34	0034	0084
	FAO	TEMP1		43	0084	32	0006	0233
ASKNG	BMI	ASKNG	CHECK	44	0233	46	0236	0137
	FAD	PRES1		45	0236	32	0000	0177
	NZU	NOSIR	OK	46	0177	44	0031	0082
NOSIR	BMI	MODFY	OK	47	0031	46	0134	0082
MODFY	RAU	LOCTO		48	0134	60	0072	0277
	FAD	INCOT		49	0277	32	0050	0277
	STU	LOCTO	COMPU	50	0277	21	0072	0025
CHECK	FSB	PRES1		51	0137	33	0000	0327
	NZU	TESTO	OK	52	0327	44	0081	0082
TESTO	BMI	OK	RETUN	53	0081	46	0082	0022
RETUN	RAU	INCOT		54	0022	60	0050	0055
	FOV	FPTWO		55	0055	34	0136	0286
	FSB	FP043		56	0286	33	0200	0377
	BMI	MINUS		57	0377	46	0180	0131
	FAO	FP043		58	0131	32	0200	0427
	STU	INCOT		59	0427	21	0050	0103
	RAU	LOCTO		60	0103	60	0072	0477
	FSB	INCOT		61	0477	33	0050	0527
OK	STU	LOCTO	COMPU	62	0527	21	0072	0025
	RAU	LOCTO		63	0082	60	0072	0577
	FMP	LOCTO		64	0577	39	0072	0122
	STU	SQR D		65	0122	21	0026	0029
	RAU	1951		66	0029	60	1951	0105
	FMP	1951		67	0105	39	1951	0101
	STU	MSORE		68	0101	21	0056	0059
	FAO	SQR D		69	0059	32	0026	0153
	FOV	1951		70	0153	34	1951	0151
	FOV	FPTWO		71	0151	34	0136	0336
	FMP	1977		72	0336	39	1977	0627
	STU	80X 2		73	0627	21	0132	0035
	RAU	1951		74	0035	60	1951	0155
	FSB	80X 2		75	0155	33	0132	0109
	LDO		EOOLR	76	0109	69	0012	0015
	STU	1982		77	0012	21	1982	0085
	RAU	1951		78	0085	60	1951	0205
	FAO	MSORE		79	0205	32	0056	0283
	FAO	SQR D		80	0283	32	0026	0203
	FMP	COSIN		81	0203	39	0034	0184
	STU	80X1		82	0184	21	0038	0041
	RSU	SQR D		83	0041	61	0026	0181
	FOV	80X1		84	0181	34	0038	0088
	FMP	1982		85	0088	39	1982	0182
	STU	1981		86	0182	21	1981	0234
	FAO	1978		87	0234	32	1978	0255
	STU	1978		88	0255	21	1978	0231
	RAU	1981		89	0231	60	1981	0135
	FOV	1978		90	0135	34	1978	0028
	RAM	8003		91	0028	67	8003	0185
	RAU	8002		92	0185	60	8002	0143
	FSB	ACUCY		93	0143	33	0250	0677
	BMI	THANK		94	0677	46	0230	0281
	RAU	1979		95	0281	60	1979	0333
	FAO	FPONE		96	0333	32	0040	0167
	STU	1979		97	0167	21	1979	0232
	LDO	DECOT		98	0232	69	0100	0253
THANK	STO	INCOT	AK100	99	0253	24	0050	0032
	LDO	LOCTO		100	0230	69	0072	0075
	STO	1980		101	0075	24	1980	0383
	LDO	COSIN		102	0383	69	0034	0187
	STO	1983		103	0187	24	1983	0386
	PCH	1977	AOTAU	104	0386	71	1977	0727
AOTAU	LDO	1999		105	0727	69	1999	0002
	STD	1978		106	0002	24	1978	0331
	RAU	1977		107	0331	60	1977	0381
				108				

	FAD D TAU		109	0381	32	0400	0777
	FSB ONE S		110	0777	33	0500	0827
	BMI	GD1	111	0827	46	0280	0431
	FAD ONE S	SET	112	0280	32	0500	0005
GD1	FAD ONE S		113	0431	32	0500	0877
	FAD D TAU		114	0877	32	0400	0927
	FSB MXTAU		115	0927	33	0350	0977
	BMI	CONN1	116	0977	46	0330	0481
	FAD MXTAU	SET	117	0330	32	0350	0005
CONN1	LDD DECO1		118	0481	69	0100	0303
	STD INCDT	START	119	0303	24	0050	0600
MINUS	LDD DIVGE		120	0180	69	0450	0353
	STD 1983	PUNCH	121	0353	24	1983	0436
PUNCH	PCH 1977		122	0436	71	1977	0078
	LDD 1999		123	0070	69	1999	0052
	STD 1983	AOTAU	124	0052	24	1983	0727
E00C8	STD CSV02	BEGIN CDS	125	0133	24	0486	0039
	NZE CSV03	CSV04	002	126	0039	45	0042
CSV04	RAU FPONE	CSV02	003	127	0193	60	0040
CSV03	BMI CSV05	CSV06	004	128	0042	46	0045
CSV05	FAD TWDPI	CSV04	005	129	0045	32	0048
	NZE		006	130	0125	45	0128
	BMI CSV05		007	131	0128	46	0045
	FSB ONEPI		008	132	0282	33	0036
	NZE CSV07	CSV08	009	133	0013	45	0016
CSV08	RSU FPONE	CSV02	010	134	0217	61	0040
CSV06	FSB TWDPI		011	135	0046	33	0048
	NZE	CSV04	012	136	0175	45	0178
	BMI	CSV06	013	137	0178	46	0531
	FAD ONEPI		014	138	0531	32	0036
	NZE CSV07	CSV08	015	139	0063	45	0016
CSV07	STU CSV09		016	140	0016	21	0020
	RSU FPONE		017	141	0023	61	0040
	STU CSV10		018	142	0095	21	0700
	STU CSV11		019	143	0403	21	0008
	STD CSV12	CSV13	020	144	0011	20	0065
E00CA	STD CSV02	BEGIN S	021	145	0113	24	0486
	NZE	CSV02	022	146	0089	45	0092
	BMI CSV14	CSV15	023	147	0092	46	0145
CSV14	FAD TWDPI		024	148	0145	32	0048
	NZE	CSV02	025	149	0225	45	0228
	BMI CSV14		026	150	0228	46	0145
	FSB ONEPI		027	151	0332	33	0036
	NZE CSV16	CSV02	028	152	0113	45	0066
CSV15	FSB TWDPI		029	153	0096	32	0048
	NZE	CSV02	030	154	0278	46	0486
	BMI	CSV15	031	155	0278	46	0581
	FAD ONEPI		032	156	0581	32	0036
	NZE CSV16	CSV02	033	157	0163	45	0066
CSV16	STU CSV09		034	158	0066	21	0020
	RSU 8003		035	159	0073	61	8003
	STU CSV10		036	160	0631	21	0700
	STU CSV11		037	161	0453	21	0008
	LDD		038	162	0061	69	0040
	STD CSV12	CSV13	039	163	0243	24	0065
CSV13	RAU CSV12		040	164	0018	60	0065
	FAD FPONE		041	165	0019	32	0040
	STU CSV17		042	166	0267	21	0172
	FAD FPONE		043	167	0325	32	0040
	STU CSV12		044	168	0317	21	0065
	RSU CSV10		045	169	0068	61	0700
	FMP CSV09		046	170	0305	39	0020
	FMP CSV09		047	171	0070	39	0020
	FDV CSV17		048	172	0120	34	0172
	FDV CSV12		049	173	0222	34	0065
	STU CSV10		050	174	0115	21	0700
	FDV CSV11		051	175	0503	34	0008
	RAM 8003		052	176	0058	67	8003
	RAU 8002		053	177	0165	60	8002
	FSB 81258		054	178	0323	33	0076
	BMI CSV18		055	179	0553	46	0106
	RAU CSV11		056	180	0007	60	0008
	FAD CSV10		057	181	0213	32	0700
	STU CSV11	CSV13	058	182	0328	21	0008
CSV18	RAU CSV11	CSV02	059	183	0106	60	0008
FPONE	10 0000	0051	060	184	0040	10	0000
FTWD	20 0000	0051	061	185	0136	20	0000
TWOP	62 83351	5351	062	186	0286	62	83351
SIZE8	10 0000	0043	063	187	0076	10	0000
ONEPI	31 4159	2751	064	188	0036	31	4159
E00LR	STO EXP02		065	189	0015	24	0118
	STU EXP03		066	190	0071	21	0126
	LDD 8007		067	191	0079	69	8007
	STD EXP04		068	192	0235	24	0138
	RAC 0000		069	193	0091	88	0000
	RAU EXP03		070	194	0047	60	0126
	BMI EXP05	EXP06	071	195	0681	46	0284
EXP06	FSB LNTEN		072	196	0285	33	0188
	BMI EXP07		073	197	0215	46	0168
	AXC 0001	EXP06	074	198	0069	58	0001
EXP07	FAD LNTEN	EXP08	075	199	0168	32	0188
EXP05	SXC 0001		076	200	0284	59	0001
	FAD LNTEN		077	201	0140	32	0188
	BMI EXP05	EXP08	078	202	0315	46	0284
EXP08	STU EXP03		079	203	0265	21	0126
	STU EXP09		080	204	0129	21	0334
	RAL EXP04		081	205	0237	65	0138
	LDD 8007		082	206	0293	69	8007
	STD EXP04		083	207	0049	24	0138
	RAC 8002		084	208	0191	88	8002
	LDD FPONE		085	209	0099	69	0040
	STD EXP10		086	210	0343	24	0146
	STO EXP11	EXP12	087	211	0149	24	0102
EXP12	RAU EXP10		088	212	0355	60	0146
	FAD EXP09		089	213	0201	32	0334
	STU EXP10		090	214	0111	21	0146
	RAU EXP11		091	215	0199	60	0102
	EAD FPONE		092	216	0057	32	0040
	STU EXP11		093	217	0367	69	0102
	RAU EXP09		094	218	0405	60	0334
	FMP EXP03		095	219	0139	39	0126
	FDV EXP11		096	220	0176	34	0102
	STU EXP09		097	221	0152	21	0334
	EDV EXP10		098	222	0287	34	0146
	FSB SIZE8		099	223	0196	33	0076
	BMI	EXP12	100	224	0355	46	0156
	RAU EXP10		101	225	0156	60	0146
	AUP EXP04	EXP02	102	226	0251	10	0138
FPONE	10 0000	0051	103	227	0040	10	0000
SIZE8	10 0000	0043	104	228	0076	10	0000
LNTEN	23 0258	5151	105	229	0188	23	0258

Table 6-d. Computer program for use with
IBM 650 for computing equations (28) & (45).

	BLR	1000	1999	1	0000	00	0000	0000
PRES1	10	0000	0048	2	0000	10	0000	0048
INCDT	10	0000	0050	3	0050	10	0000	0050
DECO1	10	0000	0050	4	0100	10	0000	0050
HAFPI	15	7079	6351	5	0150	15	7079	6351
FP043	10	0000	0043	6	0200	10	0000	0043
ACUCY	10	0000	0044	7	0250	10	0000	0044
1STAU	50	0000	0049	8	0300	50	0000	0049
MXTAU	36	0000	0051	9	0350	36	0000	0051
D TAU	50	0000	0049	10	0400	50	0000	0049
DIVGE	99	9999	9999	11	0450	99	9999	9999
ONE 5	10	5000	0051	12	0500	10	5000	0051
FP 3	30	0000	0051	13	0550	30	0000	0051
START	RCD	1951		14	0600	70	1951	0001
	RAU	1STAU	SET	15	0001	60	0300	0005
SET	STU	1977		16	0005	21	1977	0030
	LOD	DECO1		17	0030	69	0100	0003
	STD	INCDT		18	0003	24	0050	0053
	LOO	1951		19	0053	69	1951	0004
	STO	1984		20	0004	24	1984	0037
	LOD	FPONE		21	0037	69	0040	0043
	STD	1979	AK100	22	0043	24	1979	0032
AK100	RAU	1979		23	0032	60	1979	0033
	FMP	ONEPI		24	0033	39	0036	0086
	STU	N PI		25	0086	21	0090	0093
	RAU	1979		26	0093	60	1979	0083
	FMP	FPTWO		27	0083	39	0136	0186
	FS8	FPONE		28	0186	33	0040	0017
	FMP	HAFPI		29	0017	39	0150	0650
	FAO	INCDT		30	0650	32	0050	0027
	FS8	N PI		31	0027	33	0090	0067
	NZU		RETUN	32	0067	44	0021	0022
	8MI		RETUN	33	0021	46	0024	0022
	FAD	N PI		34	0024	32	0090	0117
COMPU	STU	LOCTD	COMPU	35	0117	21	0072	0025
	FDV	1951		36	0025	34	1951	0051
	STU	TEMP1		37	0051	21	0006	0009
	RAU	LOCTD		38	0009	60	0072	0077
	LDO		EOOCB	39	0077	69	0080	0133
	STU	COSIN		40	0080	21	0034	0087
	RAU	LOCTD		41	0087	60	0072	0127
	LDD		EOOCA	42	0127	69	0130	0183
				43	0130	21	0039	0137
	FDV	COSIN		44	0137	34	0034	0134
	FAD	TEMP1		45	0134	32	0006	0233
ASKNG	8MI	ASKNG	CHECK	46	0233	46	0236	0187
	FAD	PRES1		47	0236	32	0000	0177
	NZU	NOSIR	OK	48	0177	44	0031	0082
NOSIR	8MI	MODFY	OK	49	0031	46	0184	0082
MOOFY	RAU	LOCTD		50	0184	60	0072	0227
	FAD	INCDT		51	0227	32	0050	0277
	STU	LOCTD	COMPU	52	0277	21	0072	0025
CHECK	FS8	PRES1		53	0187	33	0000	0327
	NZU	TESTQ	OK	54	0327	44	0081	0082
TESTQ	8MI	OK	RETUN	55	0081	46	0082	0022
RETUN	RAU	INCDT		56	0022	60	0050	0055
	FDV	FPTWO		57	0055	34	0136	0286
	FS8	FP043		58	0286	33	0200	0377
	8MI	MINUS		59	0377	46	0180	0131
	FAD	FP043		60	0131	32	0200	0427
	STU	INCDT		61	0427	21	0050	0103
	RAU	LOCTD		62	0103	60	0072	0477
	FS8	INCDT		63	0477	33	0050	0527
	STU	LOCTD	COMPU	64	0527	21	0072	0025
OK	RAU	LOCTD		65	0025	60	0072	0577
	FMP	LOCTD		66	0577	39	0072	0122
	STU	SORD		67	0122	21	0026	0029
	RAU	1951		68	0029	60	1951	0105
				69	0105	39	1951	0155
	STU	MSQRE		70	0155	21	0010	0013
	FAO	SORD		71	0013	32	0026	0153
	FDV	1951		72	0153	34	1951	0101
	FDV	FPTWO		73	0101	34	0136	0336
	FMP	1977		74	0336	39	1977	0627
	STU	BOX 2		75	0627	21	0132	0035
	RAU	1951		76	0035	60	1951	0205
	FS8	BOX 2		77	0205	33	0132	0059
	LDD		EOOLR	78	0059	69	0012	0015
	STU	1982		79	0012	21	1982	0085
	RAU	1951		80	0085	60	1951	0255
	FAD	MSQRE		81	0255	32	0010	0237
	FAD	SORD		82	0237	32	0026	0003
	STU	BOX 3		83	0203	21	0008	0011
	RAU	LOCTD		84	0011	60	0072	0677
	FMP	FPTWO		85	0677	39	0136	0386
	FMP	SIGN		86	0386	39	0039	0089
	FDV	BOX 3		87	0089	34	0008	0058
	FMP	1982		88	0058	39	1982	0182
	STU	1981		89	0182	21	1981	0234
	FAO	SUM		90	0234	32	0287	0063
	STU	SUM		91	0063	21	0287	0140
	RAU	1981		92	0140	60	1981	0135
	FDV	SUM		93	0135	34	0287	0337
	RAM	8003		94	0337	67	8003	0045
	RAU	8002		95	0045	60	8002	0253
	FS8	ACUCY		96	0253	33	0250	0727
	8MI	THANK		97	0727	46	0230	0181
	RAU	1979		98	0181	60	1979	0283
	FAD	FPONE		99	0283	32	0040	0167
	STU	1979		100	0167	21	1979	0232
	LDO	DECO1		101	0232	69	0100	0303
	STD	INCDT	AK100	102	0303	24	0050	0032
THANK	LDD	LOCTD		103	0230	69	0072	0075
	STD	1980		104	0075	24	1980	0333

	LOO	SUGN		105	03333	69	0084	0387
	STD	1983		106	0387	24	1983	0436
	RAU	FPONE		107	0436	60	0040	0095
	FSB	SUM		108	0095	33	0287	0113
	STU	1978		109	0113	21	1978	0231
	PCH	1977	AOTAU	110	0231	71	1977	0777
AOTAU	LOO	1999		111	0777	69	1999	0002
	STD	SUM		112	0002	24	0287	0190
	RAU	1977		113	0190	60	1977	0281
	FAD	O TAU		114	0281	32	0400	0827
	FSB	ONE 5		115	0827	33	0500	0877
	BMI		G01	116	0877	46	0280	0331
	FAD	ONE 5	SET	117	0280	32	0500	0005
G01	FAD	ONE 5		118	0331	32	0500	0927
	FAO	O TAU		119	0927	32	0400	0977
	FSB	MXTAU		120	0977	33	0350	0028
	BMI	MXTAU	CONN1	121	0028	46	0381	0282
	LOO	DEC01	SET	122	0381	32	0350	0805
CONN1	STD	INCOT	START	123	0282	69	0100	0353
MINUS	LOD	OIVGE		124	0353	24	0050	0600
	STD	1983	PUNCH	125	0180	69	0450	0403
PUNCH	PCH	1977		126	0403	24	1983	0486
	LOO	1999		127	0486	71	1977	0078
	STD	1983	AOTAU	128	0078	69	1999	0052
E00C8	STU	CSV02	BEGIN COS	129	0052	24	1983	0777
	NZE	CSV03		130	0133	24	0536	0139
CSV04	RAU	FPONE	CSV04	131	0139	45	0042	0143
CSV03	BMI	CSV05	CSV06	003	132	60	0040	0536
CSV05	FAD	TWOPI		004	133	46	0145	0046
	NZE		CSV04	005	134	32	0048	0125
	BMI	CSV05		006	135	45	0128	0143
	FSB	ONEPI		007	136	46	0145	0332
	NZE	ONEPI		008	137	33	0036	0163
	RSU	FPONE	CSV08	009	138	45	0010	0217
CSV08	RSU	FPONE	CSV02	010	139	61	0040	0536
CSV06	FSB	TWOPI		011	140	33	0048	0175
	NZE		CSV04	012	141	45	0178	0143
	BMI		CSV06	013	142	46	0431	0046
	FAO	ONEPI		014	143	32	0036	0213
	NZE	CSV07	CSV0H	015	144	45	0016	0217
CSV07	STU	CSV09		016	145	21	0020	0023
	RSU	FPONE		017	146	61	0040	0195
	STU	CSV10		018	147	21	0750	0453
	STU	CSV11		019	148	45	0108	0061
	STL	CSV12	CSV13	020	149	20	0065	0018
E00CA	STD	CSV02	BEGIN SIN	150	0183	24	0536	0189
	NZE			021	151	45	0092	0536
	BMI	CSV14	CSV02	022	152	46	0245	0096
CSV14	FAD	TWOPI	CSV15	023	153	32	0048	0225
	NZE		CSV02	024	154	45	0228	0536
	BMI	CSV14		025	155	46	0245	0382
	FSB	ONEPI		026	156	33	0036	0263
	NZE	CSV16	CSV02	027	157	45	0066	0536
CSV15	FSB	TWOPI		028	158	33	0048	0275
	NZE		CSV02	029	159	45	0278	0536
	BMI		CSV15	030	160	46	0481	0096
	FAO	ONEPI		031	161	32	0036	0313
	NZE	CSV16	CSV02	032	162	45	0066	0536
CSV16	STU	CSV09		033	163	21	0020	0073
	RSU	8003		034	164	61	8003	0531
	STU	CSV10		035	165	21	0750	0503
	STU	CSV11		036	166	21	0108	0111
	LOO	FPONE		037	167	69	0040	0193
	STD	CSV12	CSV13	038	168	24	0065	0018
CSV13	RAU	CSV12		039	169	60	0065	0019
	FAD	FPONE		040	170	32	0040	0267
	STU	CSV17		041	171	21	0172	0325
	FAD	FPONE		042	172	32	0040	0317
	STU	CSV12		043	173	21	0065	0068
	RSU	CSV10		044	174	61	0750	0305
	FMP	CSV09		045	175	39	0020	0070
	FMP	CSV09		046	176	39	0020	0120
	FOV	CSV17		047	177	34	0172	0222
	FOV	CSV12		048	178	34	0065	0115
	STU	CSV10		049	179	21	0750	0533
	FOV	CSV11		050	180	34	0108	0158
	RAM	8003		051	181	67	8003	0165
	RAU	8002		052	182	60	8002	0123
	FSB	SIZE8		053	183	33	0076	0603
	BMI	CSV18		054	184	46	0056	0007
	RAU	CSV11		055	185	60	0108	0363
	FAD	CSV10		056	186	32	0750	0328
	STU	CSV11	CSV13	057	187	21	0108	0018
CSV18	RAU	CSV11	CSV02	058	188	60	0108	0536
FPONE	10	0000	0051	059	189	10	0000	0051
FPTWO	20	0000	0051	060	190	20	0000	0051
TWOPI	62	8318	5351	061	191	62	8318	5351
SIZE8	10	0000	0043	062	192	10	0000	0043
ONEPI	31	4159	2751	063	193	31	4159	2751
E00LR	STD	EXP02		064	194	21	0110	0071
	STU	EXP03		065	195	21	0126	0079
	LOO	8U07		066	196	69	8007	0185
	STD	EXP04		067	197	24	0038	0041
	RAC	0000		068	198	88	0000	0047
	RAU	EXP03		069	199	60	0126	0581
	BMI	EXP05	EXP06	070	200	46	0284	0235
	FSB	LNTEN		071	201	33	0088	0215
	BMI	EXP07		072	202	46	0168	0069
	AXC	0001	EXP06	073	203	58	0001	0235
EXP06	FAO	LNTEN	EXP08	074	204	32	0088	0265
	SXC	0001		075	205	59	0001	0240
EXP07	FAO	LNTEN		076	206	32	0000	0315
EXP05	BMI	EXP05	EXP08	077	207	46	0284	0265
	STU	EXP03		078	208	21	0126	0189
	STU	EXP09		079	209	21	0314	0437
	RAL	EXP04		080	210	65	0038	0843
	LOO	8007		081	211	69	8007	0049
	STD	EXP04		082	212	34	0038	0091
	RAC	8002		083	213	88	8002	0099
	LOO	FPONE		084	214	69	0040	0293
	STD	EXP10		085	215	24	0146	0149
	STD	EXP11	EXP12	086	216	24	0108	0355
EXP12	RAU	EXP10		087	217	60	0146	0181
	FAD	EXP09		088	218	32	0334	0161
	STU	EXP10		089	219	21	0146	0199
	RAU	EXP11		090	220	60	0102	0057
	FAO	FPONE		091	221	32	0040	0367
	STU	EXP11		092	222	21	0102	0405
	FMP	EXP03		093	223	60	0334	0239
	FOV	EXP11		094	224	39	0126	0176
	STU	EXP09		095	225	34	0102	0172
	FOV	EXP10		096	226	21	0334	0487
	FSB	SIZE8		097	227	34	0146	0196
	BMI		EXP12	098	228	33	0076	0653
	RAU	EXP10		099	229	46	0106	0355
	AUP	EXP04	EXP04	100	230	60	0146	0201
FPONE	10	0000	0051	101	231	10	0038	0118
SIZE8	10	0000	0043	102	232	10	0000	0051
LNTEN	23	0258	5151	103	233	10	0000	0043
				104	234	23	0258	5151

RESULTS AND DISCUSSION

Steady-state Tubular Flow Reactor

The concentration profiles for the various combinations of M and R for the first-order chemical reactions were calculated by use of equations (9) through (12), (16) and (18). Typical results corresponding to $R=1/2$ for B.C.(I), (II) and (III) are shown in Fig. 1 and those corresponding to $R=2$ for B.C.(IV) are shown in Fig. 2. The figures show that the concentration profiles obtained by use of B.C.(II), (III) and (IV) were inconsistent with the limits of both a plug flow reactor and a uniformly stirred tank reactor for all values of M and R ; whereas the profiles computed by use of B.C.(I) were consistent with such limits.

The word "inconsistent" means that a computed conversion or concentration profile either less than that of a uniformly stirred tank reactor or greater than that of a plug flow reactor is unreasonable. Models which predict this are therefore "inconsistent". The numerical values of the limits of inconsistency were obviously calculated by use of equations (16) and (18) for a uniformly stirred tank reactor and a plug flow reactor respectively.

Figure 3, however, indicates that the over-all conversions computed by use of the different sets of boundary conditions were numerically in agreement for large values of M (small degree of axial dispersion) and deviate increasingly from each other as the degree of axial dispersion increases. Again the overall conversions computed by use of B.C.(II) and (III) gave values lower than those obtained for a uniformly stirred tank reactor, and those

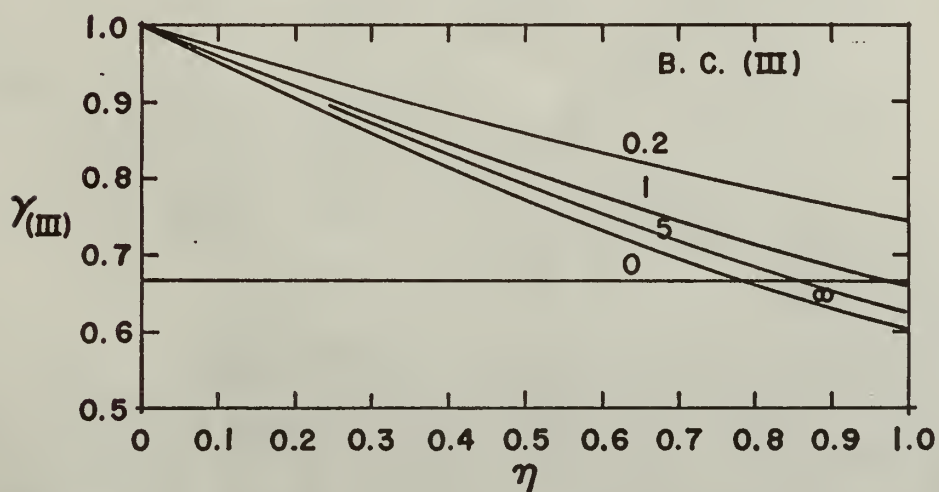
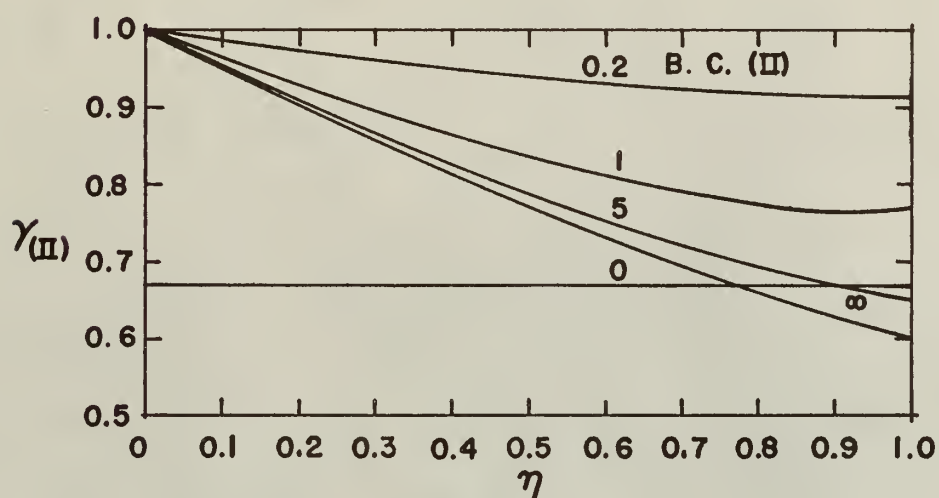
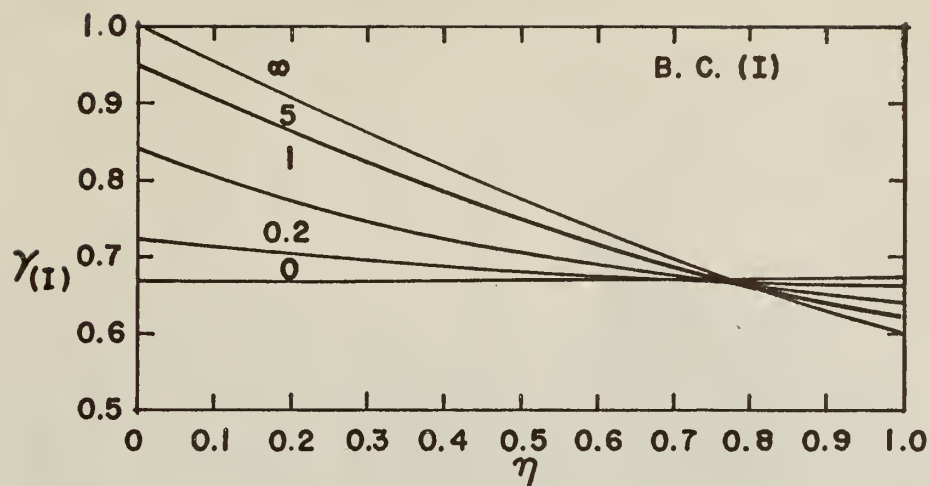


FIG. 1. Concentration profile of reactant for various M for first-order reaction with $R = \frac{1}{2}$.

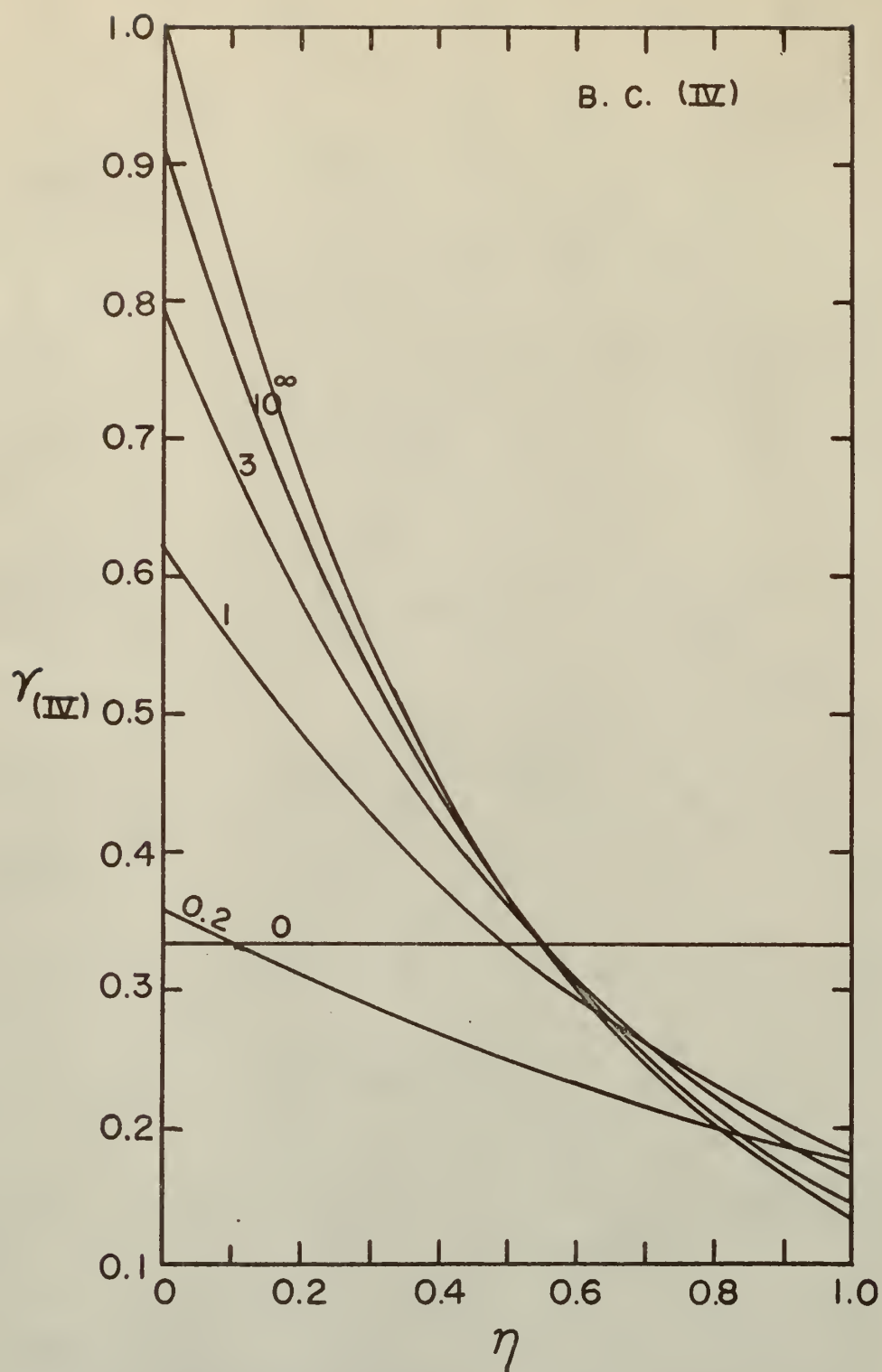


Fig. 2. Concentration profile of reactant for various M for first-order reaction with $R=2$

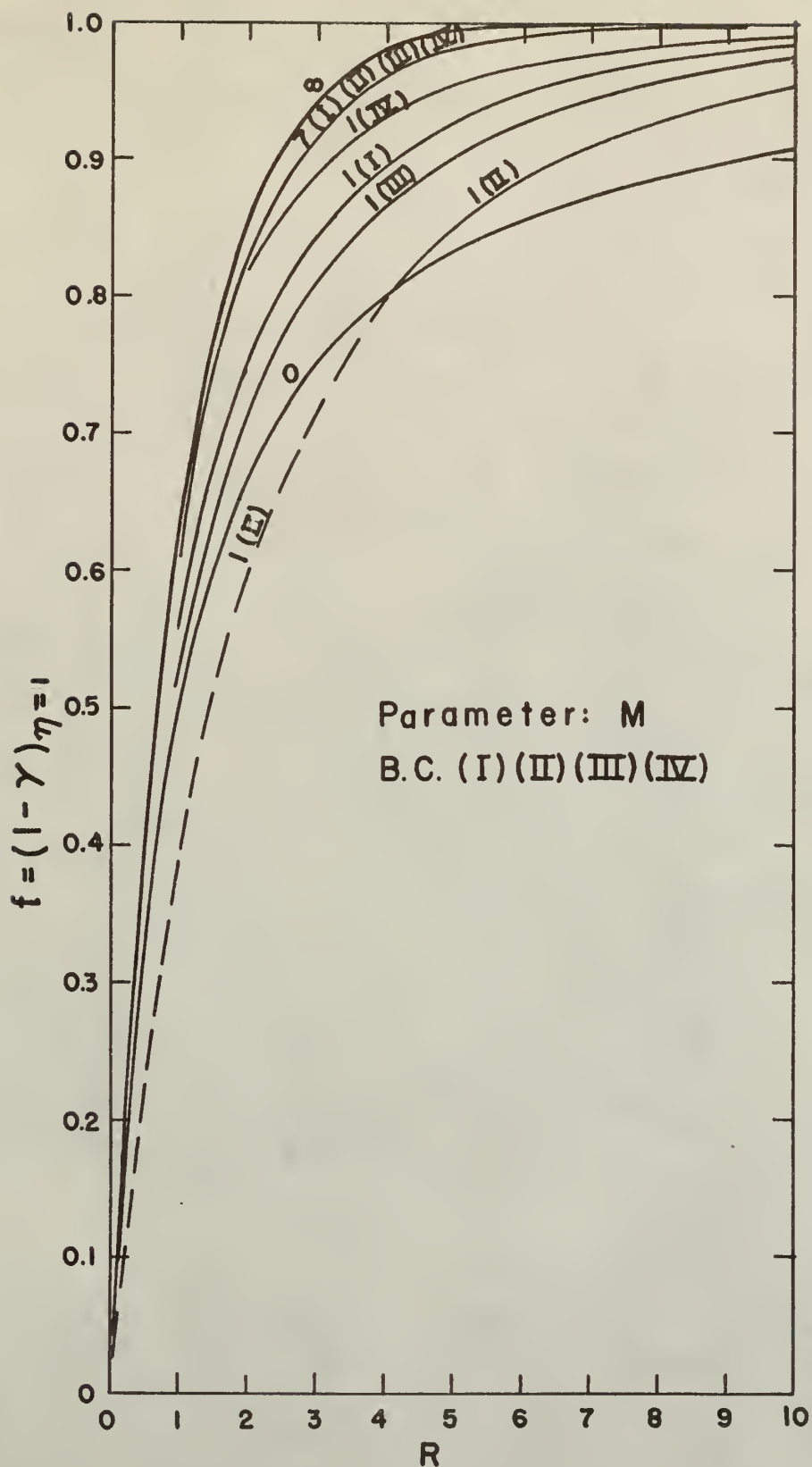


FIG.3. Comparison of final conversions as a function of R for first-order reaction.

computed by use of B.C.(IV) gave values higher than those obtained in a plug flow reactor, especially for low conversions and a high degree of axial dispersion, i. e., low values of R and M. Reactors falling in this category are sometimes met in laboratory and pilot plant experimentation, though not commonly required for design purposes.

The inconsistency shown in Figs. 1 and 2 led to the establishment of the domain for which any combinations of M and R are inconsistent, with the limits of the domain, as shown in Figs. 4 and 5. The curves in Fig. 4 show the locus representing:

$$(f_{II}) \eta_{=1} - f_{M=0} = 0$$

$$(f_{III}) \eta_{=1} - f_{M=0} = 0$$

The domain below the curves, then, corresponds to the condition:

$$(f_{II}) \eta_{=1} - f_{M=0} < 0$$

$$(f_{III}) \eta_{=1} - f_{M=0} < 0$$

The curve in Fig. 5 indicates the locus representing:

$$(f_{IV}) \eta_{=1} - f_{M=\infty} = 0$$

The domain below the curve, then, corresponds to the condition:

$$(f_{IV}) \eta_{=1} - f_{M=\infty} > 0$$

Where f represents the overall conversions.

The inconsistency of the overall conversions using B.C.(II) has been indicated by Wehner and Wilhelm (24). However, it has

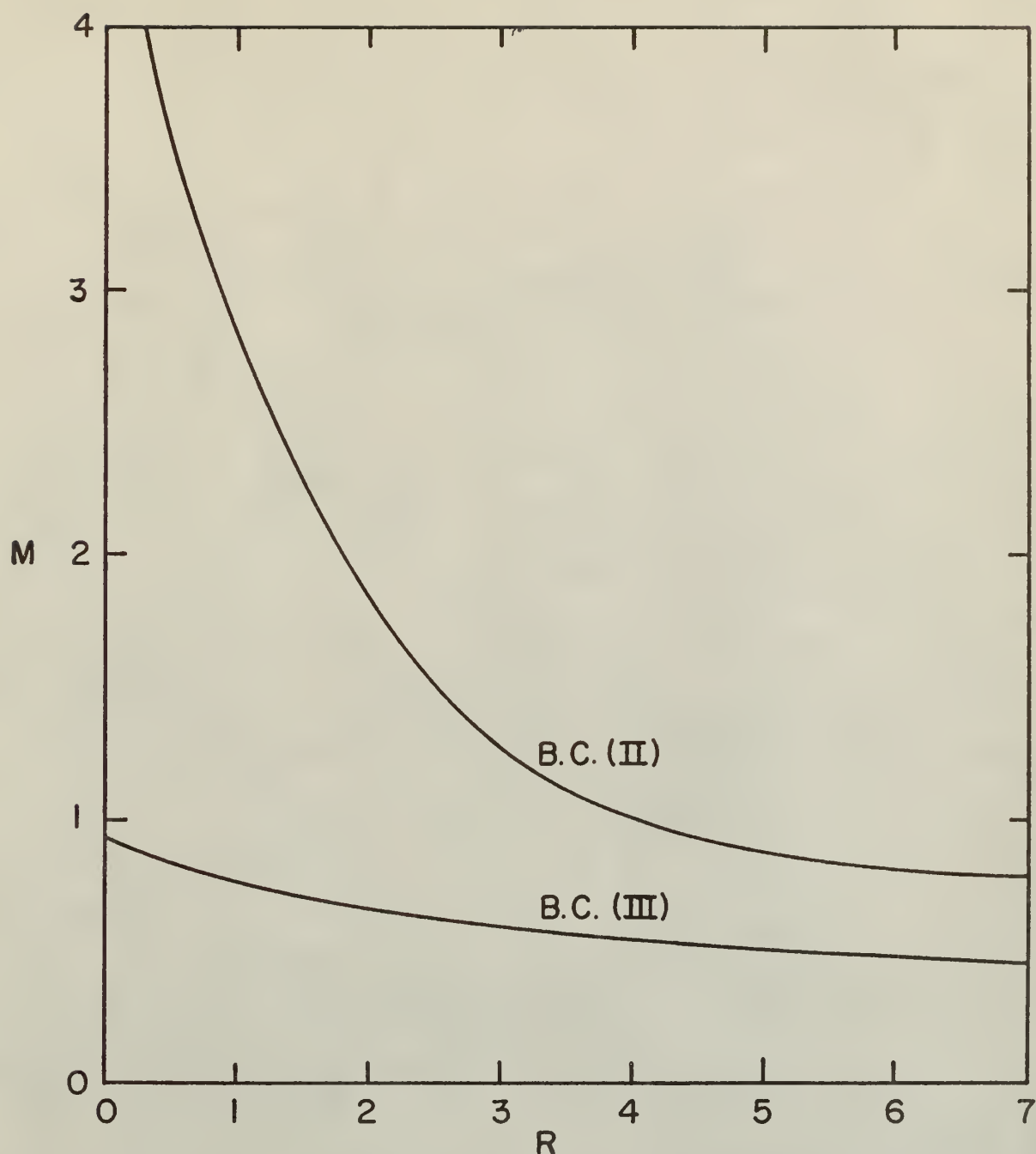


Fig. 4. Locus indicating regions in which conversions are inconsistent with limit of an uniformly stirred tank reactor for first-order reaction.

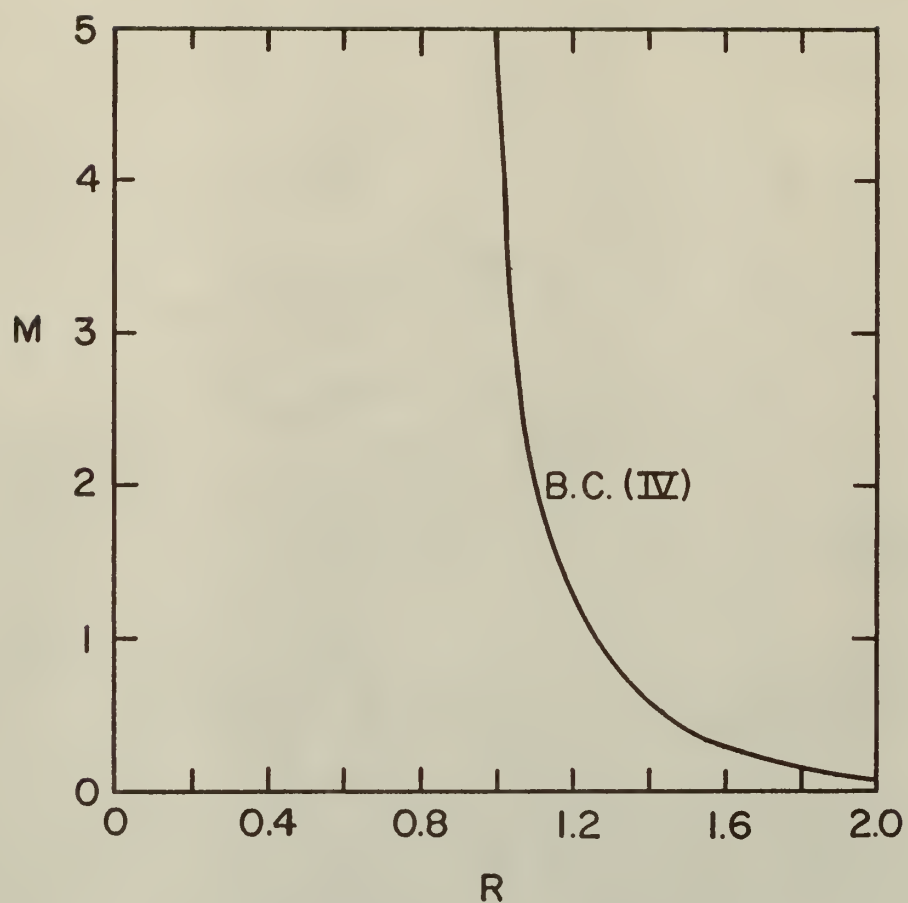


Fig.5. Locus indicating regions in which conversions are inconsistent with limit of a plug flow reactor for first-order reaction.

never been established quantitatively as a function of the system characteristics, i. e., M and R . The cases of B.C.(III) and (IV) have been hitherto unknown. The overall conversions computed by use of B.C.(II) and (III) are consistent with the upper limit of a plug flow reactor, and those by use of B.C.(IV) are consistent with the lower limit of a stirred tank reactor.

The design of an isothermal reactor should not be based only on the overall conversion when the reactions to be carried out in the reactor are exothermic or endothermic. The estimation of the heat transfer area to be provided requires a knowledge of the concentration profile, or the conversion along the length of reactor. Figs. 6 and 7 show the comparison of concentration profiles based on the different sets of boundary conditions at both high and low degree of dispersions. The figures supplement the previous statement that the deviations among the profiles decrease as M increases. This means that the effect of boundary conditions become negligible for high values of M , say of the order of 15, and the reactor can be assumed to behave as a plug flow reactor with reasonable accuracy. The percentage deviations of the overall conversions computed by use of B.C.(II), (III) and (IV) from those computed by use of B.C.(I) are shown in Figs. 8, 9 and 10. The reason that B.C.(I) is used as a basis for the computation is that, as pointed out previously, this yields conversions which are consistent with both limits of the reactor. These figures show that the deviations increase as M and R decrease, and B.C.(II) always yields larger deviations than B.C.(III) and (IV). Since the solutions to B.C.(III) and (IV) are in the form of a simple exponential

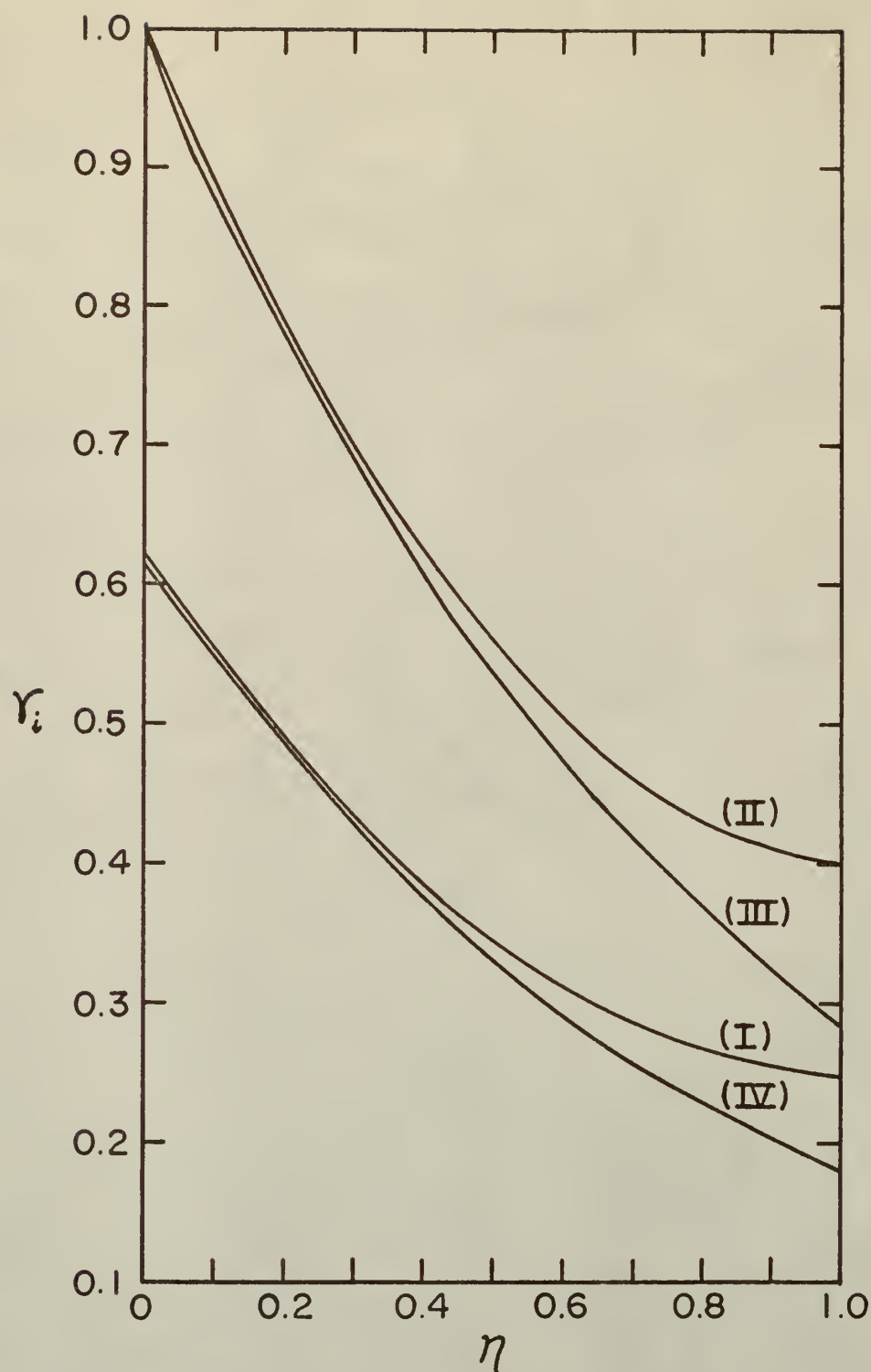


Fig. 6. Comparison of concentration profile of reacting systems ($R=2$) based on different boundary conditions at high degree of dispersion ($M=1$).

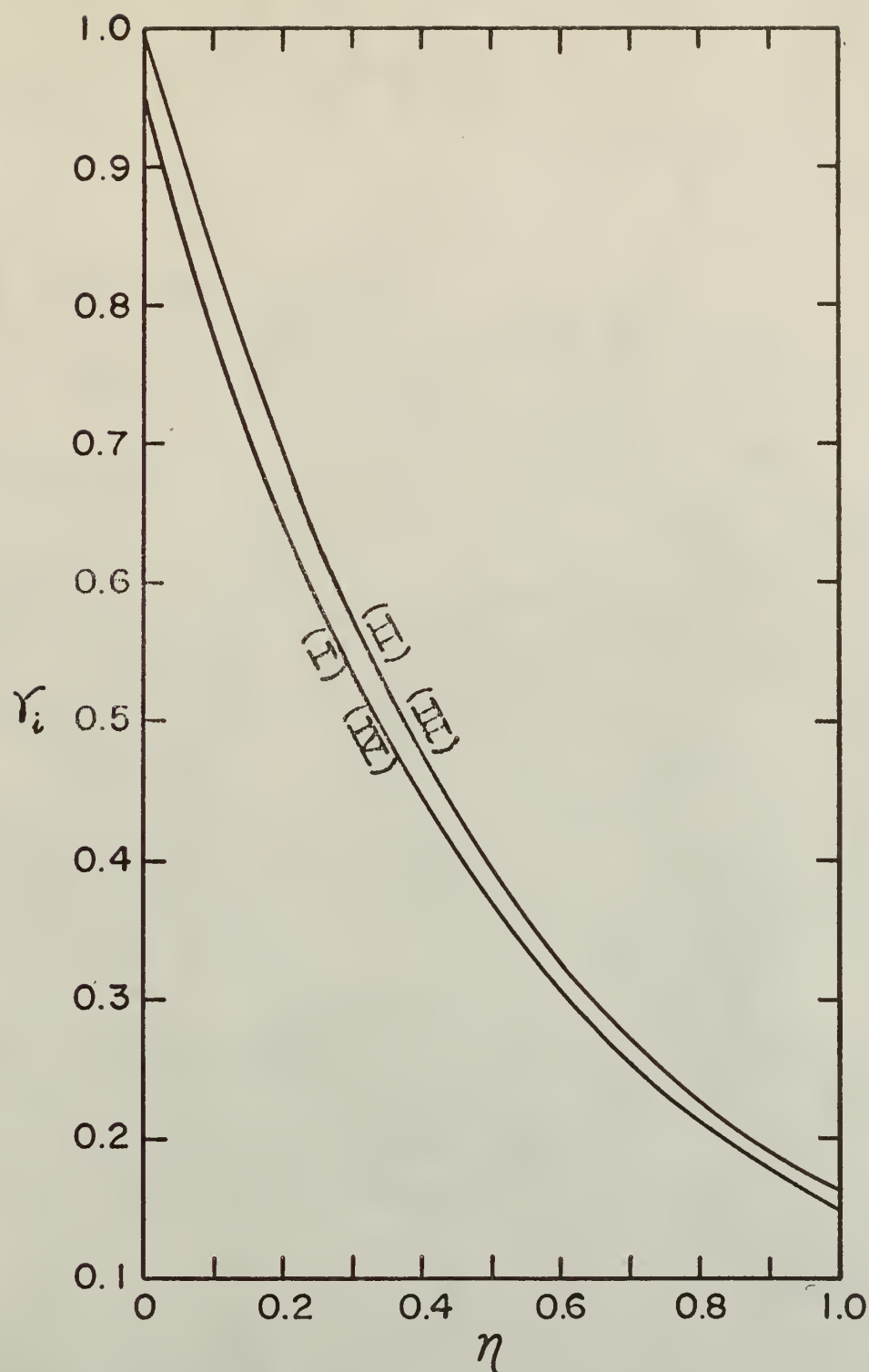


Fig. 7. Comparison of concentration profile of reacting systems ($R=2$) based on different boundary conditions at low degree of dispersion ($M=15$).

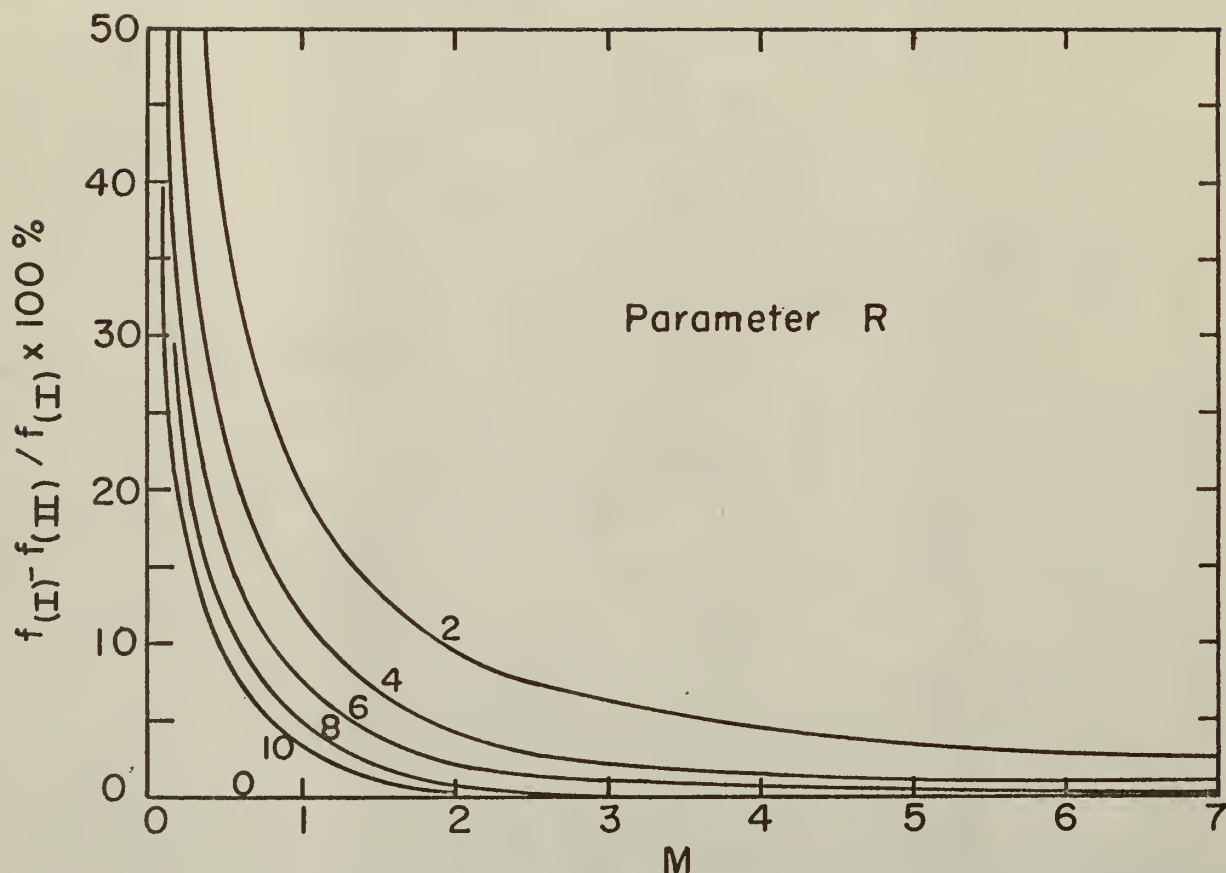


Fig. 8. Percentage deviations of final conversions computed by use of boundary conditions (II) from those computed by use of boundary conditions (I) as functions of M for first-order reactions.

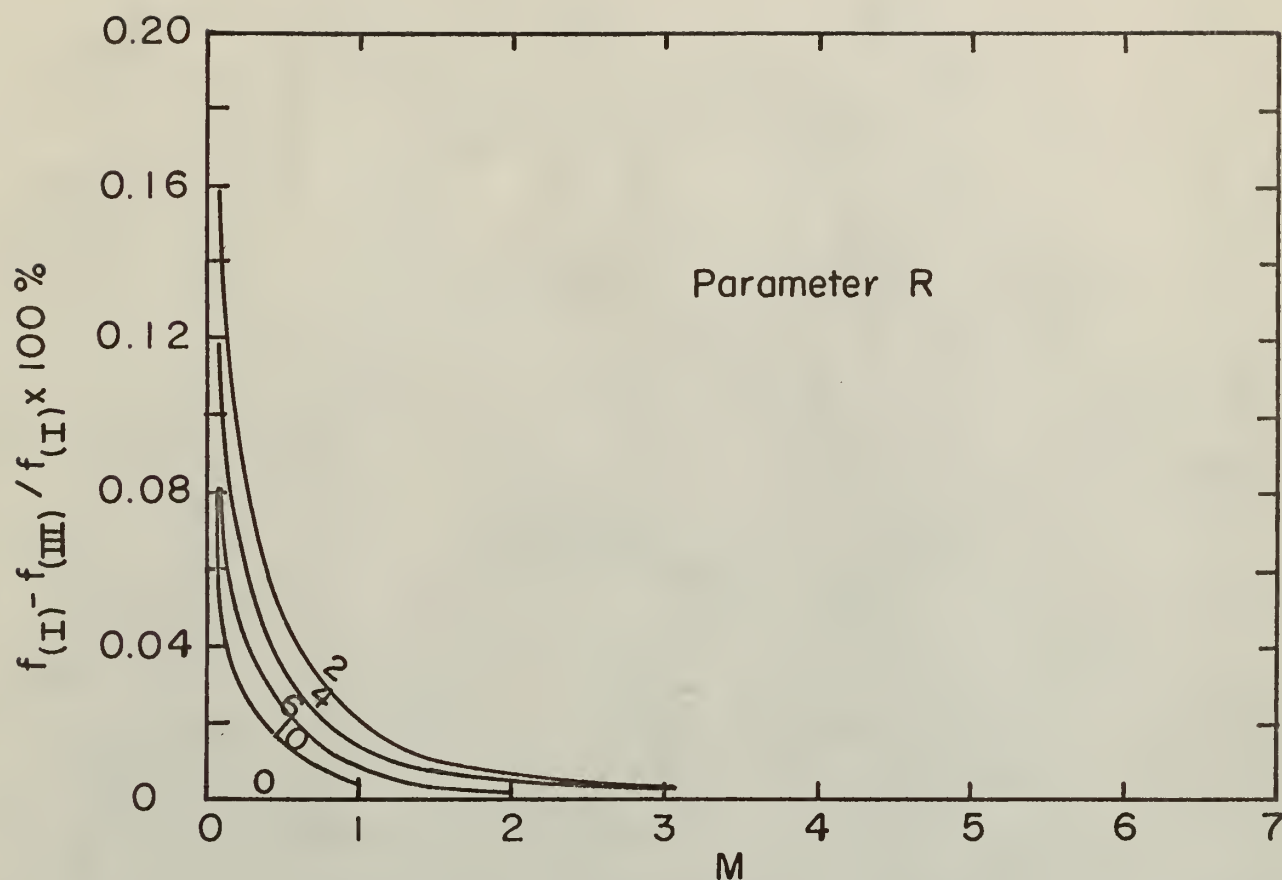


Fig. 9. Percentage deviations of final conversions computed by use of boundary conditions (III) from those computed by use of boundary conditions (I) as functions of M for first-order reactions.

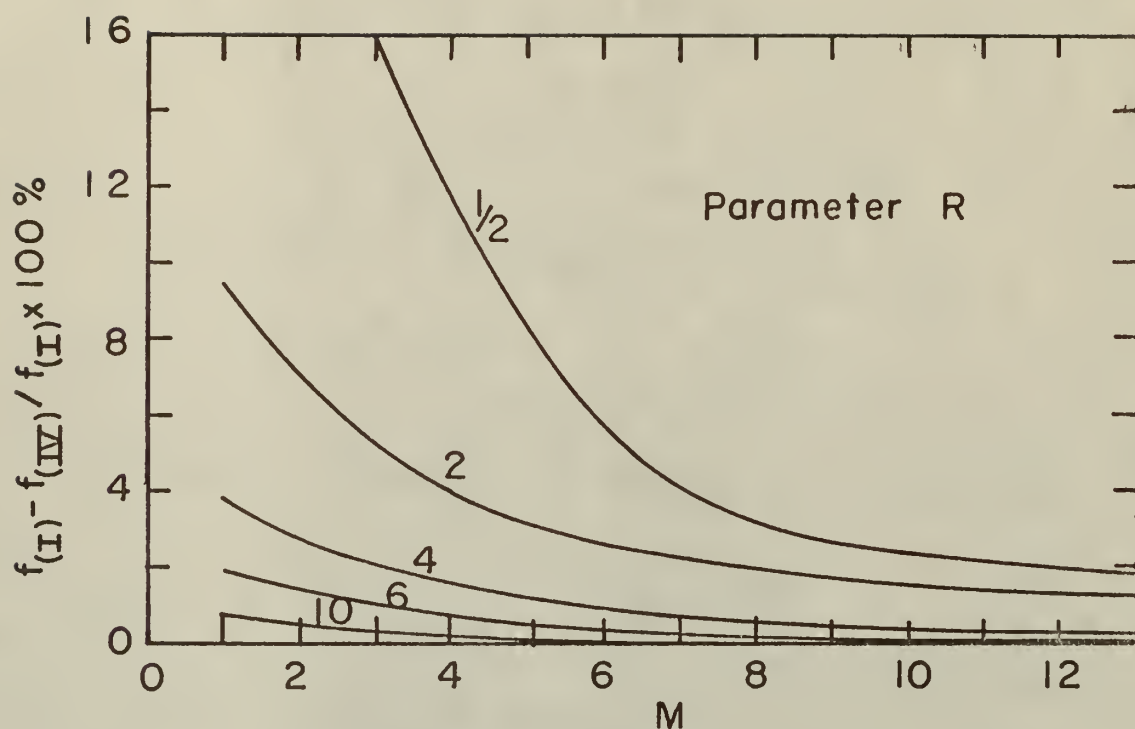


Fig. 10. Percentage deviations of final conversions computed by use of boundary conditions (IV) from those computed by use of boundary conditions (I) as functions of M for first-order reaction.

function, their use to approximate B.C.(I) is always more convenient, and also yields more accurate results than the use of B.C.(II).

The establishment of regions of inconsistency for the use of different sets of boundary conditions gives a set of criteria for the use of B.C. (II), (III) and (IV). These boundary conditions should be avoided in the domain of inconsistency for computing final conversions. Due to ease of solving simple exponential functions, the use of B.C.(III) or (IV), if their limits are recognized, offers the most convenient means for computing the final conversions. The recognition of the limit, of course, refers to the domain above the curves in Figs. 4 and 5. The use of B.C.(II) for a homogenous tubular flow reactor with axial dispersion is not desirable from the view point of both accuracy and convenience. Though the use of B.C.(IV) has never been attempted previously, it is useful in the computation of conversions around the inlet and center of the reactor, since as shown in Figs. 6 and 7, the concentration profile computed by use of B.C.(IV) follows closely that computed by the use of B.C.(I), except at the exit of reactor. As stated previously, the simple analytical solution in exponential form for B.C.(IV) provides a reasonably accurate and convenient method of evaluating conversions.

Unsteady-state Tubular Flow Reactor

The numerical calculations of the impulse response function tabulated in Table 4 are relatively easy in comparison with the direct numerical solution of the original partial differential equation, equation (2). Some computational results of the impulse

response functions, both with and without chemical reaction, are presented in Figs. 11 through 18. These figures definitely indicate the differences of response based on different sets of boundary conditions. Four sets of impulse response curves at $R=0$ for high and low degree of dispersion are replotted in Figs. 19 and 20 to show that as the degree of dispersion becomes small, say of the order of $M=15$, all the response curves representing the four boundary conditions approximately converge to one and the deviations among the four response curves increase as degree of dispersion increases. This fact supplements the previous statement that the effect of boundary conditions becomes negligible above the value of $M=15$. Exactly the same observations are given in Figs. 15 through 18 where the response curves of first-order chemical reaction are considered, except that the magnitude of the response curves decrease rapidly with increase in R . This is due to the fact that the magnitude of R is proportional to the rate of reaction; thus the faster the reaction the faster the concentration decreases.

One application of such computed response curves is the determination of the axial dispersion coefficient from a dynamic experiment. A survey of published literature shows that most of the methods used to determine the mean dispersion coefficient from transient response experiments makes use of only a single point of an entire experimental response curve, such as the maximum point or the slope at unit time (4) (25). It is known that the use of a single datum in drawing conclusions for an entire experiment may lead to serious error because of probable experimental error.

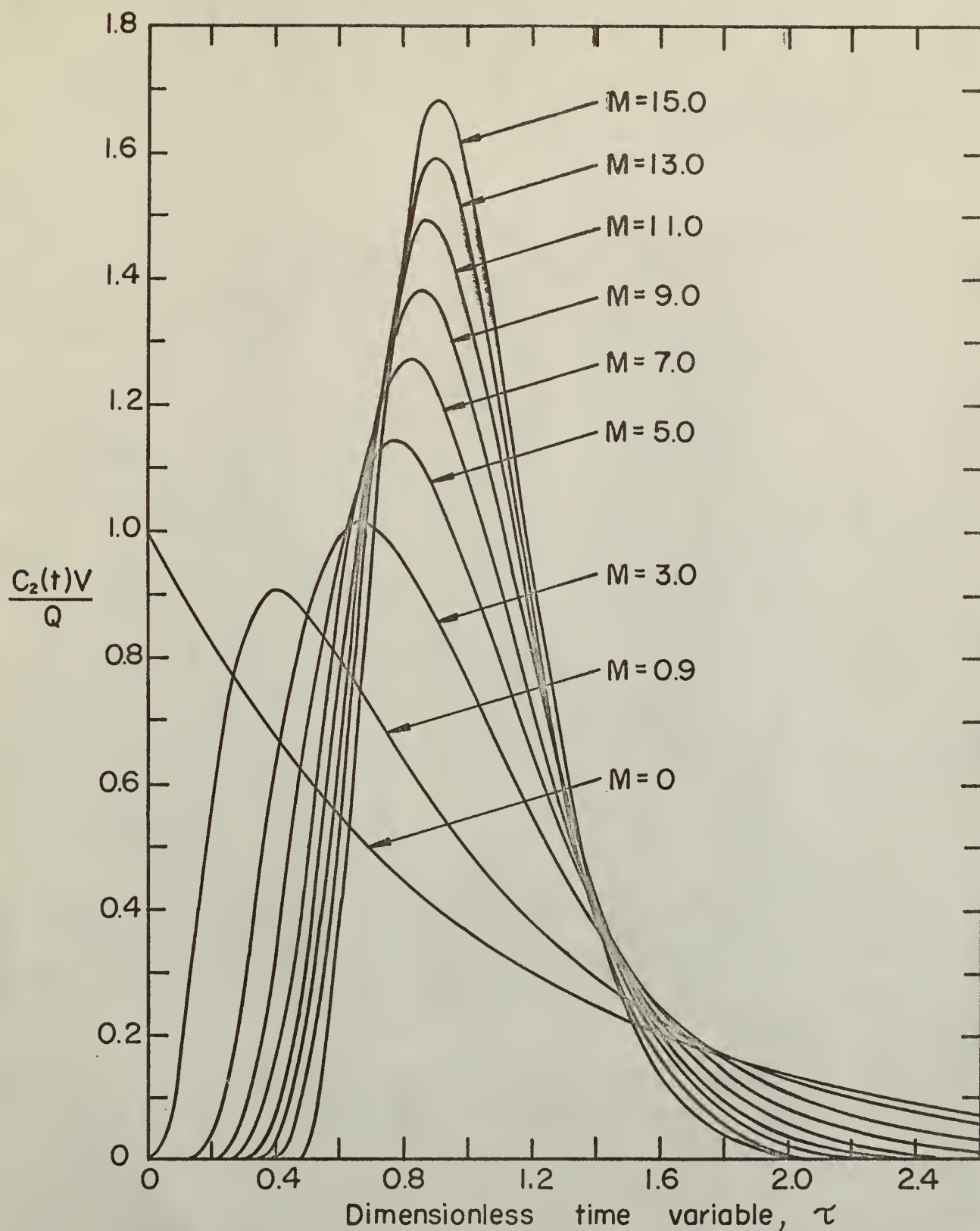


Fig. 11. Calculated response curves of exit concentration against impulse for non-reacting systems ($R=0$) based on B. C. (I) with M as parameter.

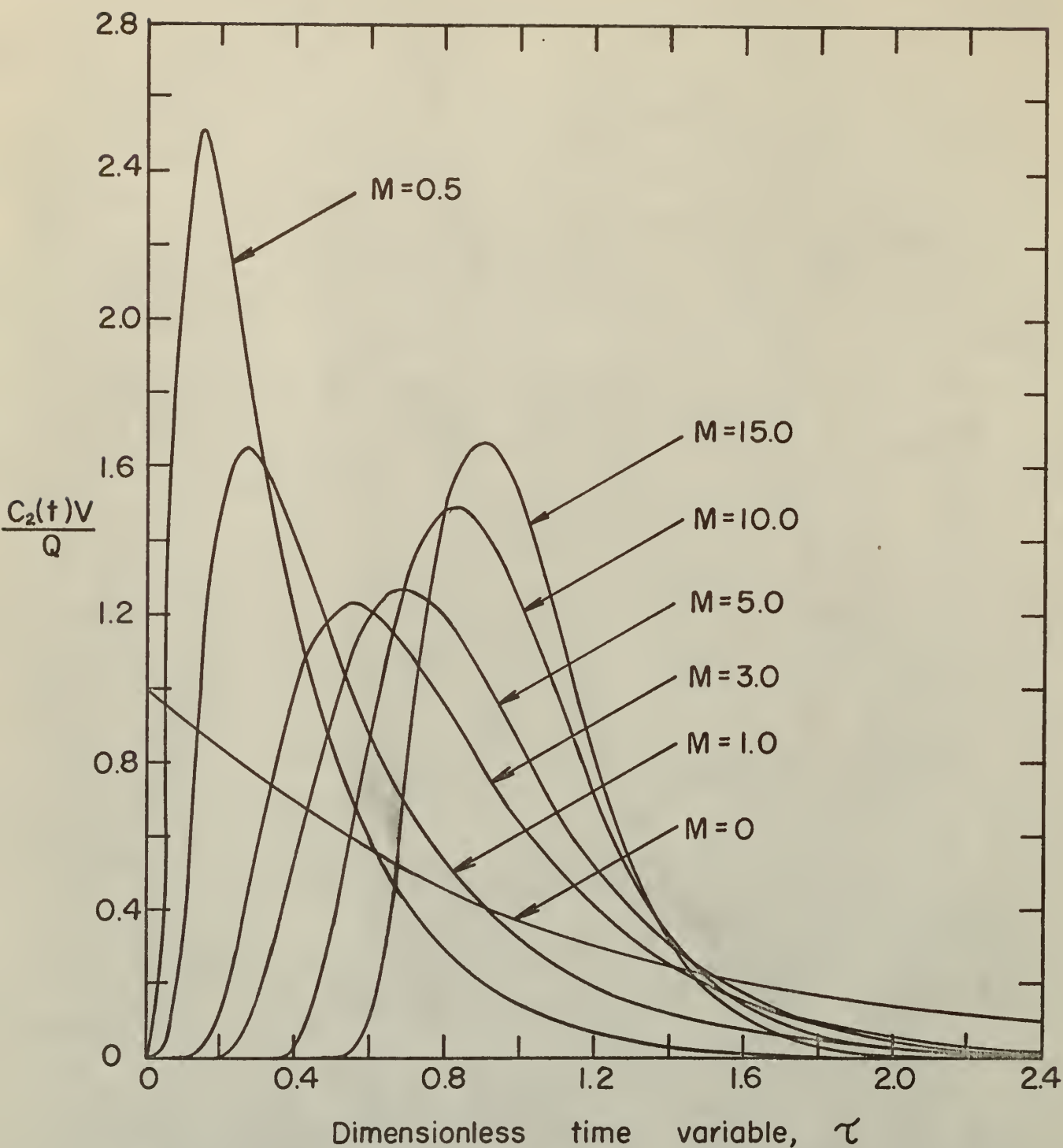


Fig. 12. Calculated response curves of exit concentration against impulse for non reacting systems ($R=0$) based on B. C. (II) with M as parameter.

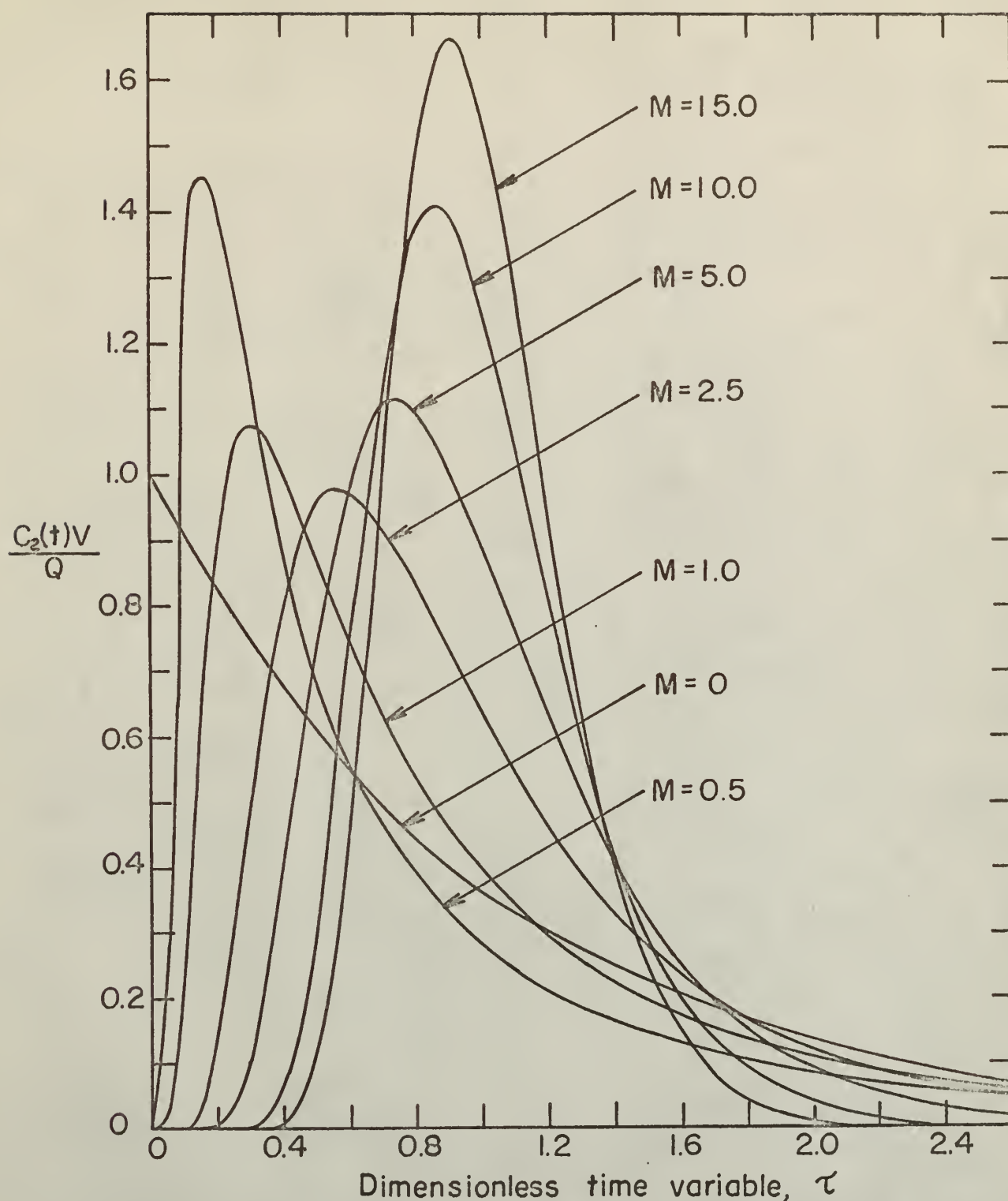


Fig. 13. Calculated response curves of exit concentration against impulse for non-reacting systems ($R=0$) based on B. C. (III) with M as parameter.

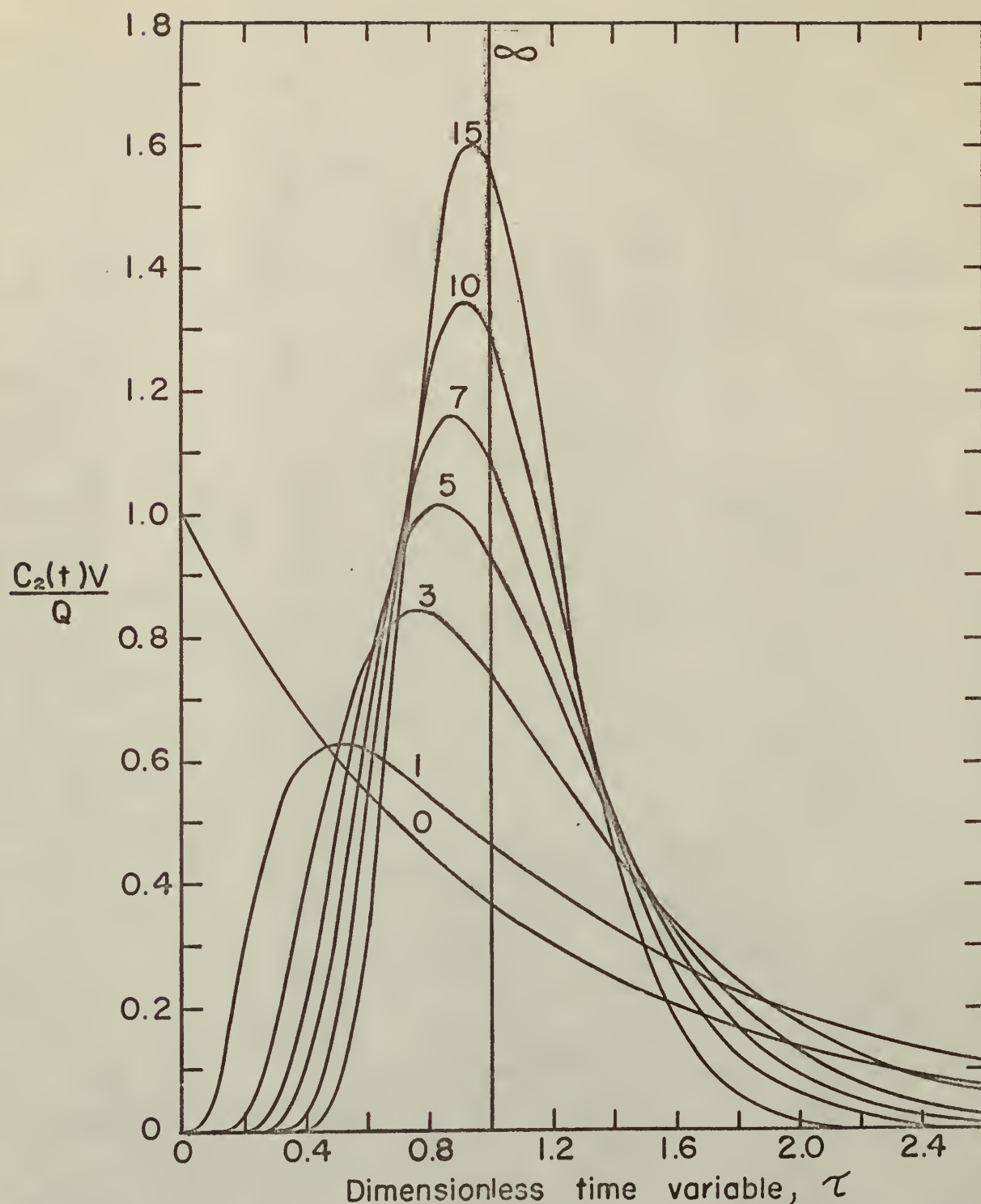


Fig. 14. Calculated response curves of exit concentration against impulse for non-reacting systems ($R=0$) based on B.C.(IV) with M as parameter.

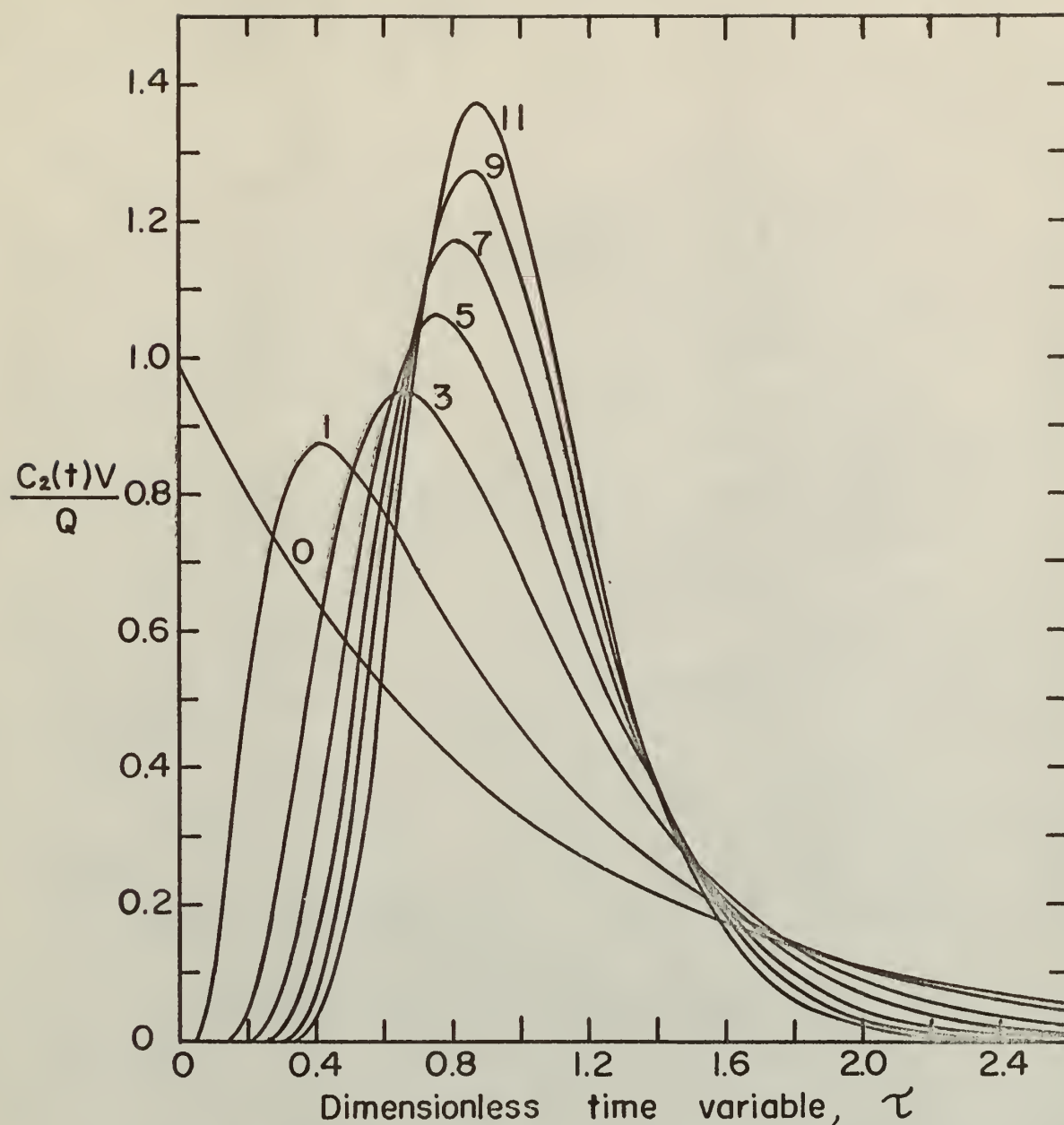


Fig.15. Calculated response curves of exit concentration against impulse for reacting systems ($R=0.1$) based on B.C.(I) with M as parameter.

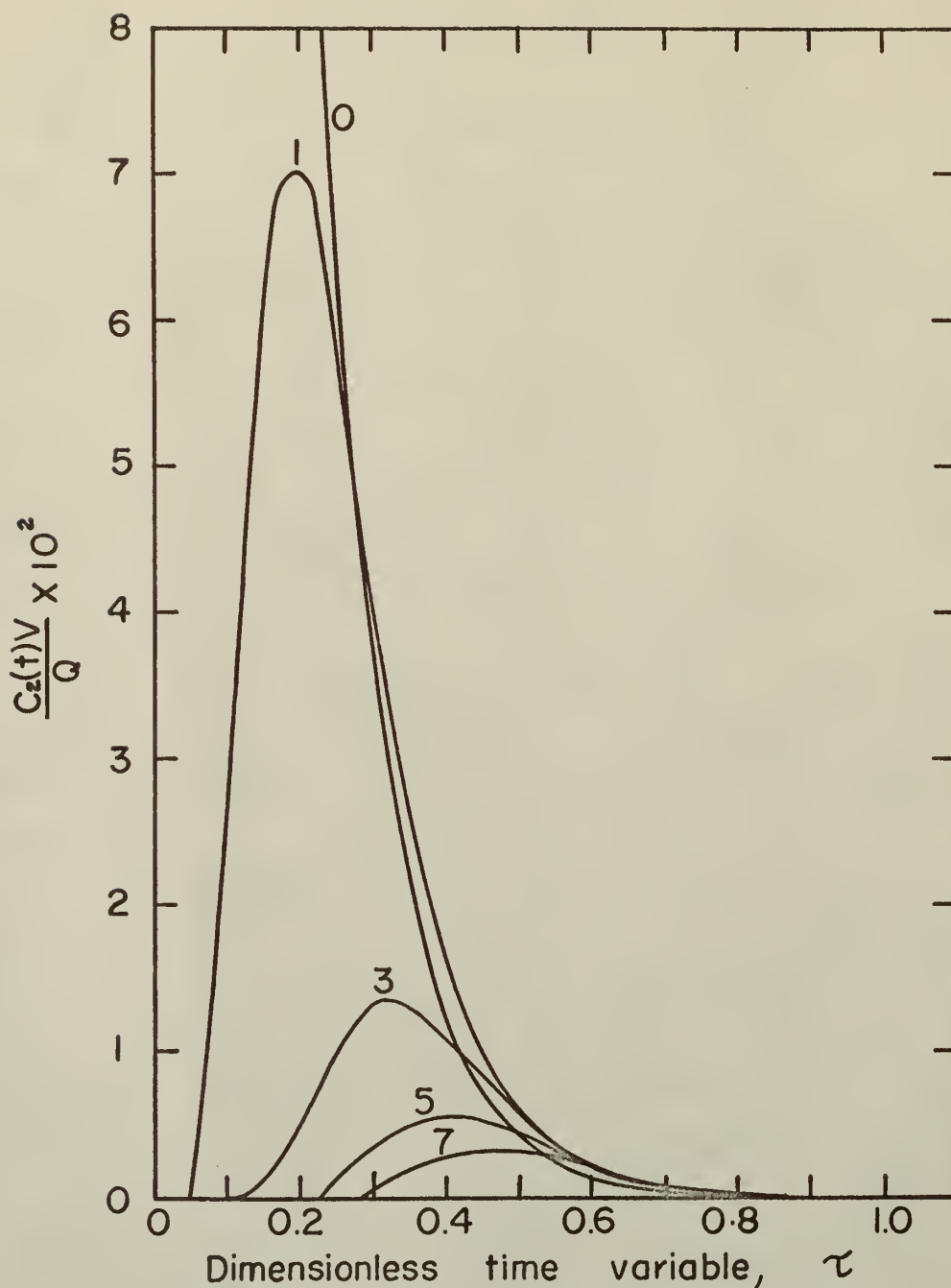


Fig. 16. Calculated response curves of exit concentration against impulse for reacting systems ($R=10$) based on B.C. (I) with M as parameter.

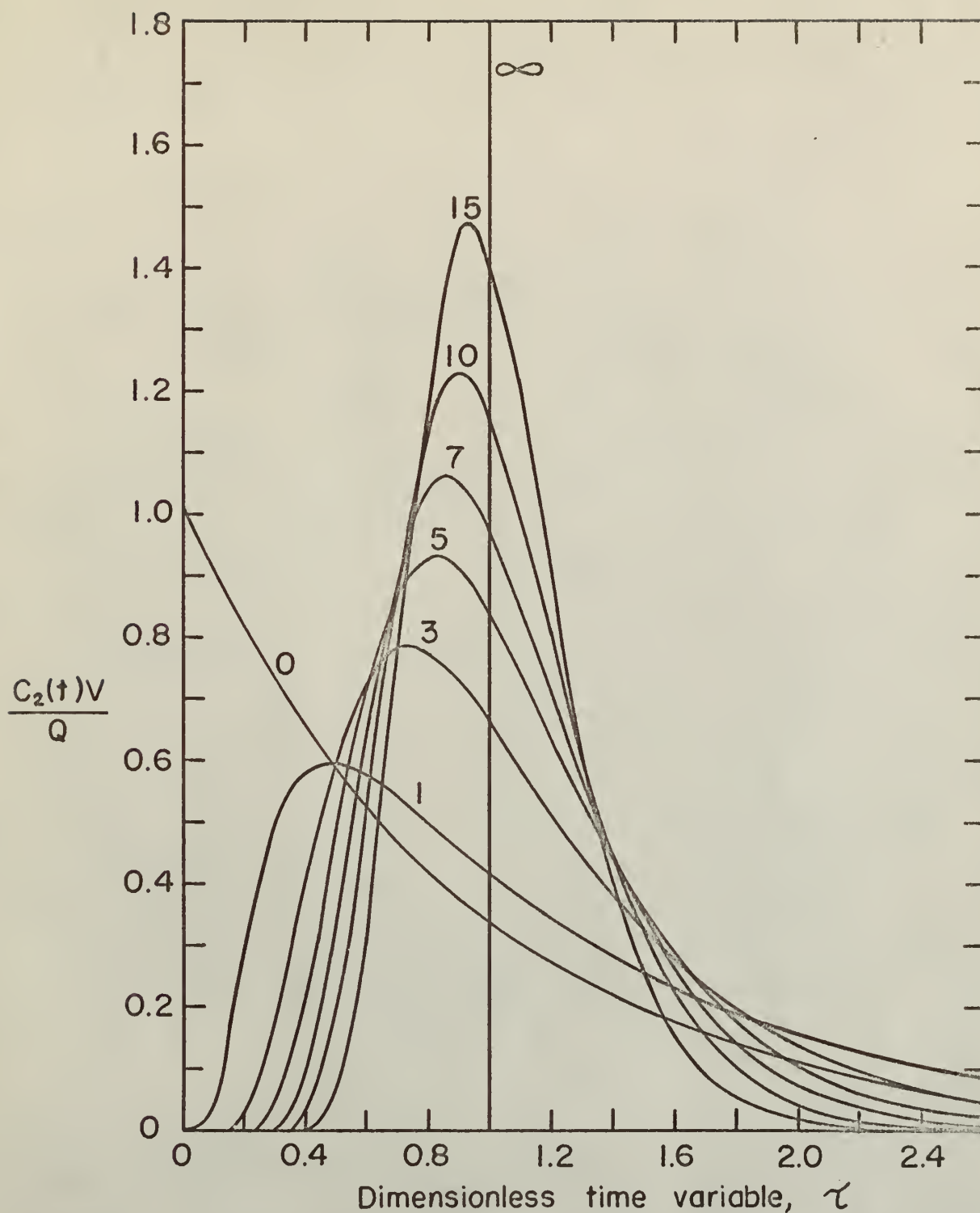


Fig. 17. Calculated response curves of exit concentration against impulse for reacting systems ($R=0.1$) based on B.C. (IV) with M as parameter.

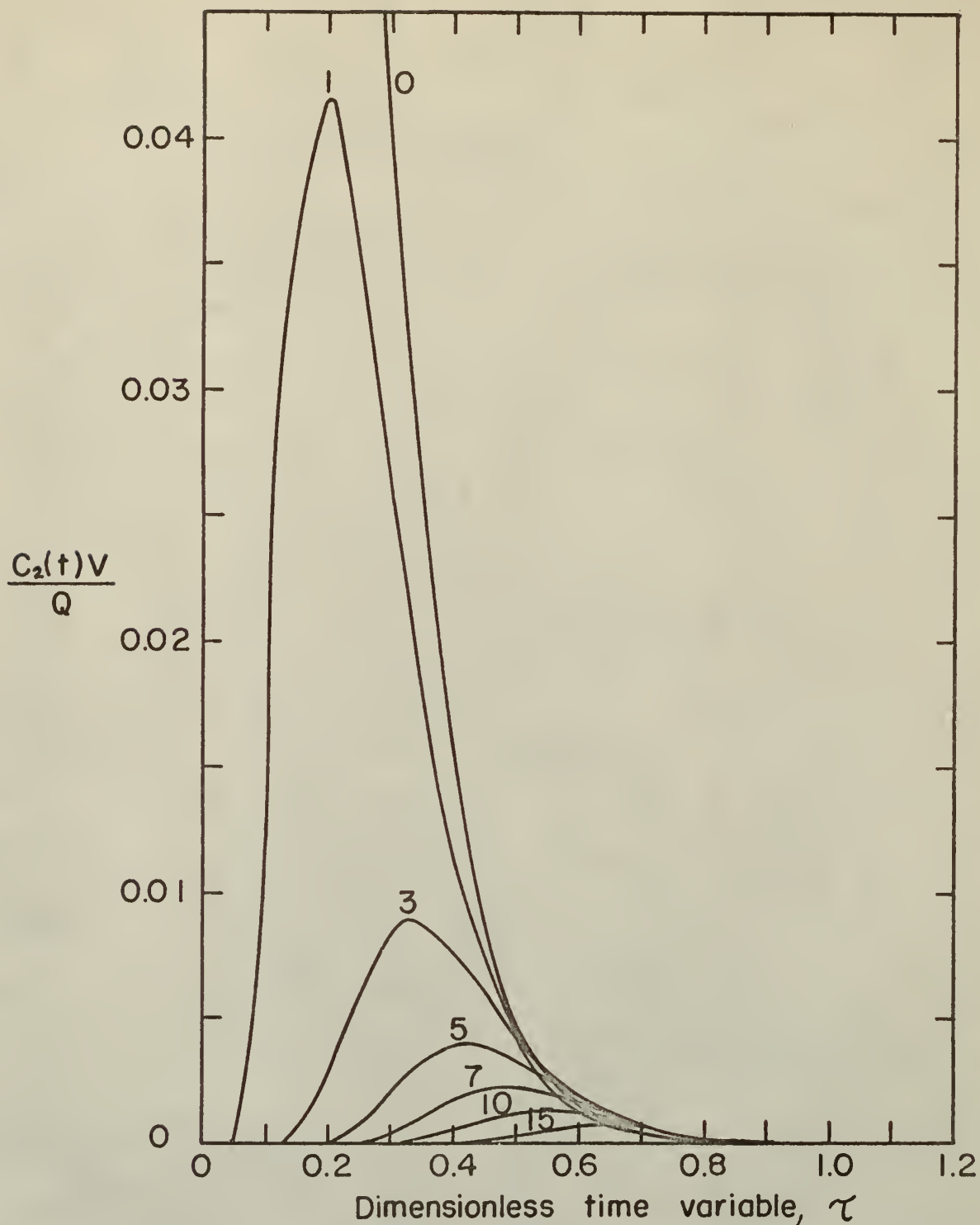


Fig. 18. Calculated response curves of exit concentration against impulse for reacting systems ($R=10$) based on B.C. (IV) with M as parameter.

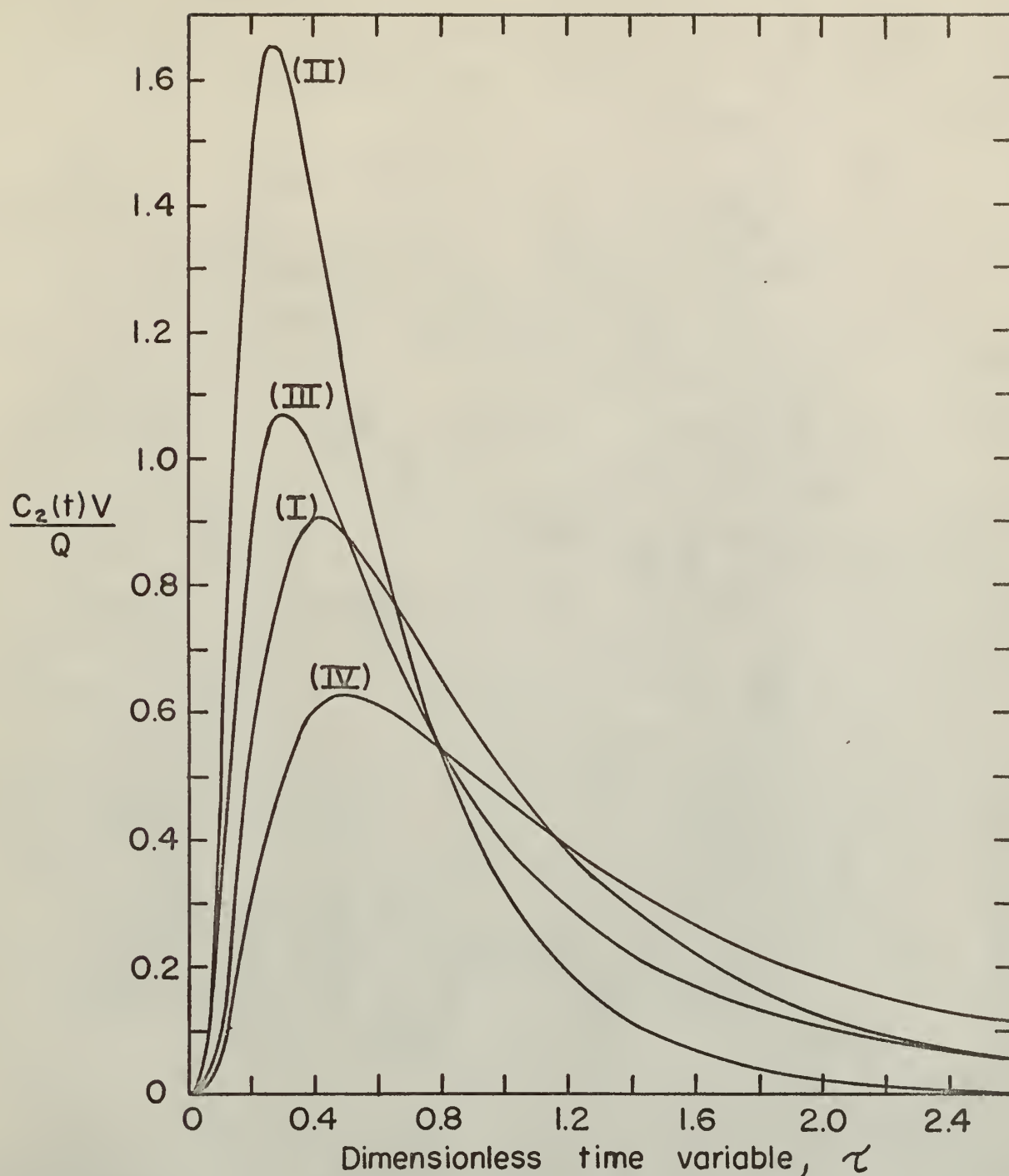


Fig. 19. Comparisons of response to impulse for non reacting systems ($R=0$) based on different boundary conditions at high degree of dispersion ($M=1$).

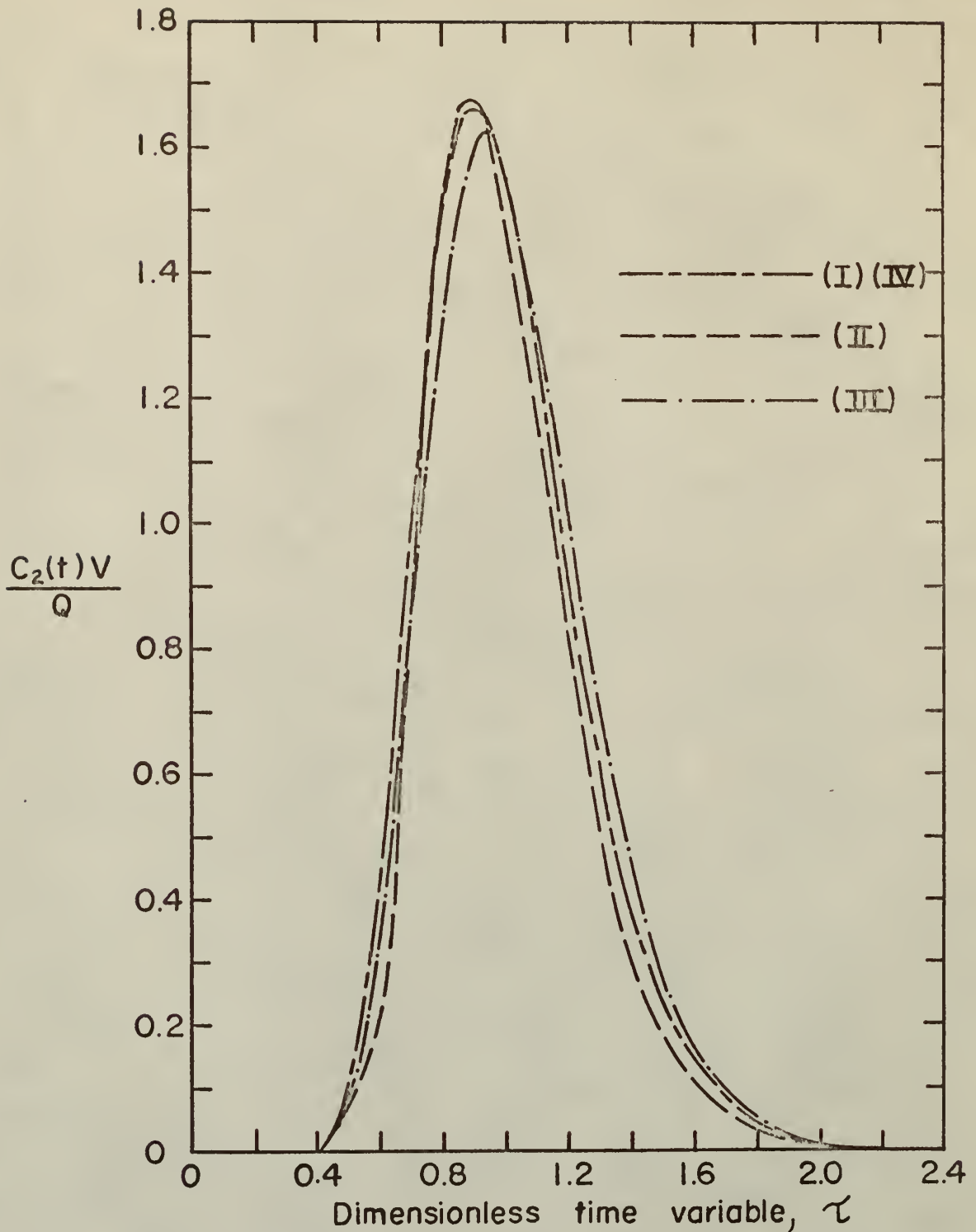


Fig. 20. Comparisons of response to impulse for non-reacting systems ($R=0$) based on different boundary conditions at low degree of dispersion ($M=15$).

Another approach in estimating \bar{D} is reported by van der Laan (21), Levenspiel and Smith (12) who introduced the concept of the variance of the concentration distribution by momentum analysis. This method has the advantage by avoiding the cumbersome back transformation of the Laplace transform. The disadvantage, however, is that one cannot explain how distribution functions for residence times can be defined and measured for actual systems. Since the overall flow behavior problem is often of prime concern in the design of continuous processing equipment, one should compare response curves calculated from the assumed models and boundary conditions with the experimental curves.

Since any model of fluid dispersions, including the diffusion-type model, is approximate in nature, the best procedure for estimating the mean dispersion coefficient based on any one of the models should be the comparison of the calculated response curves and an experimental response curve. Such a procedure is illustrated in Fig. 21.

The experimental data used in Fig. 21 were obtained from Run 194 of Fowler and Brown (6). Levenspiel and Smith (12) obtained the value of mean dispersion coefficient of $2.28 \text{ ft}^2/\text{sec.}$ for this experiment by the methods of both maximum point and variance of concentration distribution curve. Inspection of Fig. 21, however, shows that the comparison of the entire response curves yields mean dispersion coefficient of $2.7 \text{ ft}^2/\text{sec.}$ if the response curve based on B.C.(III) is used. It also indicates that the system represented by the experimental data may not be represented by B.C.(II) and (IV). In other words, the comparison of the

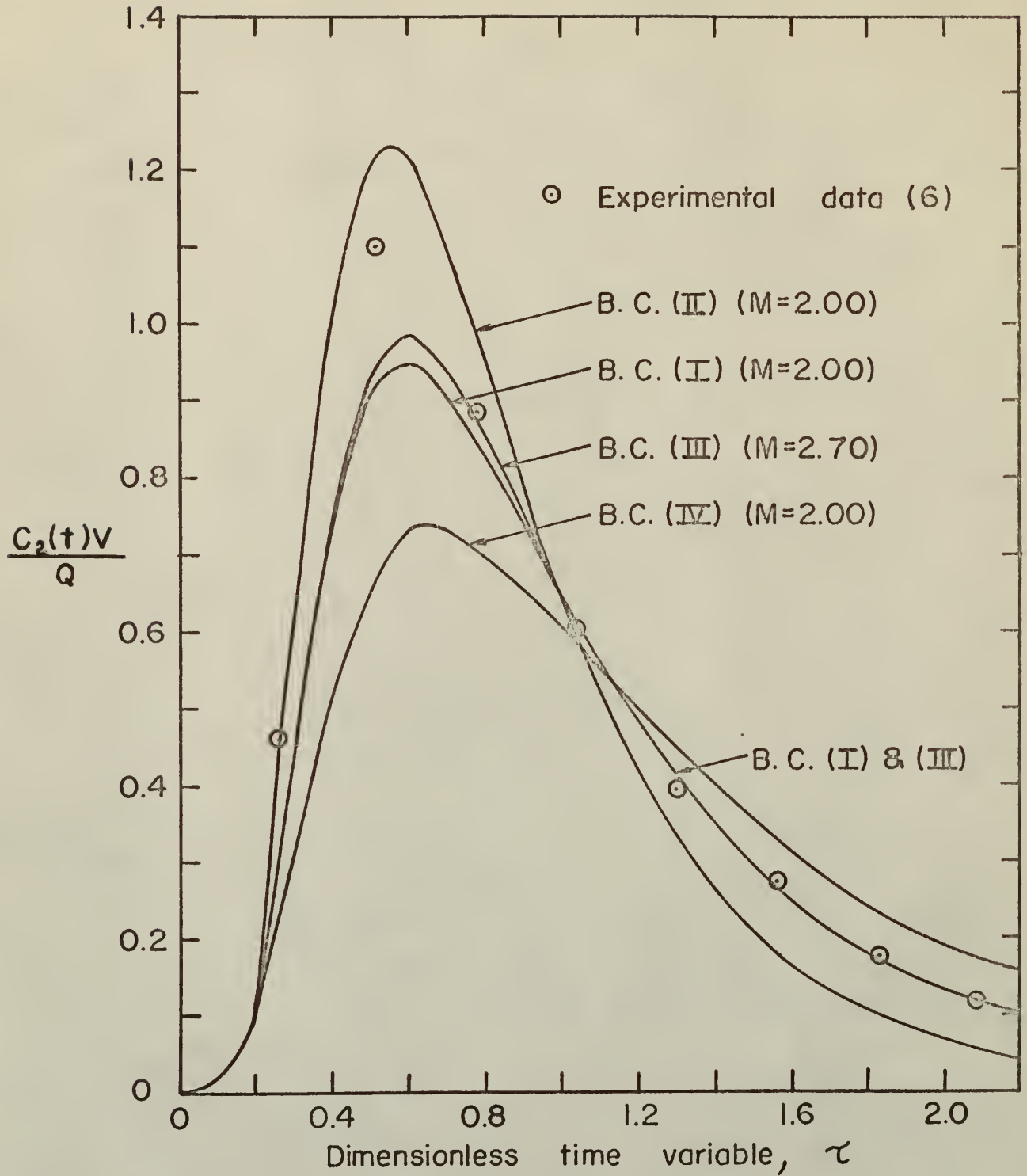


Fig. 21. Comparison of experimental and calculated response curves.

experimental and calculated response curves should lead to the establishment of the appropriate models or boundary conditions.

There are no data available on impulse experiments accompanied by first-order chemical reaction. If the data were available, the calculated response curves based on first-order chemical reaction (Figs. 15 through 18) could be useful to extract transport parameters from a dynamic experiment.

The typical results for step input response curves are shown in Figs. 22 through 25 both with and without chemical reaction. Due to the consistency observed in the steady-state concentration profile with the limits of both stirred tank and plug flow reactor, only the response curves based on B.C.(I) are presented here. The response curves with first-order chemical reaction do not approach one as τ approaches infinity because of the steady-state conversion, i.e., the response curves approach steady-state conversion as τ approaches infinity.

One method of determining \bar{D} from step input response curves is illustrated by Levenspiel and Smith (12). They replotted the data of Von Rosenberg (23) on probability paper to obtain a straight line. From the property of normal distribution that the interval between the 16th and 84th percentile points is two standard deviations, they evaluated the mean dispersion coefficient. This method requires an extra step, and also involves appreciable error because the actual experimental data does not give a symmetric normal distribution curve. The best procedure for determining \bar{D} , again therefore, should be by matching the calculated response curves and an experimental response curve.

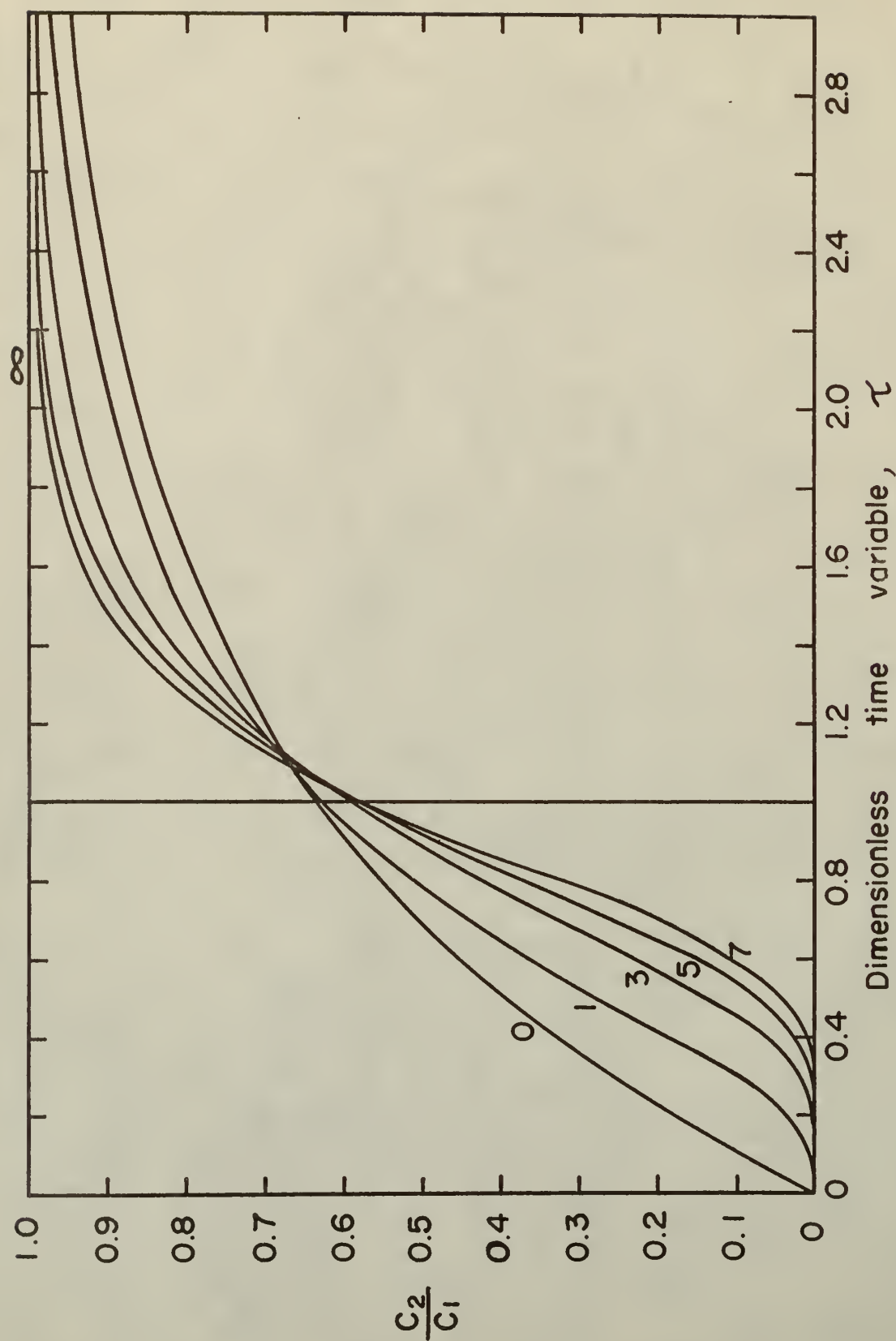


Fig. 22. Calculated response curves of exit concentration against step input for non-reacting systems ($R=0$) based on B. C. (I) with M as parameter.

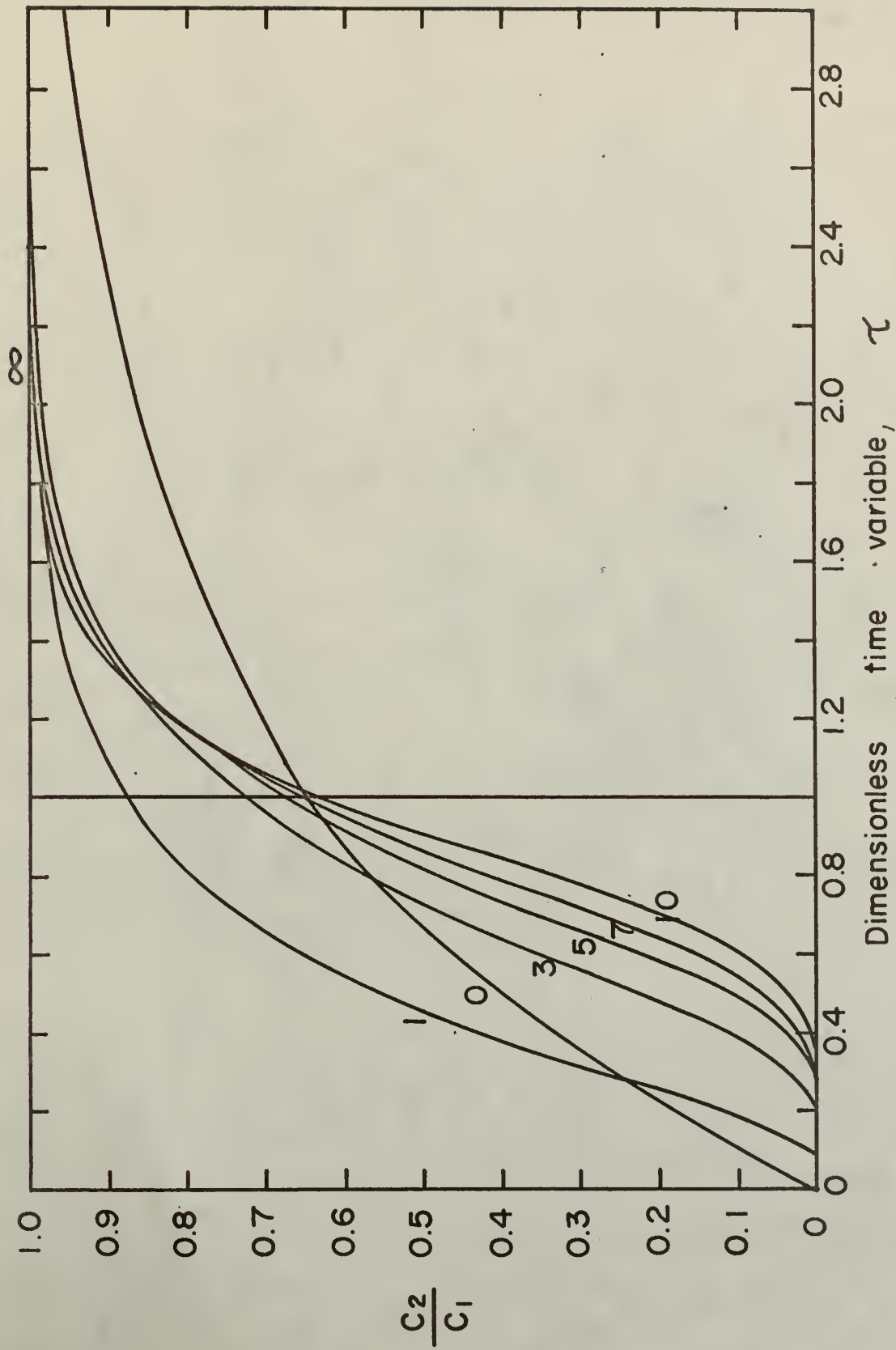


Fig. 23. Calculated response curves of exit concentration against step input for non reacting systems ($R=0$) based on B. C. (II) with M as parameter.

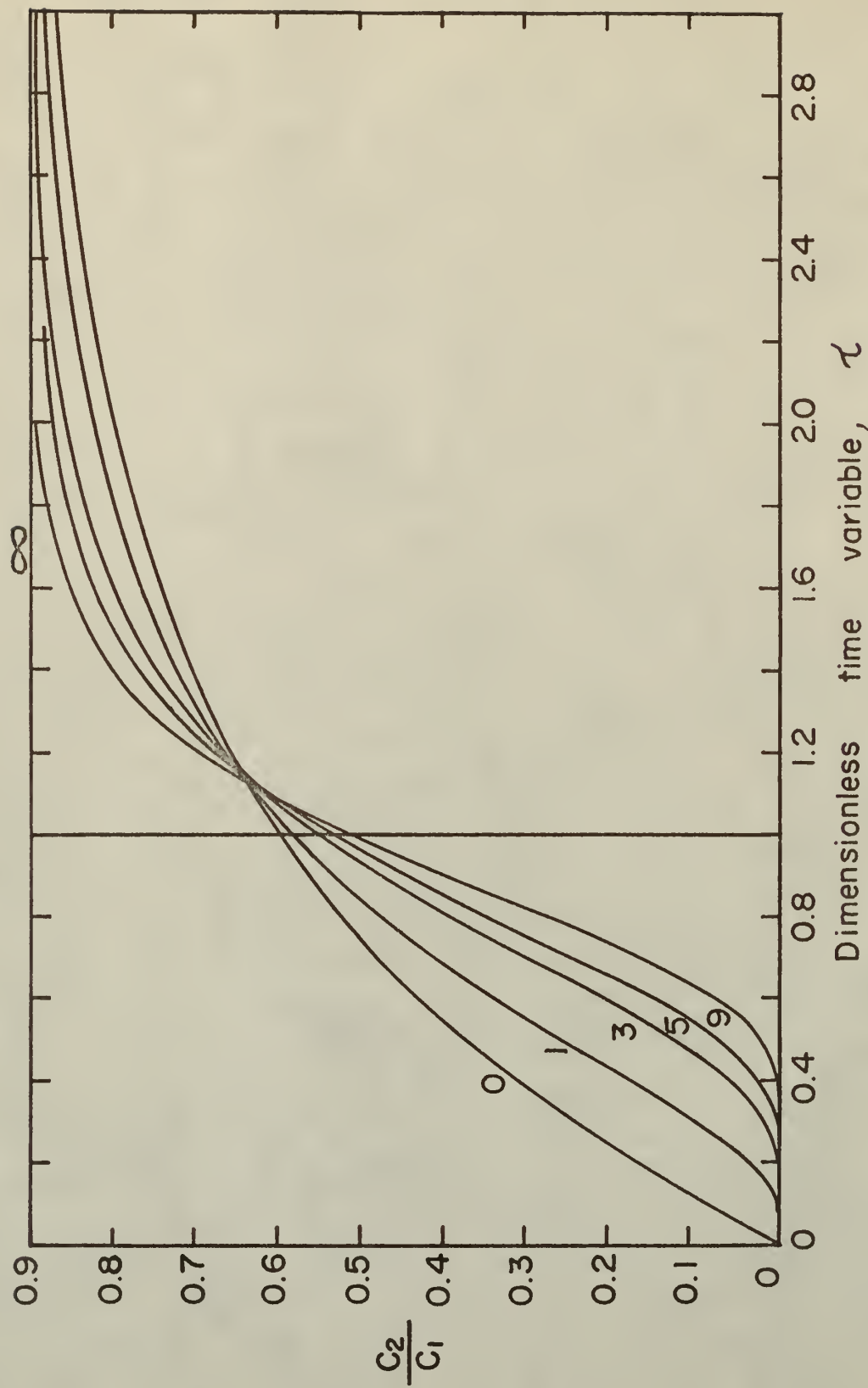


Fig. 24. Calculated response curves of exit concentration against step input for reacting systems ($R=0.1$) based on B. C. (I) with M as parameter.

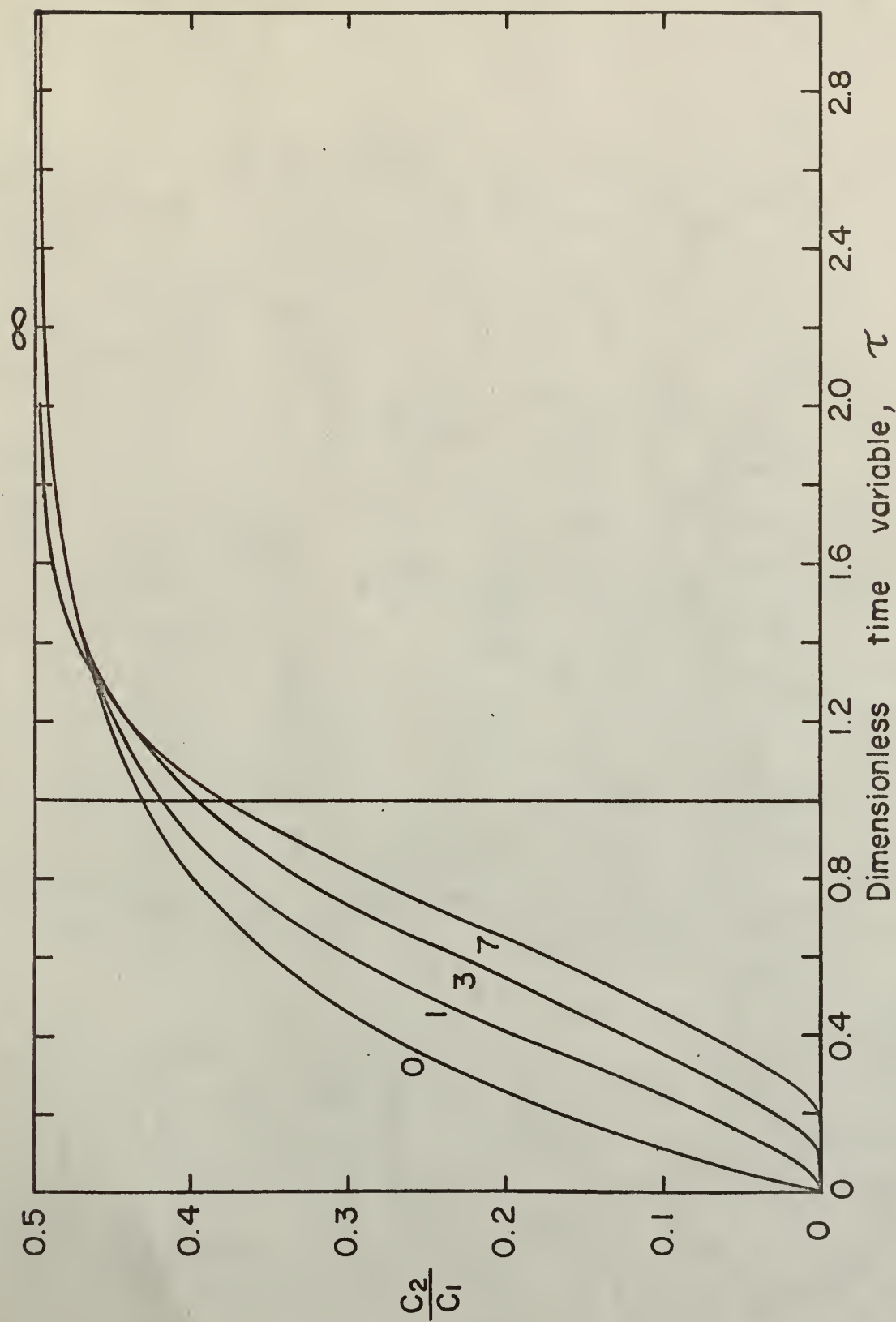


Fig. 25. Calculated response curves of exit concentration against step input for reacting systems ($R=1$) based on B.C. (I) with M as parameter.

The results of frequency response curves based on B.C.(I) (equation 31) are shown in Figs. 26 through 28 in the form of dimensionless Bode plots. Frequency responses of the two limiting conditions of the dispersion model based on equations (38) and (43) are also shown in the figures for the sake of comparison. Only the frequency response curves based on B.C.(I) are presented herewith because of the consistency observed in the previous section with the two limiting cases. Other frequency curves based on different sets of boundary conditions are presented only for the sake of comparison in Figs. 29 through 34.

Figure 26 shows that the plots of amplitude ratios and phase shifts change monotonically between the two limits for the nonreacting systems in the range of $\omega\theta$ which is of practical interest (15). This, however, is no longer true for the reacting systems, as shown in Figs. 27 and 28. As the relative reaction rates characterized by R increase, amplitude ratio decreases rapidly, especially at large values of M . For example, at the conditions for which R and M are greater than 6 and 10 respectively, the amplitude ratio is very small and practically is negligible. Under these conditions of rapid reaction and small degree of dispersion, a cyclic unsteady-state system has little effect on the performance of the system; and therefore, the frequency response data are of little use. In contrast, the phase shift data are only slightly affected by the change in the values of M and R . Thus, it can be said that the phase shift data are of little value in determining the characteristics of the system.

The performance function represented by equation (118) (See

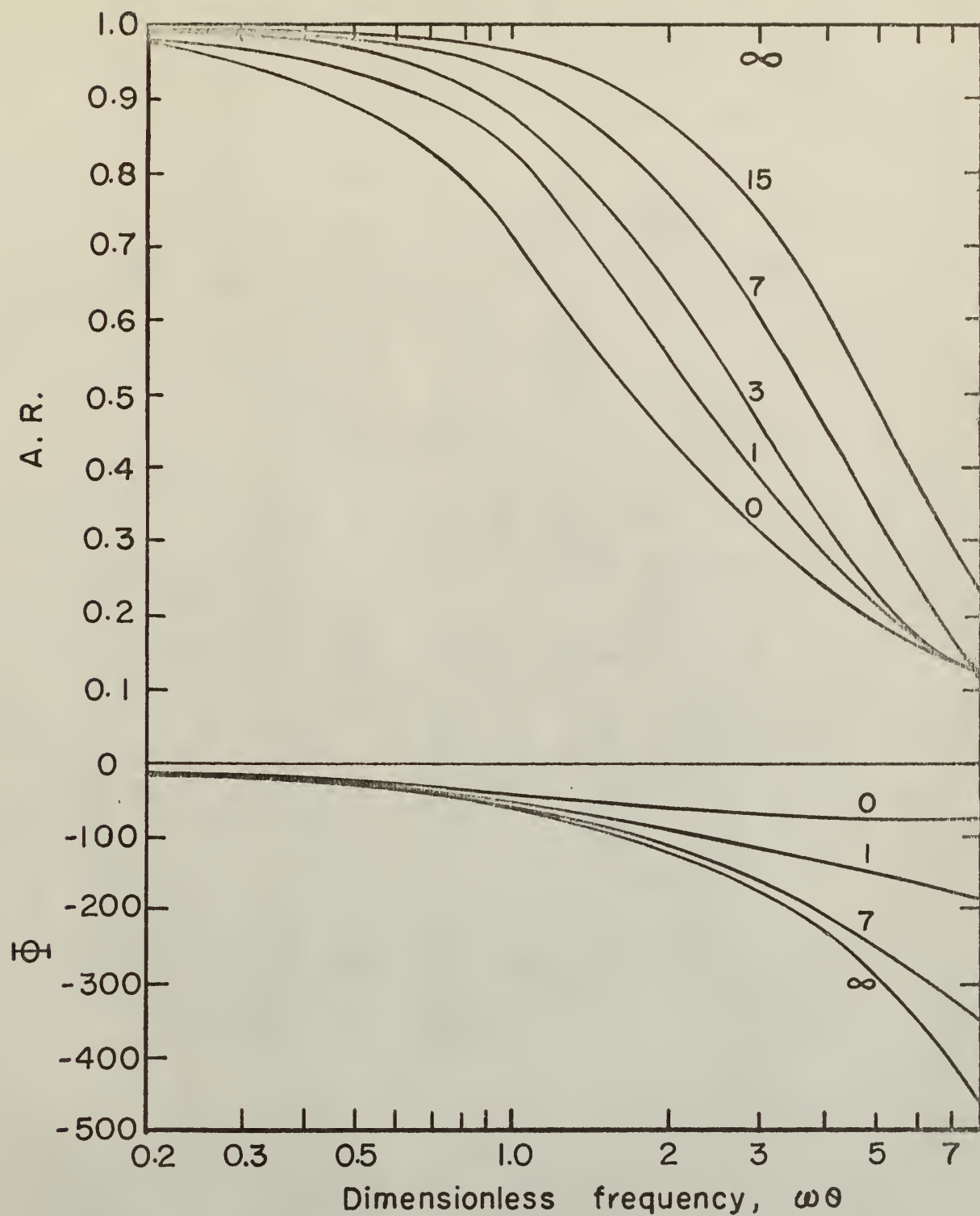


Fig. 26. Frequency response diagram of non-reacting system ($R=0$) based on B.C. (I) with M as parameter.

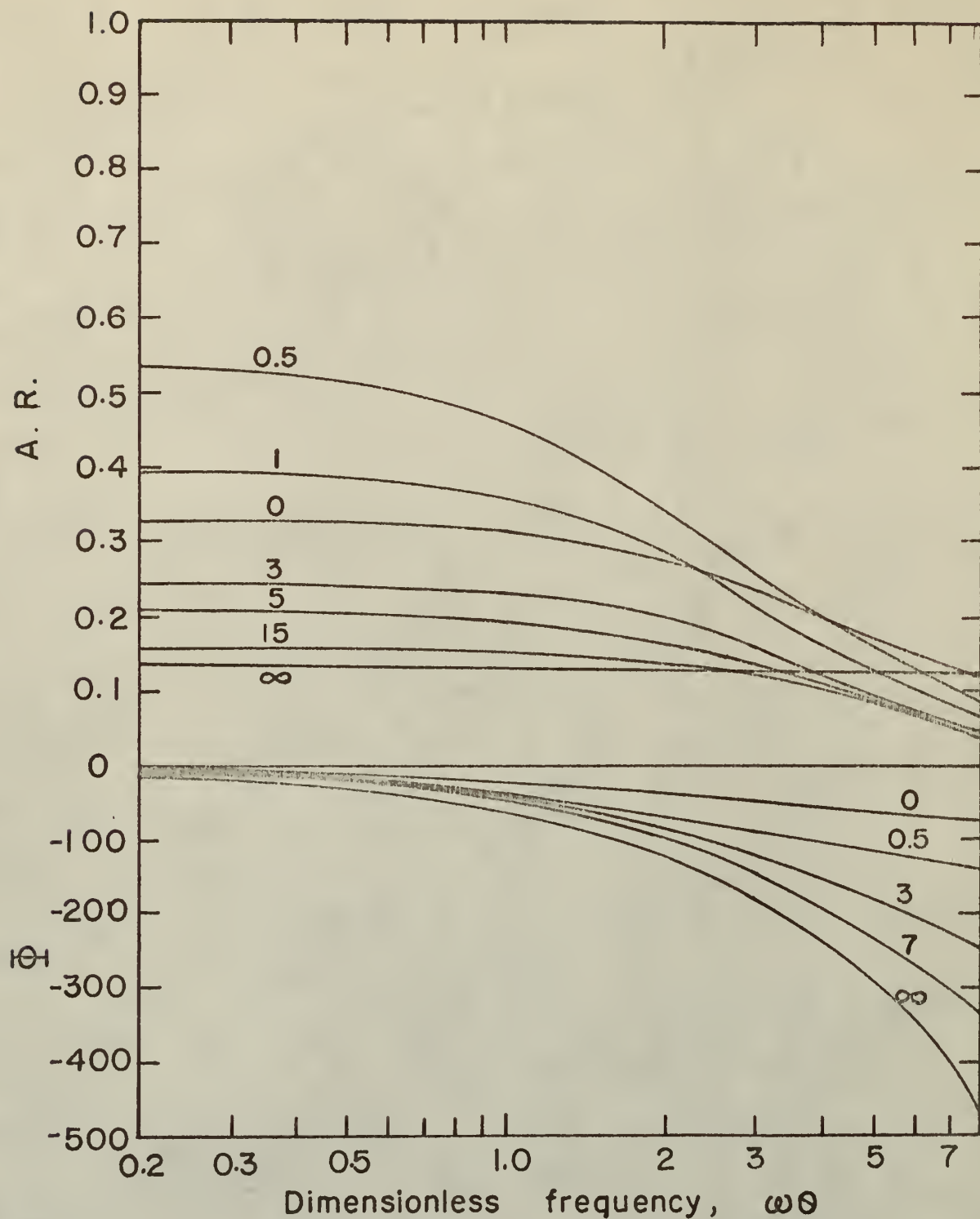


Fig. 27. Frequency response diagram of reacting system ($R=2$) based on B. C. (I) with M as parameter.

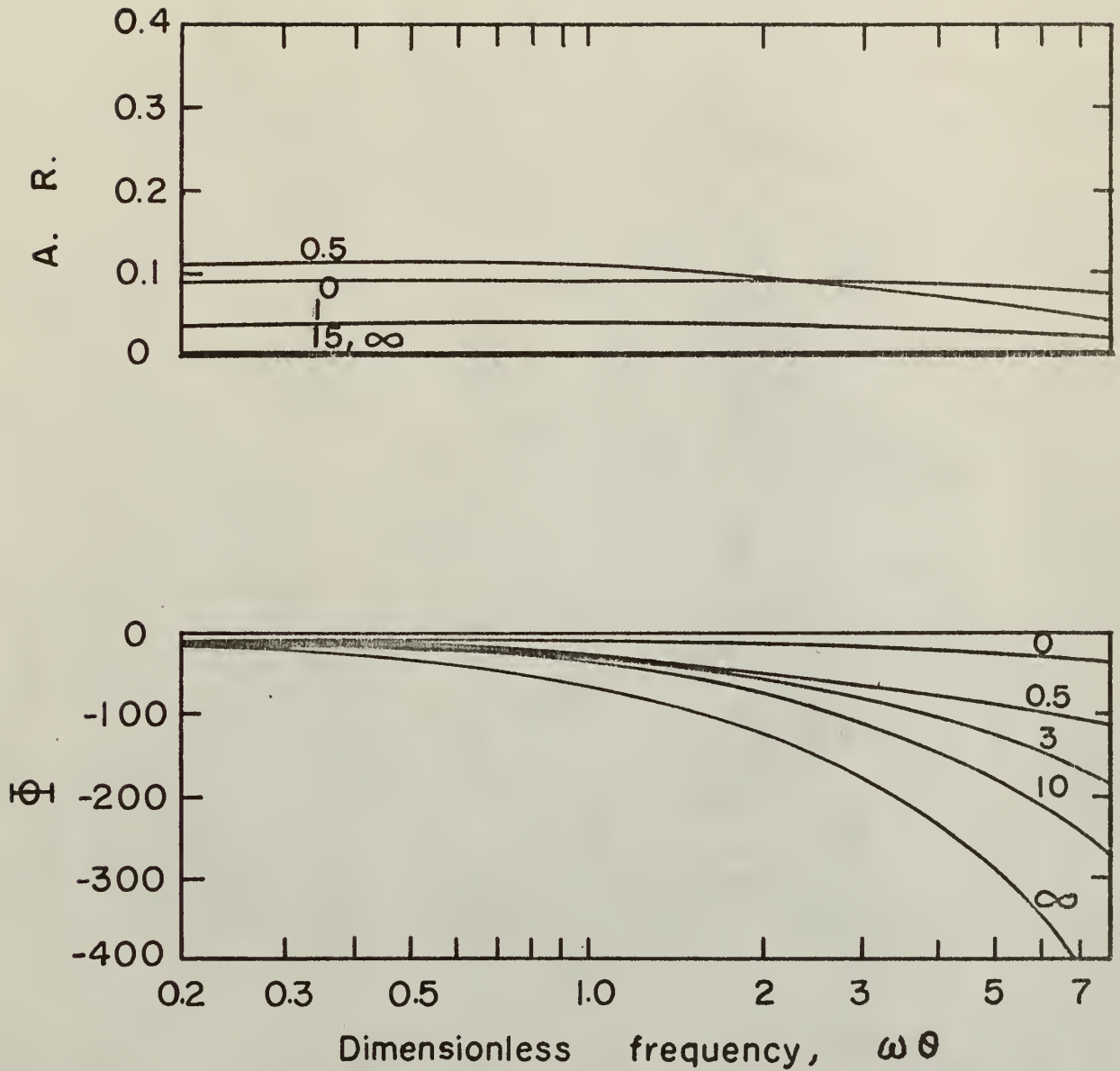


Fig. 28. Frequency response diagram of reacting system ($R=10$) based on B.C. (I) with M as parameter.

Appendix IV) was originally derived by Kramers and Alberda (10) for non-reacting systems. In order to avoid complexities in transforming equation (118) to the computable form of equation (31), it was suggested that equation (31) be approximated by the performance function of n numbers of the perfectly mixed vessels connected in series. Such approximation is good only for large values of M (c.a. $M = 15$). Very frequently approximate boundary conditions which do not conserve the continuity requirements, i.e., B.C.(II) and (III), are employed. Kramers and his coworkers (11) employed B.C.(II) and (III), and the evaluation of the performance function based on B.C.(III) (equation 108 of Appendix IV) has been frequently presented in the literature (7) (22).

The amplitude ratios and phase shift based on these different sets of boundary conditions are compared in Figs. 29 through 34 both with and without chemical reaction at high and low degree of dispersion. As can be seen from Fig. 29, the agreement of the amplitude ratios and phase shift curves for non-reacting systems based on approximate boundary conditions are quite poor at low values of M (large degree of dispersion). An increasing degree of agreement, however, is observed in Fig. 30 as M increases. The deviations among the phase shift curves are very small when M is approximately 10 or greater, and the amplitude ratio curves practically coincide when M is approximately 15 or greater. It is also seen that phase shift curves are generally less affected than amplitude ratio curves. The performance function curves of the reacting system, shown in Figs. 31 through 34, reveal the effects of using different sets of boundary conditions more clearly than

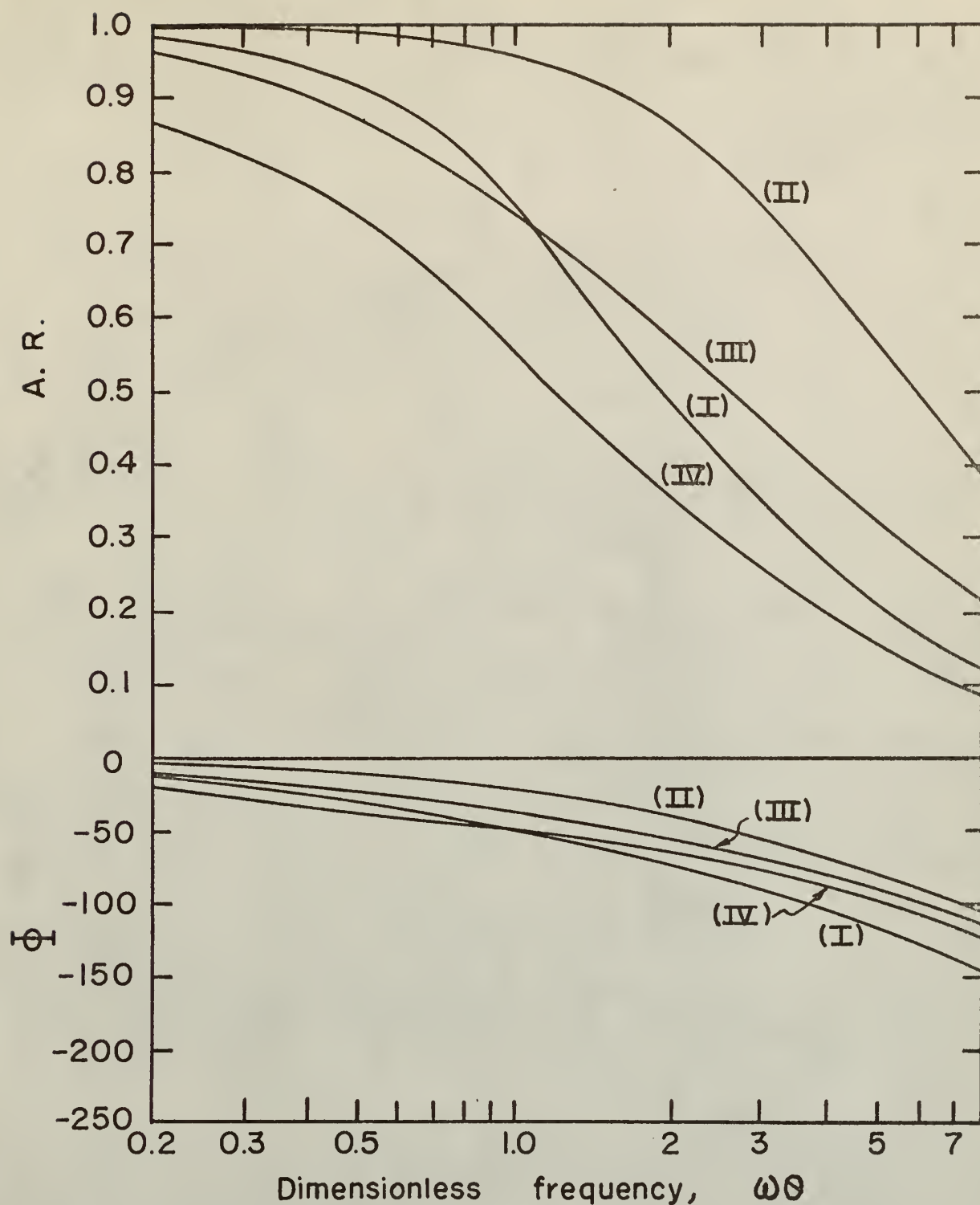


Fig. 29. Comparison of frequency response diagrams of non-reacting system ($R=0$) based on different boundary conditions at high degree of dispersion, $M=\frac{1}{2}$.

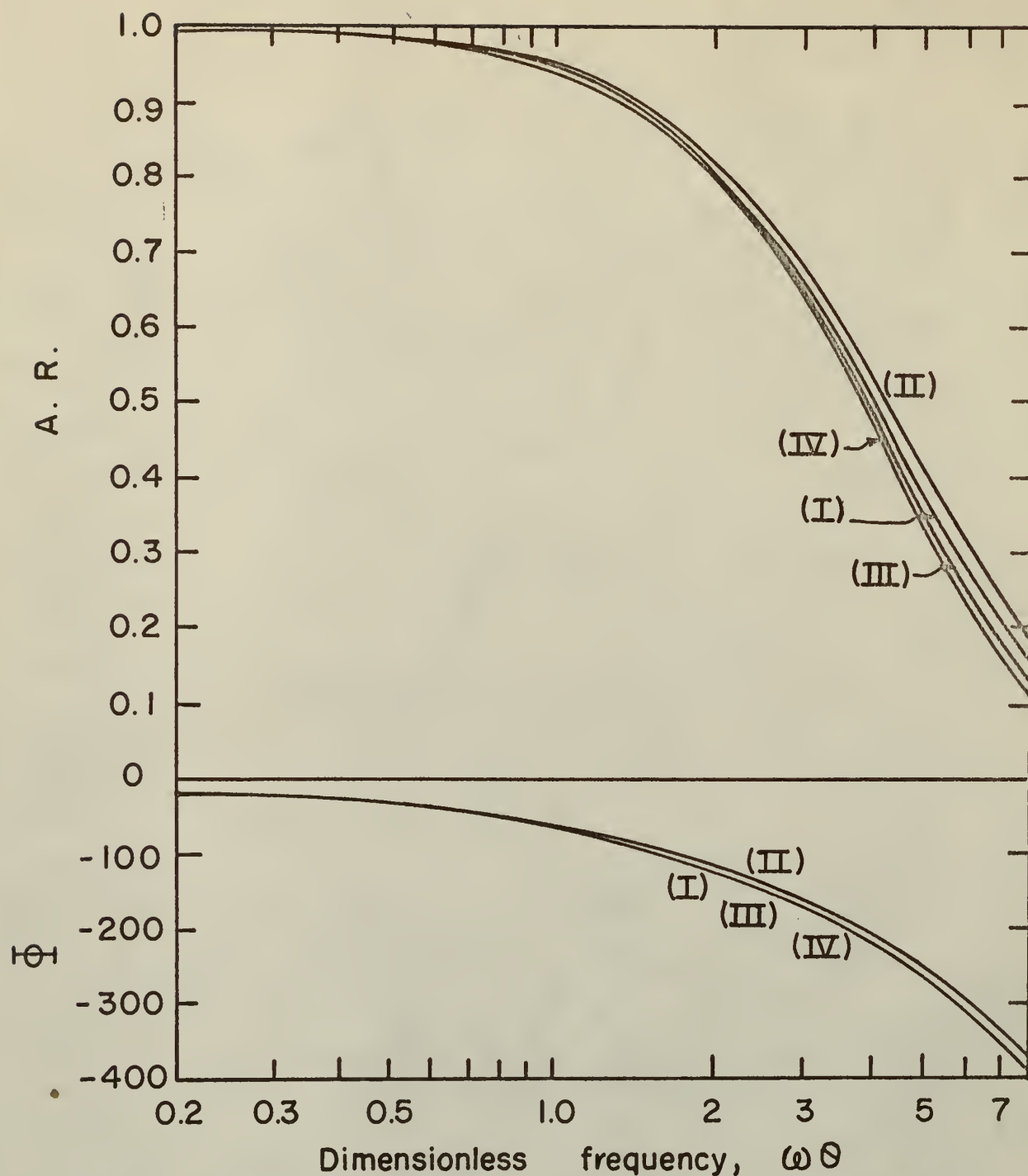


Fig. 30. Comparison of frequency response diagrams of non-reacting system ($R=0$) based on different boundary conditions at low degree of dispersion, $M=10$.

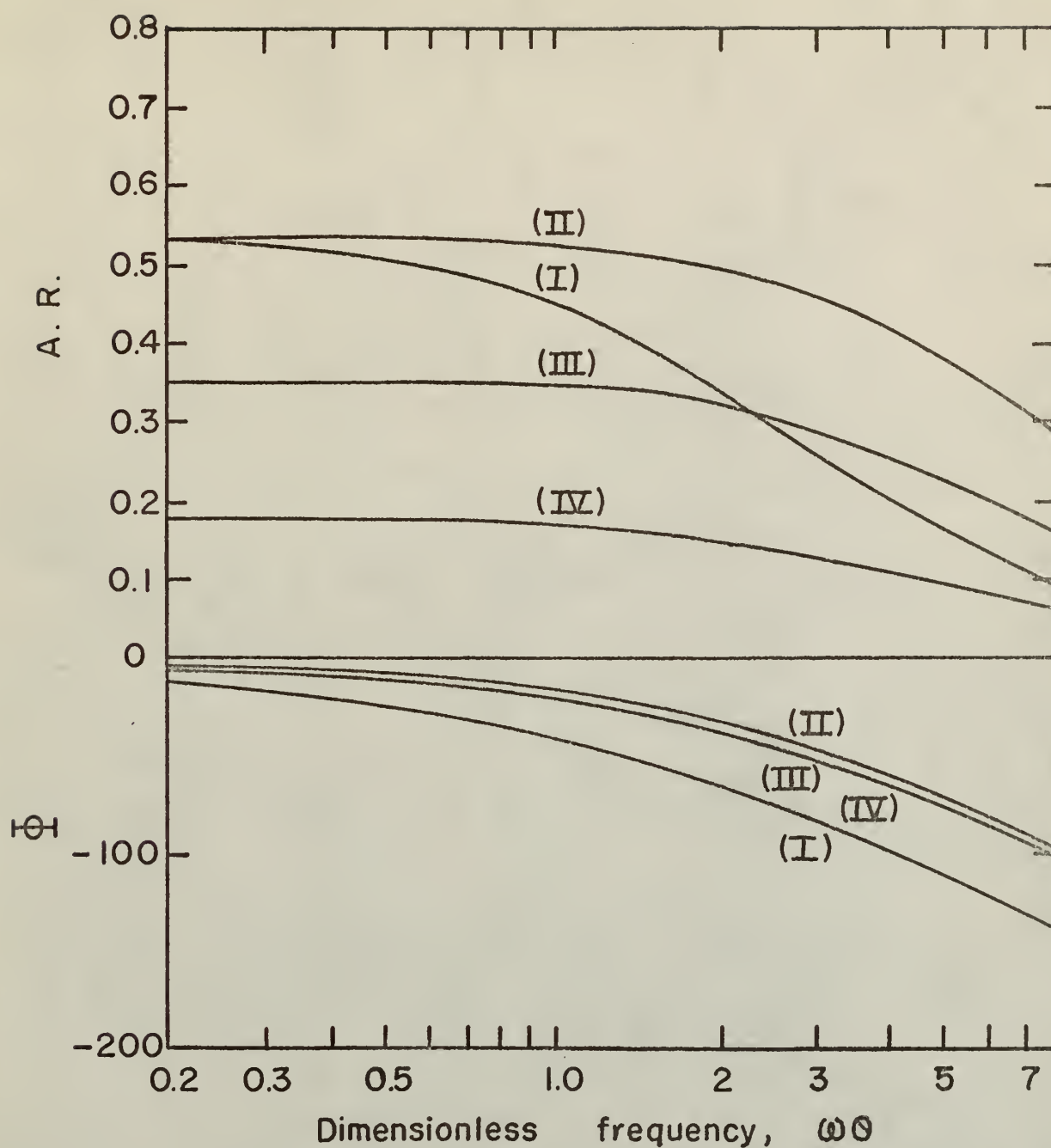


Fig. 31. Comparison of frequency response diagrams of reacting system ($R=2$) based on different boundary conditions at high degree of dispersion, $M = \frac{1}{2}$.

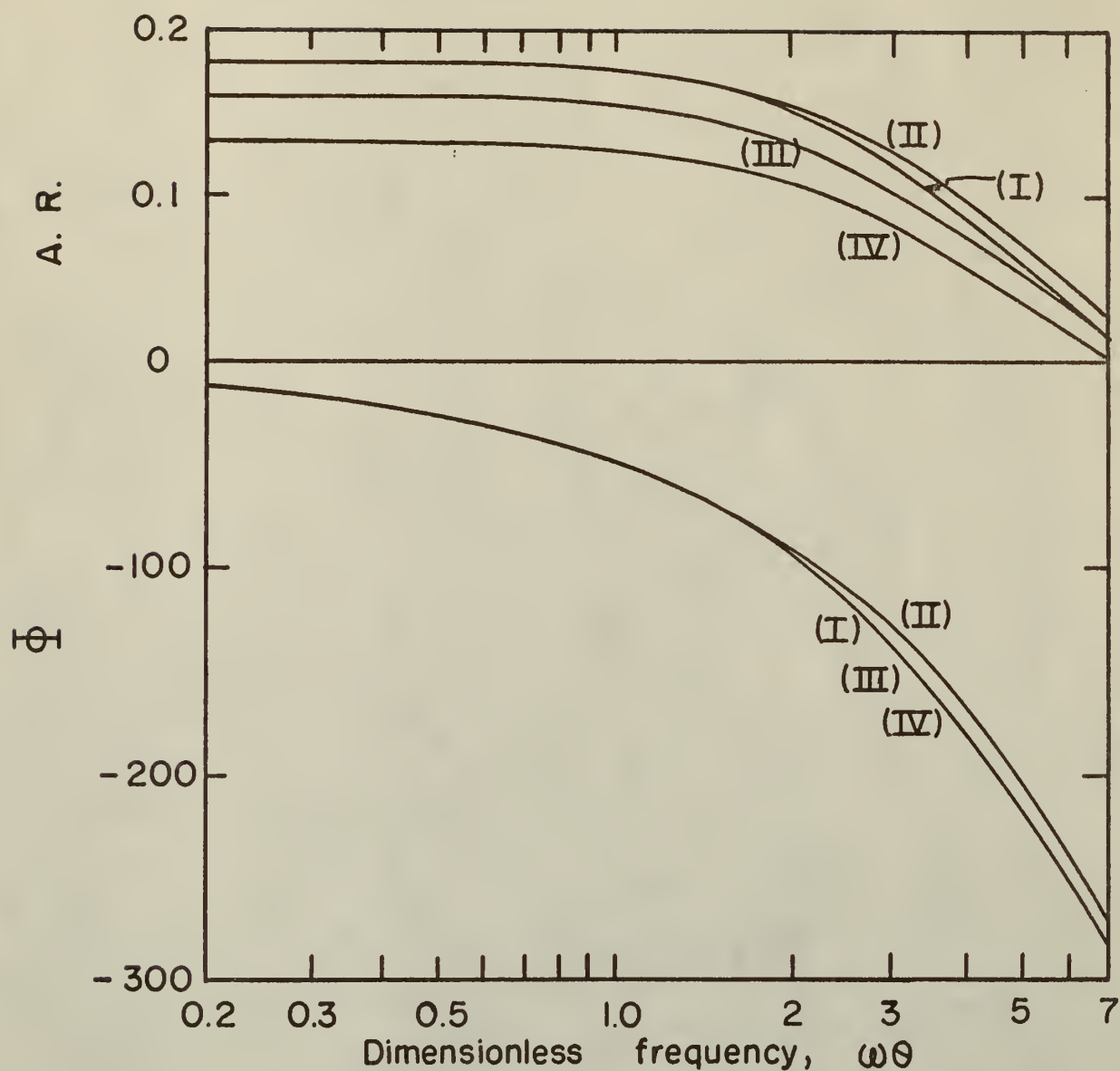


Fig. 32. Comparison of frequency response diagrams of reacting system ($R=2$) based on different boundary conditions at low degree of dispersion, $M=7$.

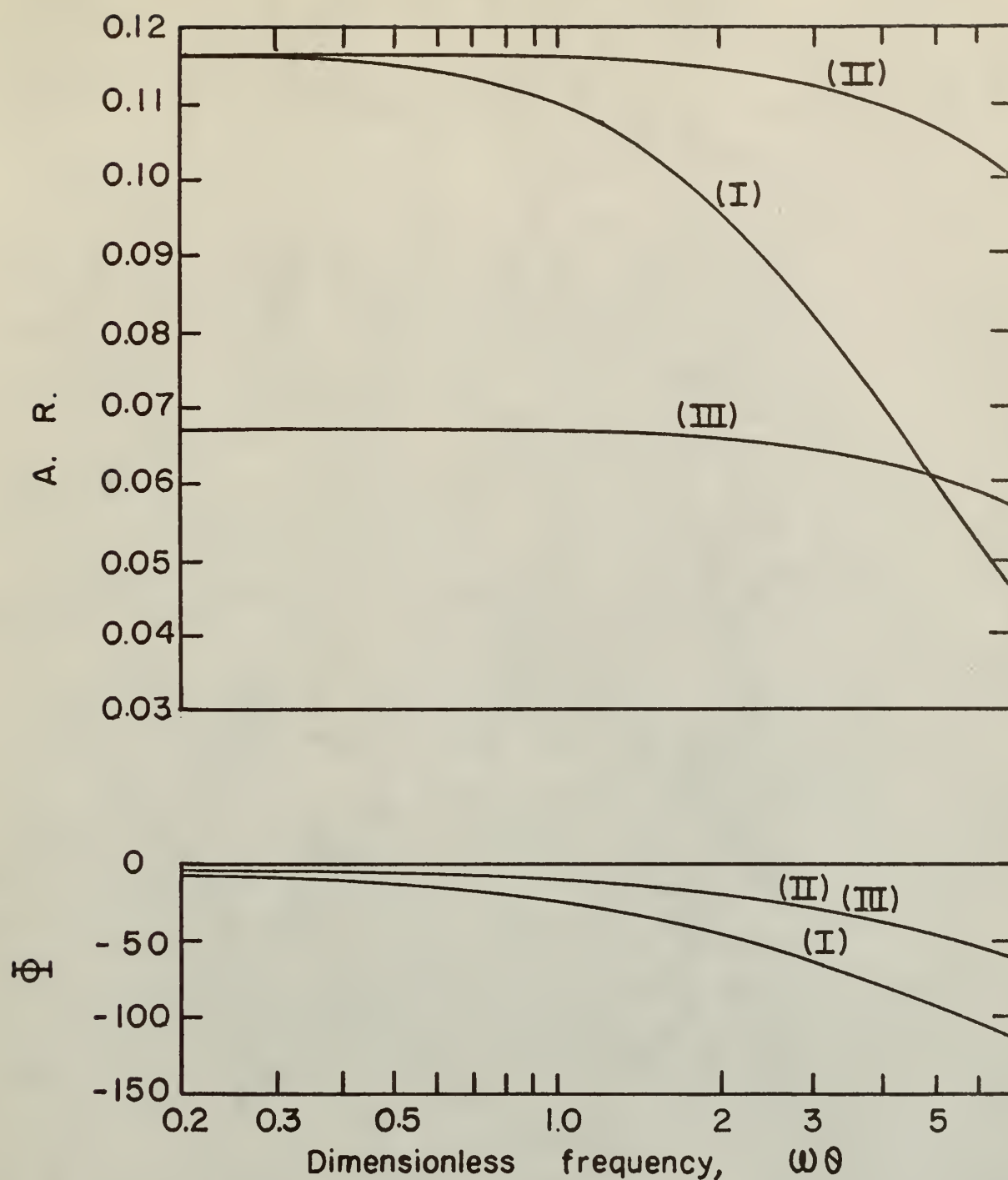


Fig. 33. Comparison of frequency response diagrams of reacting system ($R=10$) based on different boundary conditions at high degree of dispersion, $M=\frac{1}{2}$.

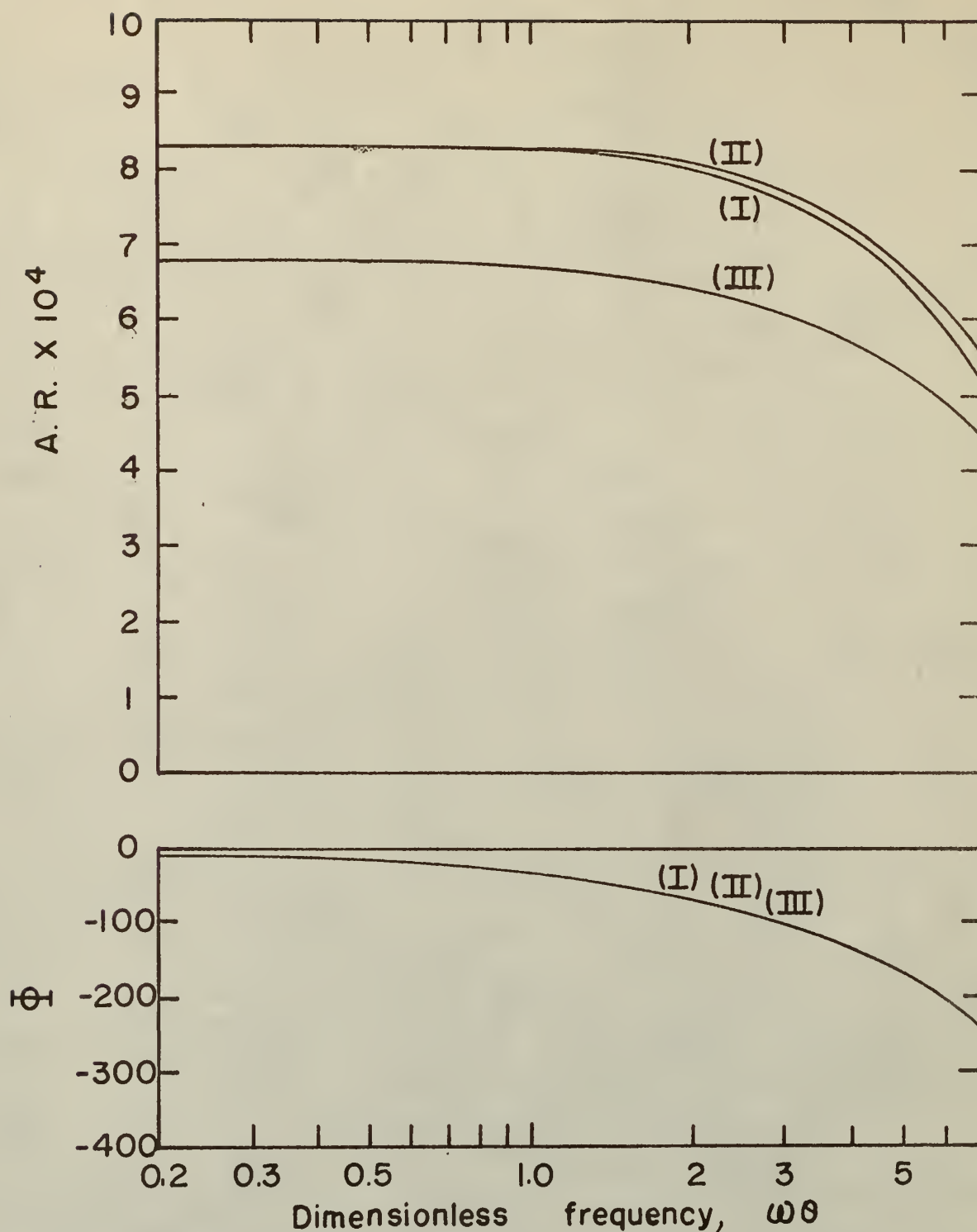


Fig. 34. Comparison of frequency response diagrams of reacting system ($R=10$) based on different boundary conditions at low degree of dispersion, $M=10$.

the corresponding curves for the non-reacting systems -- the greater the value of R , the greater the effect of the boundary conditions. It is also seen from these figures that, at any given relative reaction rate characterized by R , the effect of different sets of boundary conditions becomes smaller as the degree of mixing decreases.

An illustration of the determination of the mean dispersion coefficient from a frequency response experiment for non-reacting systems by comparing available experimental data (7) (10) (22) with calculated values is presented in Figs. 35 through 37. These figures show that the use of approximate boundary conditions can give rise to an appreciable error in determining the degree of dispersion, especially at low values of M . As has been discussed in the previous section, the use of B.C.(III) is quite desirable under some circumstances due to the simplicity of equation (33). For the preliminary design of a tubular flow system and the synthesis of a complex control system in which the use of a complicated performance function is involved, the use of approximate boundary conditions should be avoided.

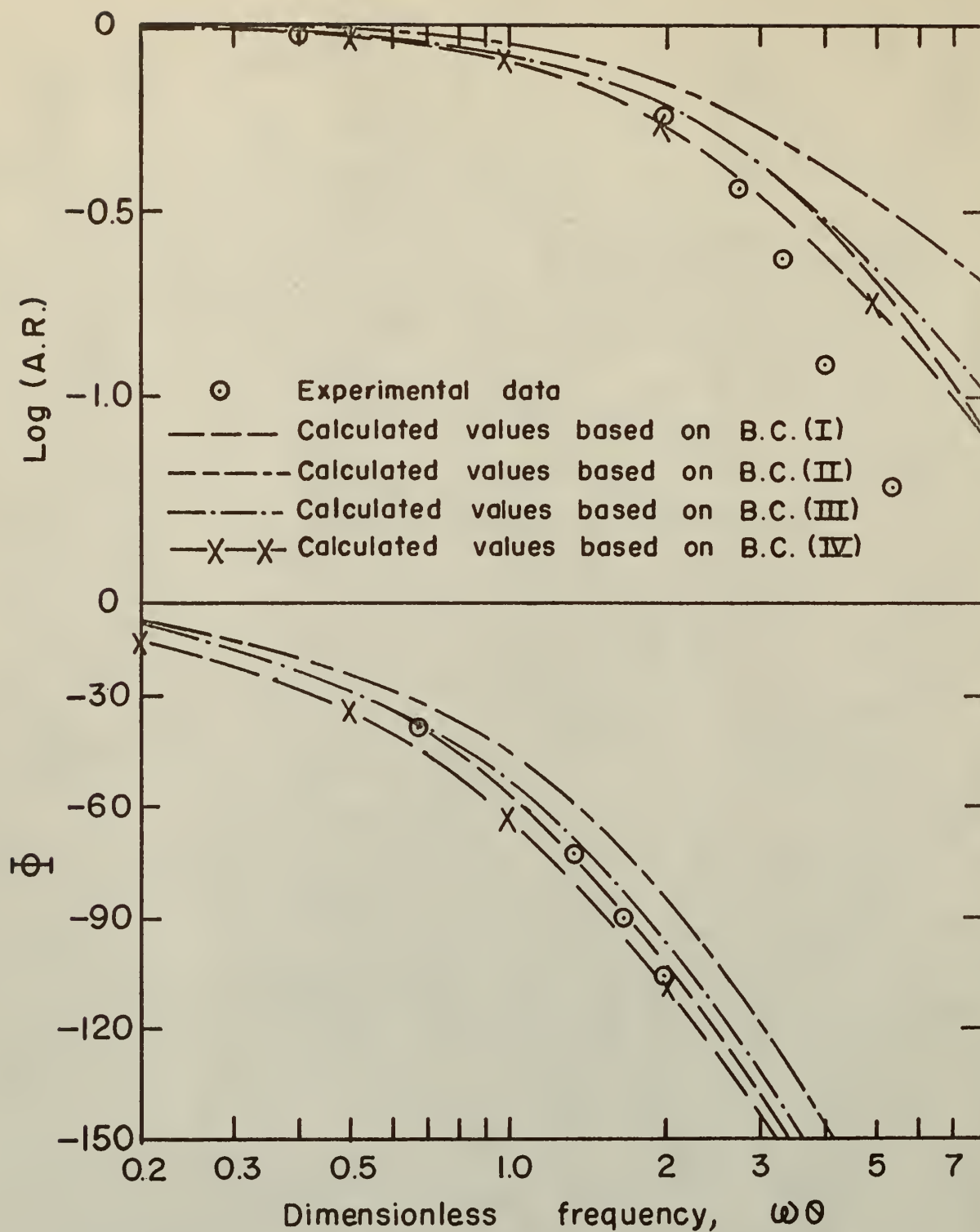


Fig. 35. Comparison of experimental and calculated frequency response diagrams.

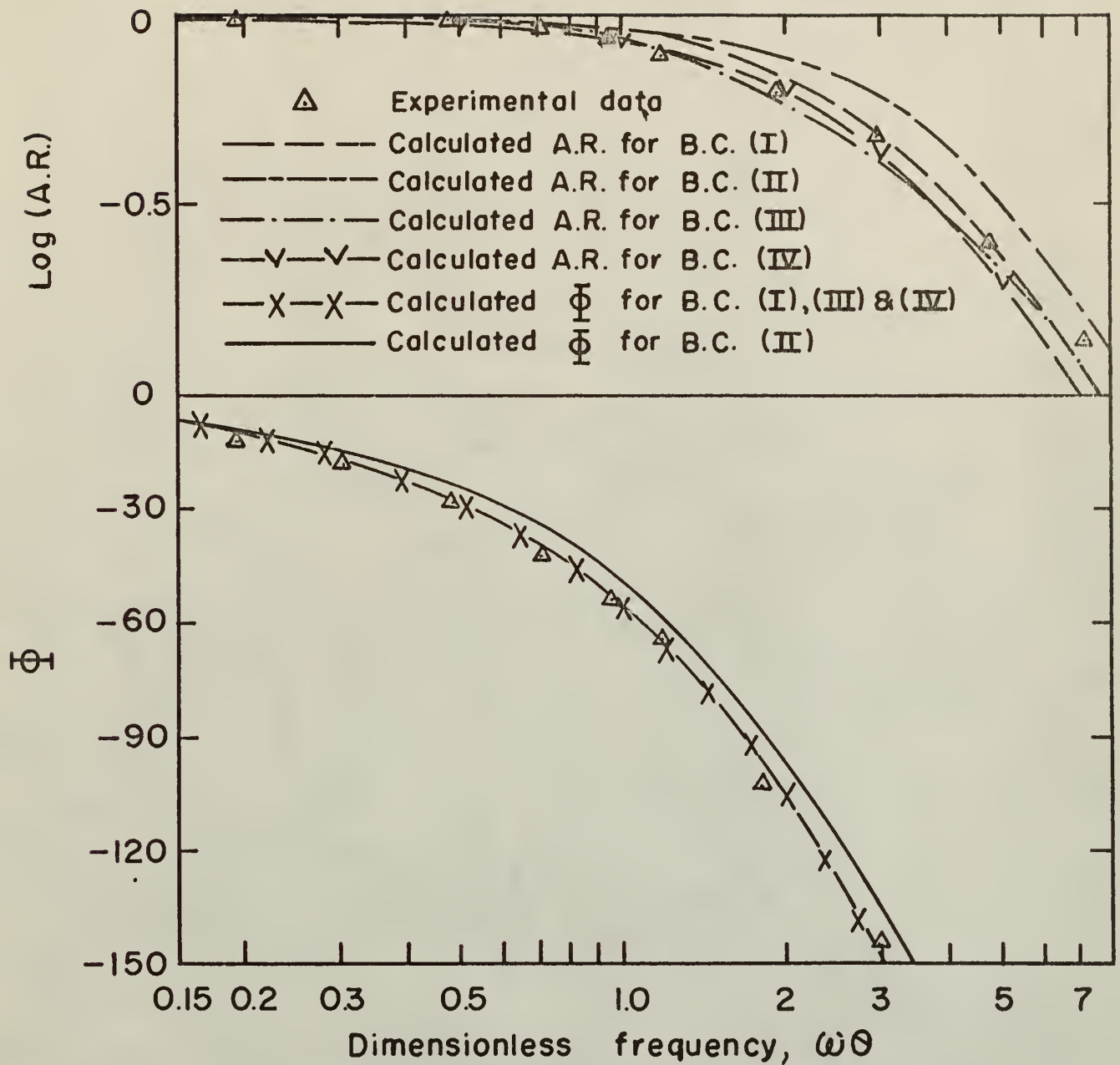


Fig. 36. Comparison of experimental and calculated frequency response diagrams.

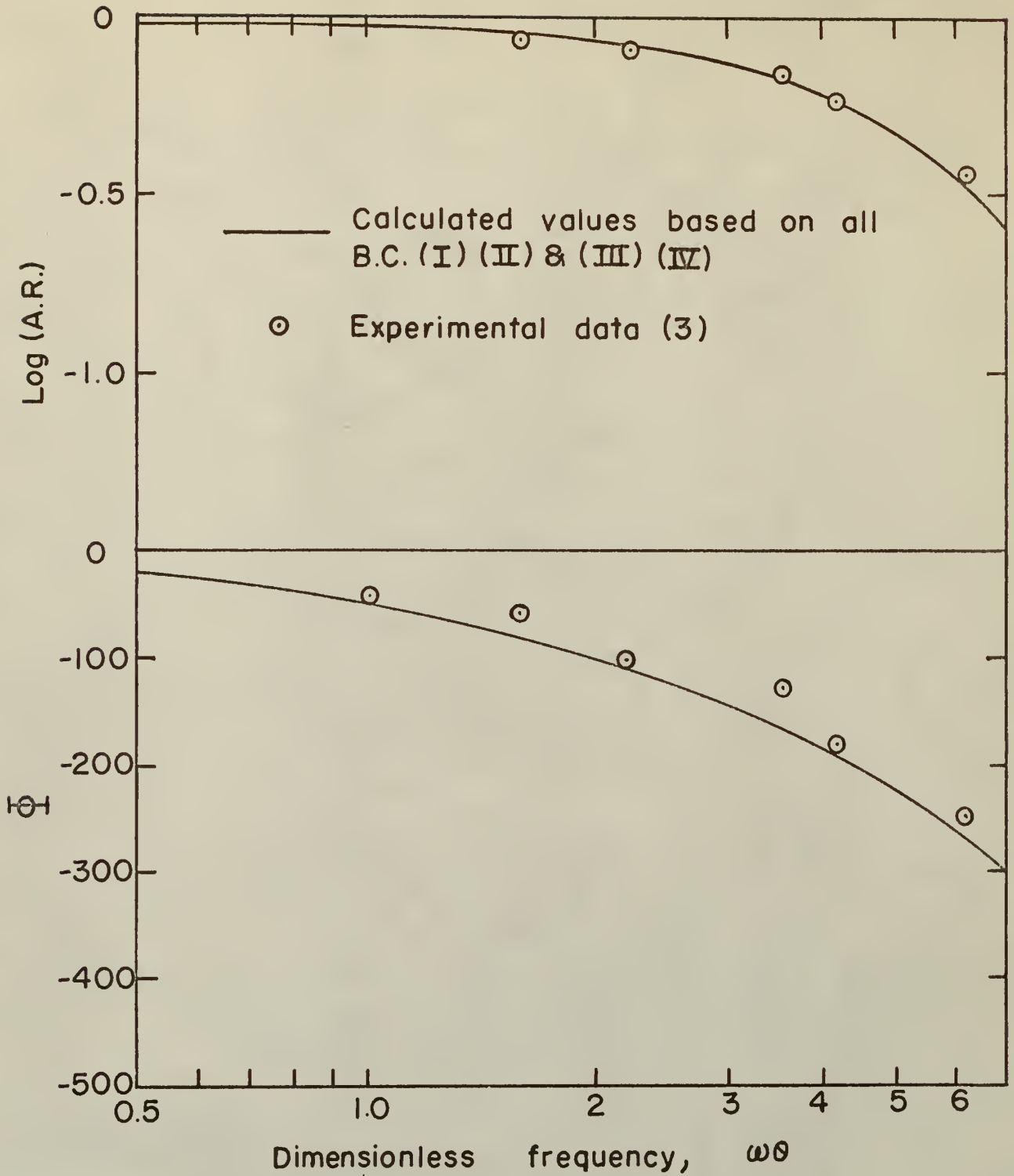


Fig. 37. Comparison of experimental and calculated frequency response diagrams.

CONCLUSION

The general extent of axial dispersion and its effect on the performance of isothermal tubular flow reactors can be characterized by a diffusion-type model with appropriate boundary conditions. Among the four boundary conditions considered, only B.C.(I) yields conversions and concentration profiles consistent with the theoretical limits of both a uniformly stirred tank reactor and a plug flow type reactor, and for this reason, the use of B.C.(II), (III) and (IV) should be avoided in and around the domain of inconsistent conversions based on a one-dimensional transport equation.

There are no data available on transient response experiments accompanied by first-order chemical reaction. If the data were available, the calculated response curves based on first-order chemical reaction could be used to determine the dispersion coefficient by the same method presented in this research for the non-reacting case.

The response curves presented in this research for various types of transient input should be useful for practicing process control engineers and for those who wish to extract the transport parameters from dynamic response experiments.

ACKNOWLEDGMENT

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NOMENCLATURE

a	= angle defined as $\frac{1}{2} \tan^{-1} \left[\frac{2(\omega \theta)}{2R + M} \right]$, degree.
A	= a constant in equation (27) which is numerically equal to the steady-state conversion, dimensionless.
(A.R.)	= amplitude ratio, dimensionless.
b	= $\left[\left(1 + \frac{2R}{M}\right)^2 + \frac{4(\omega \theta)^2}{M^2} \right]^{\frac{1}{4}}$
B	= a constant in equation (28) which is numerically equal to the steady-state conversion, dimensionless.
C	= concentration of reactant, m/l.
$C_1(t)$	= inlet concentration of reactant, m/l.
$C_2(t)$	= exit concentration of reactant, m/l.
D	= axial dispersion coefficient, l^2/t .
\bar{D}	= mean value of D, l^2/t .
f_i	= conversion calculated based on B.C.(i), dimensionless.
$f_{M=0}$	= conversion calculated based on a stirred tank reactor, dimensionless.
$f_{M=\infty}$	= conversion calculated based on a plug flow reactor, dimensionless.
j	= $\sqrt{-1}$
k	= reaction rate constant, t^{-1} .
L	= length of reactor, l.
M	= $\bar{u}L/2\bar{D}$, 1/2 axial Peclet number, dimensionless.
n	= order of reaction.
P.F.	= performance function defined as the ratio of Fourier transform of output to input.

Q	=	amount of tracer injected in impulse function, m.
R	=	kL/\bar{u} , reaction number, dimensionless.
$r(C)$	=	rate of reaction.
s	=	function of Laplace transform.
t	=	time.
\bar{u}	=	average flow velocity of reacting mixture, l/t.
v	=	volumetric flow rate, l ³ /t.
V	=	volume of reactor, l ³ .
$W(s)$	=	transfer function defined as the Laplace transform of output to input.
x	=	$Mb \cos a$
y	=	$Mb \sin a$
z	=	co-ordinate in the direction of flow, l.

Greek Letters

$$\alpha_1 = \sinh x \cos y - \frac{(\omega \theta)}{M} \cosh x \sin y + b(\cos a \cosh x \cos y - \sin a \sinh x \sin y).$$

$$\alpha_2 = \left(\frac{\omega \theta}{M}\right) \sinh x \cos y + \cosh x \sin y + b(\cos a \sinh x \sin y + \sin a \cosh x \cos y).$$

$$\alpha_3 = \sinh x \cos y + b(\cos a \cosh x \cos y - \sin a \sinh x \sin y).$$

$$\alpha_4 = \cosh x \sin y + b(\sin a \cosh x \cos y + \cos a \sinh x \sin y).$$

$$\beta = (1 + 2R/M)^{\frac{1}{2}}.$$

$$\beta' = [1 + (2\theta/M)(s + k)]^{\frac{1}{2}}$$

$$\beta'' = [1 + (2\theta/M)(j\omega + k)]^{\frac{1}{2}}$$

$$\delta_n = \text{roots of } \cot \delta = \frac{1}{2} \left(\frac{\delta}{M} - \frac{M}{\delta} \right).$$

$$\delta'_n = \text{roots of } \tan \delta' = - \frac{\delta'}{M}.$$

$$\phi = \text{phase shift, degree.}$$

$$\eta = z/L, \text{ dimensionless distance from the inlet of reactor.}$$

$$\gamma = C/C_1, \text{ dimensionless concentration.}$$

$$\tau = \text{the dimensionless time variable, } t/\theta.$$

$$\tau_0 = \text{thickness of impulse in time domain, dimensionless.}$$

$$\theta = L/\bar{u}, t.$$

$$\omega = \text{angular frequency, rad./t.}$$

APPENDIX

APPENDIX I

Derivations of Steady-state Solutions

The solutions of equation (8) which satisfy the boundary conditions (I), (II), (III) and (IV) are obtained by means of Laplace transformations.

Laplace transform of equation (8) with respect to η is:

$$\mathcal{Y}(s) = \frac{s \mathcal{Y}(0^+) + \mathcal{Y}'(0^+) - 2M \mathcal{Y}'(0^+)}{s^2 - 2Ms - 2MR} \quad (48)$$

For the solution using B.C.(I), the Laplace transform of equation (3) is substituted in equation (48) to obtain:

$$\mathcal{Y}(s) = \frac{s \mathcal{Y}(0^+) - 2M}{s^2 - 2Ms - 2MR} \quad (49)$$

By means of partial fractionation, equation (49) is separated into two parts:

$$\mathcal{Y}(s) = \frac{(1+\beta) \mathcal{Y}(0^+) - 2}{2\beta[s - M(1+\beta)]} - \frac{(1-\beta) \mathcal{Y}(0^+) - 2}{2\beta[s - M(1-\beta)]} \quad (50)$$

The inverse transform of equation (50) gives:

$$\begin{aligned} \mathcal{Y}(\eta) = & \frac{(1+\beta) \mathcal{Y}(0^+) - 2}{2\beta} \exp [M(1+\beta)\eta] \\ & - \frac{(1-\beta) \mathcal{Y}(0^+) - 2}{2} \exp [M(1-\beta)\eta] \end{aligned} \quad (51)$$

According to boundary condition represented by equation (4), differentiation of equation (51) and setting it equal to zero

gives:

$$\gamma(0^+) = \frac{2\{(1+\beta)\exp[M(1+\beta)] - (1-\beta)\exp[M(1-\beta)]\}}{(1+\beta)^2\exp[M(1+\beta)] - (1-\beta)^2\exp[M(1-\beta)]} \quad (52)$$

Substitution of equation (52) into (51) and simplification gives equation (9) of text.

For the solution using B.C.(II), the Laplace transform of equation (5) is substituted into equation (48) and simplified to:

$$\gamma(s) = \frac{(s - 2M) + \gamma'(0^+)}{s^2 - 2Ms - 2MR} \quad (53)$$

By means of partial fractionation, equation (53) is separated into two parts as:

$$\gamma(s) = \frac{-M(1-\beta) + \gamma'(0^+)}{2M\beta[s-M(1+\beta)]} - \frac{-M(1+\beta) + \gamma'(0^+)}{2M\beta[s-M(1-\beta)]} \quad (54)$$

The inverse transform of equation (54) yields:

$$\begin{aligned} \gamma(\eta) = & \frac{-M(1-\beta) + \gamma'(0^+)}{2M\beta} \exp[M(1+\beta)\eta] \\ & - \frac{-M(1+\beta) + \gamma'(0^+)}{2M\beta} \exp[M(1-\beta)\eta] \end{aligned} \quad (55)$$

According to boundary condition represented by equation (4), equation (55) is differentiated in order to solve for $\gamma'(0^+)$.

It gives:

$$\gamma'(0^+) = M(1-\beta^2) \left\{ \frac{\exp[M(1+\beta)] - \exp[M(1-\beta)]}{(1+\beta)\exp[M(1+\beta)] - (1-\beta)\exp[M(1-\beta)]} \right\} \quad (56)$$

Substitution of equation (56) into (55) and simplification gives equation (10) of text.

For the solution using B.C.(III), substitution of equation (5) into equation (48) gives equation (55). As η approaches infinity in equation (55), the second term in the right-hand side of equation (55) vanishes as:

$$\frac{M(\beta - 1) + \gamma(0^+)}{2M\beta} = 0 \quad (57)$$

Solving for $\gamma(0^+)$ in equation (57) and substituting it into equation (55), equation (11) of text is obtained.

For the solution using B.C.(IV), substitution of equation (3) into equation (48) gives equation (51). As η approaches infinity in equation (51), the first term in the right-hand side of equation (51) vanishes, hence

$$0 = \frac{(1+\beta) \gamma(0^+) - 2}{2\beta} \quad (58)$$

Solving for $\gamma(0^+)$ in equation (58) and substituting it into equation (51) yields equation (12) of text.

APPENDIX II

Derivations of the Transfer Functions, $W(s)$

The Laplace transform of equation (2) with respect to t at zero initial concentration is:

$$sC(z,s) = \bar{D} \frac{\partial^2 C(z,s)}{\partial z^2} - \bar{u} \frac{\partial C(z,s)}{\partial z} - kC(z,s) \quad (59)$$

Equation (59) is in the form of a second-order differential equation whose general solution is:

$$C(z,s) = A_1 e^{\frac{\bar{u}z}{2\bar{D}} (1+\beta)} + A_2 e^{\frac{\bar{u}z}{2\bar{D}} (1-\beta)} \quad (60)$$

where A_1, A_2 = arbitrary constants.

The transfer functions, defined as the ratio of the Laplace transforms of output to input, then become

$$W(s) = \frac{C(L,s)}{C_1(0,s)} \quad (61)$$

In order to solve for the arbitrary constants in equation (60), the boundary conditions tabulated in Table 2 are transformed as:

$$\text{B.C. (I)} \quad \left\{ \begin{array}{l} \bar{u}C(0,s) = \bar{u}C(0^+,s) - \bar{D} \frac{dC(0^+,s)}{dz} \\ \frac{\partial C(L,s)}{\partial z} = 0 \end{array} \right. \quad \begin{array}{l} (62a) \\ (62b) \end{array}$$

$$\text{B.C. (II)} \quad \left\{ \begin{array}{l} \bar{u}C(0,s) = \bar{u}C(0^+,s) \\ \frac{\partial C(L,s)}{\partial z} = 0 \end{array} \right. \quad \begin{array}{l} (63a) \\ (63b) \end{array}$$

$$\text{B.C. (III)} \quad \left\{ \begin{array}{l} \bar{u}C(0,s) = \bar{u}C(0^+,s) \\ \lim_{z \rightarrow \infty} C(z,s) = 0 \end{array} \right. \quad \begin{array}{l} (64a) \\ (64b) \end{array}$$

$$\text{B.C. (IV)} \quad \left\{ \begin{array}{l} \bar{u}C(0,s) = \bar{u}C(0^+,s) - \bar{D} \frac{\partial C(0^+,s)}{\partial z} \\ \lim_{z \rightarrow \infty} C(t,s) = 0 \end{array} \right. \quad \begin{array}{l} (65a) \\ (65b) \end{array}$$

For the transfer function corresponding to B.C.(I), equation (60) is differentiated with respect to z :

$$\frac{\partial C(z,s)}{\partial z} = \frac{A_1 \bar{u}}{2D} e^{\frac{\bar{u}z}{2D}(1+\beta')} + \frac{A_2 \bar{u}}{2D} e^{\frac{\bar{u}z}{2D}(1-\beta')} \quad (66)$$

where $\beta' = \left[1 + \left(\frac{2\theta}{M} \right) (s + k) \right]^{\frac{1}{2}}$

According to equation (62b), equation (66) is set to zero to solve for one constant in terms of another as:

$$A_1 = - \frac{A_2(1-\beta') e^{M(1-\beta')}}{(1+\beta') e^{M(1+\beta')}} \quad (67)$$

where $M = \frac{\bar{u}L}{2D}$

Using equation (62a), the Laplace transform of the input is obtained as:

$$C_1(0,s) = \frac{A_1}{2}(1-\beta') + \frac{A_2}{2}(1+\beta') \quad (68)$$

Substitution of equation (67) into equation (68) gives:

$$C_1(0,s) = \frac{A_2}{2} \left[\frac{(1+\beta')^2 e^{M(1+\beta')} - (1-\beta')^2 e^{M(1-\beta')}}{(1+\beta') e^{M(1+\beta')}} \right] \quad (69)$$

The Laplace transform of output is:

$$C(L,s) = A_1 e^{M(1+\beta')} + A_2 e^{M(1-\beta')} \quad (70)$$

Substitution of equation (67) into equation (70) gives:

$$C(L,s) = A_2 \left\{ \frac{+ 2 \beta' e^{M(1-\beta')}}{1 + \beta'} \right\} \quad (71)$$

According to equation (61), the transfer function is the ratio of equations (71) to (67), which upon simplification gives equation (19) of text.

For the transfer function corresponding to B.C.(II), equation (63a) is substituted into equation (60) to obtain:

$$C_1(0,s) = A_1 + A_2 \quad (72)$$

Since the second condition of B.C.(II) is the same as that of B.C.(I), equation (67) still holds for B.C.(II), and substitution of equation (67) into (72) gives:

$$C_1(0,s) = A_2 \left[\frac{(1+\beta') e^{M(1+\beta')} - (1-\beta') e^{M(1-\beta')}}{(1+\beta') e^{M(1+\beta')}} \right] \quad (73)$$

From equation (61), the transfer function corresponding to B.C.(II) is the ratio of equations (71) to (73), which upon simplification gives equation (20) of text.

For the transfer function corresponding to B.C.(III), equation (64b) is substituted to equation (60) to obtain:

$$\lim_{z \rightarrow \infty} C(z,s) = 0 = A_1(\infty) + A_2(0) \quad (74a)$$

So A_1 must be zero in order to satisfy equation (64b).

$$A_1 = 0 \quad (74b)$$

Since the first set of B.C.(III) is the same as that of B.C.(II), equation (72) still holds for B.C.(III). Substitution of equation (74) into (72) becomes:

$$C_1(0,s) = A_2 \quad (75)$$

The Laplace transform of the output is:

$$C(L,s) = A_2 e^{M(1-\beta')} \quad (76)$$

The transfer function corresponding to B.C.(III) is the ratio of equations (76) to (75), which upon simplification, gives equation (21) of the text.

Since B.C.(IV) is the combination of the first set of B.C.(I) and the second set of B.C.(II), equations (68) and (74) still hold for this case. Substitution of equation (74) into (68) gives:

$$C_1(0,s) = \frac{A_2(1+\beta')}{2} \quad (77)$$

The transfer function corresponding to B.C.(IV) is the ratio of equations (76) to (77), which upon simplification, gives equation (22) of text.

APPENDIX III

Systematic Derivations of The Response Functions of The Exit Concentrations Against Impulse Function and Step Inputs.

For the response functions of the exit concentration against impulse, it is known by L'Hospital rule that

$$C(L,s) = C_1 W(s) \quad (78)$$

For the solution corresponding to B.C.(I), equation (19) is substituted into equation (78) to obtain:

$$C(L,s) = C_1 \frac{4\beta'}{(1+\beta')^2 e^{-M(1-\beta')} - (1-\beta')^2 e^{-M(1+\beta')}} \quad (79)$$

The inverse transformation of equation (79) is carried out by means of the Heaviside partial fraction theorem. Since equation (79) is in the form of:

$$C(L,s) = \frac{q_1(s)}{p_1(s)} \quad (80)$$

where $q_1(s) = 4C_1\beta'$

$$p_1(s) = (1+\beta')^2 e^{-M(1-\beta')} - (1-\beta')^2 e^{-M(1+\beta')}$$

the inverse transform of equation (80) gives:

$$C(L,t) = \sum_{k=1}^{\infty} \frac{q_1(s_k)}{p_1'(s_k)} e^{s_k t} \quad (81)$$

where s_k is the root of $p_1(s) = 0$

Solving for the roots of $p_1(s) = 0$, one obtains:

$$s_k = - \frac{\delta_n^2 + M^2}{2M\theta} - k \quad (82)$$

where $\delta_n = \text{root of } \cot \delta = \frac{1}{2} \left(\frac{\delta}{M} - \frac{M}{\delta} \right)$

Differentiation of $p_1(s)$ with respect to s gives:

$$p'_1(s) = - \frac{\theta}{M\alpha} \left[M \cos M\alpha + 2 \cos M\alpha - 2M\alpha \sin M\alpha \right. \\ \left. - 2\alpha \sin M\alpha - M\alpha^2 \cos M\alpha \right] \quad (83)$$

where $\alpha = j\beta'$

Substitution of equation (82) into the expressions of $p'_1(s)$ and $q_1(s)$ gives respectively:

$$q(s_k) = C_1 \frac{2 \delta_{ne}^M}{M} \quad (84)$$

$$p'_1(s_k) = - \frac{\theta}{\delta_n} \left[M \cos \delta_n + 2 \cos \delta_n - 2 \delta_n \sin \delta_n \right. \\ \left. - \frac{2 \delta_n}{M} \sin \delta_n - \frac{\delta_n^2}{M} \cos \delta_n \right] \quad (85)$$

Substitution of equations (84) and (85) into equation (81) gives:

$$\frac{C(L,t)}{C_1} = \sum_{n=1}^{\infty} \frac{1}{M \cos \delta_n + 2 \cos \delta_{n-2} \delta_n \sin \delta_n - \frac{2 \delta_n}{M} \sin \delta_n - \frac{\delta_n^2}{M} \cos \delta_n} \quad (86)$$

$$\times \left(\frac{2 \delta_n e^M}{M} \right) \left(-\frac{\delta_n}{\theta} \right) \exp \left[- \left(\frac{\delta_n^2 + M^2}{2M\theta} \right) t - kt \right]$$

Multiplication of both numerator and denominator on the right-hand side of equation (86) by $(M \sin \delta_n + \delta_n \cos \delta_n)$ gives:

$$\frac{C(L,t)}{C_1} = \sum_{n=1}^{\infty} \left(-\frac{2 \delta_n^2}{M} \right) \left(\frac{M \sin \delta_n + \delta_n \cos \delta_n}{K} \right) \quad (87)$$

$$\times \exp \left[- \left(\frac{\delta_n^2 + M^2}{2M\theta} \right) t - kt \right]$$

$$\text{where } K = (M \cos \delta_n + 2 \cos \delta_{n-2} \delta_n \sin \delta_n - \frac{\delta_n}{M} \sin \delta_n - \frac{\delta_n^2}{M} \cos \delta_n)$$

$$\times (M \sin \delta_n + \delta_n \cos \delta_n)$$

Simplification of K gives:

$$K = -\frac{\delta_n}{M} (M^2 + \delta_n^2 + 2M) \quad (88)$$

Substitution of equation (88) into (87) and simplification, gives:

$$\frac{C_2(L,t) \theta}{C_1} = \sum_{n=1}^{\infty} \frac{2 \delta_n (M \sin \delta_n + \delta_n \cos \delta_n)}{(M^2 + \delta_n^2 + 2M)} \quad (89)$$

$$\times \exp \left[M - \left(\frac{\delta_n^2 + M^2}{2M} + R \right) \tau \right]$$

where $\tau = \frac{t}{\theta}$

The left-hand side of equation (89) is transformed in terms of the amount of tracer injected, Q , and the volume of reactor, V , by the following relationship:

$$Q = \bar{u} \int_0^\lambda C(\lambda) d\lambda = C_1 \pi d_r^2 u \quad (90)$$

$$V = \pi d_r^2 L \quad (91)$$

where d_r = diameter of reactor.

Hence, substitution of equations (90) and (91) into the left-hand side of equation (89) gives the following relationship:

$$\frac{C_2(L,t)}{C_1} = \frac{C_2(L,t)}{C_1} \times \frac{L}{u} = \frac{C_2 \pi d_r^2 L}{C_1 \pi d_r^2 u} = \frac{C_2 V}{Q} \quad (92)$$

Using equation (92) in place of equation (89) gives equation (23) of text.

For the solution corresponding to B.C.(II), equation (20) is substituted into equation (78) to obtain:

$$C(L,s) = C_1 \frac{2\beta'}{(1+\beta') e^{-M(1+\beta')} - (1-\beta') e^{-M(1-\beta')}} \quad (93)$$

Equation (93) is in the same form as equation (80), except $p_1(s)$ and $q_1(s)$ are replaced by the following two expressions:

$$q_2(s) = 2C_1 \beta'$$

$$p_2(s) = (1+\beta') e^{-M(1+\beta')} - (1-\beta') e^{-M(1-\beta')}$$

Solving for the root of $p_2(s) = 0$, one obtains:

$$s_k = - \frac{\delta_n^2 + M^2}{2M\theta} - k \quad (94)$$

where $\delta'_n = \text{root of } \tan \delta' = - \frac{\delta'}{M}$

Differentiation of p_2 with respect to s gives:

$$p'_2(s) = \left(- \frac{\theta}{M\alpha} \right) (M \cos M\alpha + \cos M\alpha - M\alpha \sin M\alpha) \quad (95)$$

Substitution of equation (94) into the expressions of $p'_2(s)$ and $q_2(s)$ gives respectively:

$$q_2(s_k) = \frac{\delta'_n}{M} e^M \quad (96)$$

$$p'_2(s_k) = - \left(\frac{\theta}{\delta'_n} \right) (M \cos \delta'_n + \cos \delta'_n - \delta'_n \sin \delta'_n) \quad (97)$$

Substitution of equations (96) and (97) into equation (81) gives:

$$\frac{C_2(L,t)}{C_1} = \sum_{n=1}^{\infty} \left(- \frac{\delta_n^2}{M} \right) \frac{1}{K_2} \exp \left[M - \left(\frac{\delta_n^2 + M^2}{2M\theta} \right) t - kt \right] \quad (98)$$

where $K_2 = M \cos \delta'_n + \cos \delta'_n - \delta'_n \sin \delta'_n$

Simplification of K_2 gives:

$$K_2 = \frac{M^2 + M + \delta_n^2}{M} \cos \delta'_n \quad (99)$$

Substitution of equation (99) into (98), upon simplification, gives:

$$\frac{C_2(L,t)\theta}{C_1} = \sum_{n=1}^{\infty} \frac{-\delta_n^2}{(M^2 + M + \delta_n'^2) \cos \delta'_n} \times \exp \left[M - \left(\frac{\delta_n^2}{2M} + R \right) \tau \right] \quad (100)$$

Substitution of equations (90) and (91) into the left-hand side of equation (100) gives equation (24) of text.

For the solution corresponding to B.C.(III), equation (21) is substituted into equation (78) to obtain:

$$C(L,s) = C_1 \exp \left[M(1-\beta') \right] \quad (101)$$

If one lets $g_1 = L/\sqrt{D}$ and $g_2 = \frac{\bar{u}^2}{4D} + k$, equation (101) is simplified to:

$$C(L,s) = C_1 e^M e^{-g_1 \sqrt{g_2 + s}} \quad (102)$$

Equation (102) is transformed to the following expression by the rule of "shift of original function" as:

$$\frac{C(L,t)}{C_1} = e^{M-g_2 t} \mathcal{L}^{-1} \left\{ e^{-g_1 \sqrt{s}} \right\} \quad (103)$$

The inverse transformations of equation (103) is carried out by means of tables of integral transforms (5):

$$\mathcal{L}^{-1} \left\{ e^{-g_1 \sqrt{s}} \right\} = \frac{g_1}{2\sqrt{\pi} t^{3/2}} \exp \left(-\frac{g_1^2}{4t} \right) \quad (104)$$

Substitution of g_1 and g_2 into equation (104) gives:

$$\frac{c_2 \theta}{c_1} = \sqrt{\frac{M}{2\pi}} (\mathcal{L})^{-1} \frac{3}{2} \exp \left[M - \frac{1}{2} \left(M \mathcal{L} + \frac{M}{\mathcal{L}} \right) - R\mathcal{L} \right] \quad (105)$$

Use of equations (90) and (91) in place of the left-hand side of equation (105) gives equation (25) of text.

For the solution corresponding to B.C.(IV), equation (22) is substituted into equation (78) to obtain:

$$C(L,s) = c_1 \frac{2 \exp [M(1-\beta')]}{(1 + \beta')} \quad (106)$$

If one lets $g_3 = \frac{L}{\sqrt{D}}$, $g_4 = \frac{\bar{u}^2}{4D} + k$, and $g_5 = \frac{M}{g_3}$,

equation (106) becomes:

$$\frac{C(L,s)}{c_1} = 2g_5 e^M \left\{ \frac{e^{-g_3 \sqrt{g_4 + s}}}{g_5 + \sqrt{g_4 + s}} \right\} \quad (107)$$

The inverse transformations of equation (107) by means of "shifts of original function" and the tables of integral transforms (5) gives:

$$\frac{C(L,t)}{C_1} = 2g_5 e^{M-g_4 t} \frac{1}{\sqrt{\pi t}} \exp\left(-\frac{g_3^2}{4t}\right) - g_5 \exp(g_3 g_5 + g_5^2 t) \\ \times \operatorname{Erfc} \frac{g_3}{2\sqrt{t}} + g_5 \sqrt{t} \quad (108)$$

Substitution of g_1 , g_2 and g_3 into equation (108), upon simplification, gives equation (26) of text.

The response functions of the exit concentration against step input are known to be:

$$C(L,s) = \frac{C_1 W(s)}{s} \quad (109)$$

Since the transfer functions, $W(s)$, corresponding to each of four boundary conditions and their inverse transforms are known by the derivations of response functions against impulse, the following relationships can be expressed from equation (109) by the method of convolution:

$$\frac{C(L,t)}{C_1} = \int_0^t \sum_{n=1}^{\infty} \frac{2 \delta_n (M \sin \delta_n + \delta_n \cos \delta_n)}{(M^2 + 2M + \delta_n^2)} \\ \times \exp \left[M - \left(\frac{\delta_n^2 + M^2}{2M} + R \right) \lambda \right] d\lambda \quad (110)$$

$$\frac{C(L,t)}{C_1} = \int_0^t \sum_{n=1}^{\infty} \frac{-\delta_n'^2}{(M^2 + M + \delta_n'^2) \cos \delta_n'} \\ \times \exp \left[M - \left(\frac{\delta_n'^2 + M^2}{2M} + R \right) \lambda \right] d\lambda \quad (111)$$

$$\frac{C(L,t)}{C_1} = \int_0^{\infty} \sqrt{\frac{M}{2\pi}} (\lambda)^{-\frac{3}{2}} \times \exp \left[M - \frac{1}{2} \left(M\lambda + \frac{M}{\lambda} + 2\lambda R \right) \right] d\lambda \quad (112)$$

$$\frac{C(L,t)}{C_1} = \int_0^{\infty} 2 \left\{ \exp \left[M - (M + 2R) \frac{\lambda}{2} \right] \operatorname{Erfc} \sqrt{\frac{M}{2\lambda}} - \exp(2M - R\lambda) \operatorname{Erfc} \left(\sqrt{\frac{M}{2\lambda}} + \sqrt{\frac{M\lambda}{2}} \right) \right\} d\lambda \quad (113)$$

The integration of equations (110) and (111) are straight forward, and their results can be obtained easily as shown by equations (27) and (28) of the text respectively. The integrations of equations (112) and (113) are difficult to carry out by the method of convolution because they contain the error functions. Heaviside's partial fraction theorem was found to be more convenient for the inverse transform of these two equations.

For the solution corresponding to B.C.(III), equation (21) is substituted into equation (109) to obtain:

$$C(L,s) = \frac{C_1 e^{M(1-\beta')}}{s} \quad (114)$$

If one lets $g_6 = \frac{L}{\sqrt{D}}$ and $g_7 = \frac{\bar{u}}{4\bar{D}} + k$, equation (114) is simplified to:

$$\frac{C(L,s)}{C_1} = \frac{e^{-g_6 \sqrt{g_7 + s}}}{s} \cdot e^{M} \quad (115)$$

Equation (115) is transformed to the following expression by the rule of "shift of original function" as:

$$\frac{C(L,t)}{C_1} = e^{M-g_7 t} \mathcal{L}^{-1} \left\{ \frac{e^{-g_6 \sqrt{s}}}{s} \right\} \quad (116)$$

The inverse transformation of equation (116) is carried out by means of the table of integral transforms (5):

$$\begin{aligned} \frac{C(L,t)}{C_1} = \frac{1}{2} \left\{ e^{M-(-g_6 \sqrt{g_7})} \operatorname{Erfc} \left(\frac{g_6}{2 \sqrt{t}} + \sqrt{g_7 t} \right) \right. \\ \left. + e^{M+(-g_6 \sqrt{g_7})} \operatorname{Erfc} \left(\frac{g_6}{2 \sqrt{t}} + \sqrt{g_7 t} \right) \right\} \end{aligned} \quad (117)$$

Substitution of g_6 and g_7 into equation (117) gives equation (29) of text.

The analytical solution corresponding to B.C.(IV) is found to be difficult because it contains the error function and can not be integrated by the method of convolution as shown in equation (113). The solution for a step input corresponding to B.C.(IV) is left in integral form as shown in equation (30) of the text. An approximate method, however, such as integration in a series form could be carried out by expanding the error functions.

APPENDIX IV

Derivations of Frequency Response Functions In Terms of Amplitude Ratios and Phase Shift.

Since the performance functions are defined as the ratio of Fourier transforms of output to input, they are systematically obtained from the transfer functions by replacing s by $j\omega$ as:

$$P.F. = \frac{4\beta''}{(1+\beta'')^2 \exp[-M(1-\beta'')] - (1-\beta'')^2 \exp[-M(1+\beta'')]} \quad (118)$$

$$P.F. = \frac{2\beta''}{(1+\beta'') \exp[-M(1-\beta'')] - (1-\beta'') \exp[-M(1+\beta'')]} \quad (119)$$

$$P.F. = \exp[M(1-\beta'')] \quad (120)$$

$$P.F. = \frac{2 \exp[M(1-\beta'')]}{(1+\beta'')} \quad (121)$$

$$\text{where } \beta'' = \sqrt{1 + \left(\frac{2\theta}{M}\right)(j\omega + k)}$$

The complex number, β'' , of equations (118) through (121) is arranged in polar form as:

$$\beta'' = \left[1 + \frac{2R}{M} + \frac{2(\omega\theta)}{M} \right] = \left[\left(1 + \frac{2R}{M} \right)^2 + \frac{4(\omega\theta)^2}{M^2} \right]^{\frac{1}{4}} e^{ja} \quad (122)$$

$$\text{where } a = \frac{1}{2} \tan^{-1} \left(\frac{2\omega\theta}{2R + M} \right)$$

For the solution corresponding to B.C.(I), equation (122) is substituted to equation (118) to obtain:

$$\frac{C_2(j\omega)}{C_1(j\omega)} = \frac{4\beta'' e^M (\cos a + j \sin a)}{4 \left[1 + \frac{(\omega\theta)}{M} j \right] \sinh(x+jy) + 4b(\cos a + j \sin a) \cosh(x+jy)} \quad (123)$$

Separation of the denominator in equation (123) into a real part and an imaginary part gives:

$$\frac{C_2(j\omega)}{C_1(j\omega)} = \frac{4be^M (\cos a + j \sin a)}{4(\alpha_1 + j \alpha_2)} \quad (124)$$

where α_1 and α_2 are given in the text.

Equation (124) is, then transformed into the polar form as:

$$\frac{C_2(j\omega)}{C_1(j\omega)} = \frac{be^M}{(\alpha_1^2 + \alpha_2^2)^{\frac{1}{2}}} \exp \left\{ j \left[a - \tan^{-1} \left(\frac{\alpha_2}{\alpha_1} \right) \right] \right\} \quad (125)$$

Transformation of equation (125) into the amplitude ratio and phase shift results in equation (31) of the text.

For the solution using B.C.(II), equation (122) is substituted into equation (119) which, upon simplification, gives:

$$\frac{C_2(j\omega)}{C_1(j\omega)} = \frac{2be^M (\cos a + j \sin a)}{\{1+b(e^{ja})\} \exp(x+jy) - \{1-b(e^{ja})\} \exp[-(x+jy)]} \quad (126)$$

Further separation of the denominator in equation (126) into real and imaginary part gives:

$$\frac{C_2(j\omega)}{C_1(j\omega)} = \frac{2be^M (\cos a + j \sin a)}{2(\alpha_3 + \alpha_4)} \quad (127)$$

where α_3 and α_4 are given in the text.

Rearrangement of equation (127) into the polar form gives:

$$\frac{C_2(j\omega)}{C_1(j\omega)} = \frac{be^M}{(\alpha_3^2 + \alpha_4^2)^{\frac{1}{2}}} \exp \left\{ j \left[a - \tan^{-1} \left(\frac{\alpha_4}{\alpha_3} \right) \right] \right\} \quad (128)$$

Equation (128) is then transformed into amplitude ratio and phase shift, which give equation (32) of the text.

Substitution of equation (122) into equation (120) gives the solution for B.C.(III) as:

$$\frac{C_2(j\omega)}{C_1(j\omega)} = \exp M \left[1 - b(\cos a + j \sin a) \right] = e^{M-x} e^{-jy} \quad (129)$$

Equation (33) of the text is obtained by transforming equation (129) in terms of amplitude ratio and the phase shift.

For the solution using B.C.(IV), equation (122) is substituted into equation (121), which upon simplification, gives:

$$\frac{C_2(j\omega)}{C_1(j\omega)} = 2 e^{M-x} \left[\frac{\cos y - j \sin y}{\left(1 + \frac{x}{M}\right) + j(\sin a)} \right] \quad (130)$$

Rearrangement of equation (130) in polar form gives:

$$\frac{C_2(j\omega)}{C_1(j\omega)} = \frac{2 e^{M-x}}{\left[\left(1 + \frac{x}{M}\right)^2 + \sin^2 a \right]^{\frac{1}{2}}} \exp \left[-y - \tan^{-1} \left(\frac{\sin a}{1 + \frac{x}{M}} \right) \right] \quad (131)$$

Equation (131) is then transformed into amplitude ratio and phase shift which give equation (34) of the text.

A DIFFUSION MODEL OF AN
ISOTHERMAL TUBULAR FLOW REACTOR

by

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The general extent of axial dispersion and its effect on the performance of isothermal tubular flow reactors are investigated by use of the diffusion-type model and four sets of following boundary conditions:

$$\text{B.C. (I)} \quad \left\{ \begin{array}{l} \bar{u}C_1 = \bar{u}C_{z \rightarrow 0^+} - \bar{D} \left(\frac{\partial C}{\partial z} \right)_{z \rightarrow 0^+} \\ \left(\frac{\partial C}{\partial z} \right)_{z=L} = 0 \end{array} \right.$$

$$\text{B.C. (II)} \quad \left\{ \begin{array}{l} \bar{u}C_1 = \bar{u}C_{z \rightarrow 0^+} \\ \left(\frac{\partial C}{\partial z} \right)_{z=L} = 0 \end{array} \right.$$

$$\text{B.C. (III)} \quad \left\{ \begin{array}{l} \bar{u}C_1 = \bar{u}C_{z \rightarrow 0^+} \\ \lim_{z \rightarrow \infty} C(z, t) = 0 \end{array} \right.$$

$$\text{B.C. (IV)} \quad \left\{ \begin{array}{l} \bar{u}C_1 = \bar{u}C_{z \rightarrow 0^+} - \bar{D} \left(\frac{\partial C}{\partial z} \right)_{z \rightarrow 0^+} \\ \lim_{z \rightarrow \infty} C(z, t) = 0 \end{array} \right.$$

The differences and the relationship among the steady-state concentration profiles and conversion and unsteady-state transient response curves obtained by use of the different sets of boundary conditions are critically re-examined and numerically evaluated. Only the values obtained by use of B.C.(I) are found to be consistent with the limits of maximum and minimum conversions, i.e. conversion in a plug-flow reactor and in a complete-

ly stirred tank reactor calculated from a one-dimensional transport equation. The solutions obtained by use of the other three sets of boundary conditions are found to be inconsistent with such limits. Hence, unless their limits are recognized, the use of B.C.(II), (III) and (IV) should be avoided. Due to the simple analytical solutions of equations in exponential form, the use of B.C.(III) and (IV), however, offers a convenient means for computing final conversions if their limits are recognized.

An illustration is presented of the method of determining mean dispersion coefficients from dynamic experiments. It is recommended that the calculated values should be compared with the experimental data to determine the feasibility of predicted models and boundary conditions.

The response curves based on different types of transient inputs are presented for all four different sets of boundary conditions both with and without chemical reaction. These curves should be useful for practicing process control engineers and for those who wish to extract the transport parameter from dynamic response experiments.

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